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**Differential Geometry** 

# Isometric deformations of the $\mathcal{K}^{\frac{1}{4}}$ -flow translators in $\mathbb{R}^{3}$ with helicoidal symmetry $\stackrel{\text{\tiny{translators}}}{\longrightarrow}$



Déformations isométriques des translateurs du flot  $\mathcal{K}^{1/4}$  dans  $\mathbb{R}^3$  à symétrie hélicoïdale

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#### ABSTRACT

The height functions of  $\mathcal{K}^{\frac{1}{4}}$ -flow translators in the Euclidean space  $\mathbb{R}^3$  solve the classical Monge–Ampère equation  $f_{xx}f_{yy} - f_{xy}^2 = 1$ . We explicitly and geometrically determine the moduli space of all helicoidal  $\mathcal{K}^{\frac{1}{4}}$ -flow translators, which are generated from planar curves by the action of helicoidal groups.

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# RÉSUMÉ

Les fonctions de hauteur des translateurs du flot  $\mathcal{K}^{1/4}$  de  $\mathbb{R}^3$  résolvent l'équation de Monge-Ampère classique  $f_{xx}f_{yy} - f_{xy}^2 = 1$ . Nous déterminons de manière géométrique explicite l'espace des modules de tous les translateurs à symétrie hélicoïdale du flot  $\mathcal{K}^{1/4}$ , qui sont engendré à partir de courbes planes par l'action de groupes hélicoïdaux.

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#### 1. Motivation and main results

#### 1.1. Introduction

The recent decades saw intensive research devoted to the study of various geometric flows and soliton solutions. The classical curve-shortening flow admits fruitful generalizations with intriguing applications. One of Huisken's theorems guarantees that an analogue of the Gage–Hamilton's shrinking-curves theorem in the plane also holds for the mean curvature flow in higher dimensional Euclidean spaces.

Chow [6] studied the  $\mathcal{K}^{\alpha}$ -flow, which is the normal deformation by powers of the Gauss curvature. Given a smooth immersion  $\mathcal{F}_0: \Sigma \to \mathbb{R}^{n+1}$  of a strictly convex hypersurface  $\Sigma$  in Euclidean space  $\mathbb{R}^{n+1}$ , the solution of the initial value problem for the  $\mathcal{K}^{\alpha}$ -flow means a one-parameter family of smooth immersions  $\{\mathcal{F}_t = \mathcal{F}(\cdot, t): \Sigma \to \mathbb{R}^{n+1}\}_{t \in [0,T)}$  satisfying the geometric evolution:

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$$\begin{cases} \frac{\partial}{\partial t} \mathcal{F}(\mathbf{p}, t) = -\mathcal{K}(\mathbf{p}, t)^{\alpha} v(\mathbf{p}, t), & (\mathbf{p}, t) \in \Sigma \times [0, T) \\ \mathcal{F}(\mathbf{p}, 0) = \mathcal{F}_0(\mathbf{p}), \end{cases}$$

where  $\alpha > 0$  is a constant,  $\nu(\mathbf{p}, t)$  denotes the outward-pointing unit normal of  $\mathcal{F}(\mathbf{p}, t)$ , and the Gauss–Kronecker curvature functional  $\mathcal{K}(\mathbf{p}, t)$  is defined as the product of the principal curvatures. The Gauss curvature flow means the  $\mathcal{K}^1$ -flow which was investigated by Tso [23]. The case when  $(n, \alpha) = (2, 1)$  was originally introduced by Firey [11] in his study of a model of the wearing process of convex stones on a beach by water waves. In 1999, Andrews [3] established Firey's conjecture that convex surfaces evolving by the Gauss curvature flow become spherical.

For several motivations for the study of hypersurfaces moving by their Gauss curvature, we refer to [4, Section 1]. In particular, the  $\mathcal{K}^{\frac{1}{n+2}}$ -flow in Euclidean space  $\mathbb{R}^{n+1}$  admits a deep and interesting geometric meaning in affine differential geometry. As indicated in [2, Section 1], the affine-invariant evolution of convex hypersurfaces in  $\mathbb{R}^{n+1}$  under the so-called affine normal flow can be reformulated as the  $\mathcal{K}^{\frac{1}{n+2}}$ -flow modulo diffeomorphisms. It is worth to mention that the geometric meaning of the mysterious factor  $\mathcal{K}^{\frac{1}{n+2}}$  in the classical affine differential geometry is also well-described in [9] and [5, Section 1] with details.

Urbas [24] investigated self-similar and translating solitons for the normal evolution by positive powers of the Gauss curvature. The simplest example of a translating soliton is Calabi's grim reaper  $y = \ln(\cos x)$ , which moves by downward translation under the  $\mathcal{K}^1$ -flow in the plane  $\mathbb{R}^2$ . In this paper, we say that a surface  $\Sigma$  in  $\mathbb{R}^3$  is a  $\mathcal{K}^{\frac{1}{4}}$ -translator when we have the geometric condition:  $\mathcal{K}_{\Sigma} = \cos^4(\theta_{\Sigma})$ . Here, the scalar function  $\mathcal{K}_{\Sigma}$  denotes the Gaussian curvature, and the third component  $\cos(\theta_{\Sigma}) = \mathbf{n}_{\Sigma} \cdot (0, 0, 1)$  of the unit normal  $\mathbf{n}_{\Sigma}$  is called the angle function on  $\Sigma$ . When the initial surface in  $\mathbb{R}^3$  is a  $\mathcal{K}^{\frac{1}{4}}$ -translator, it moves by vertical translation under the normal evolution of the  $\mathcal{K}^{\frac{1}{4}}$ -flow [24, Section 4].

The  $\mathcal{K}^{\frac{1}{4}}$ -translators in the Euclidean space  $\mathbb{R}^3$  are of significant geometrical interest. The convex graph z = f(x, y) becomes a  $\mathcal{K}^{\frac{1}{4}}$ -translator if and only if its height function f solves the classical Monge–Ampère equation:

$$f_{xx}f_{yy} - f_{xy}^2 = 1.$$

Jörgens' outstanding holomorphic resolution [16] says that, when  $f_{xx}f_{yy} - f_{xy}^2 = 1$ , the gradient graph  $(x, y, f_x, f_y)$  becomes a minimal surface in the Euclidean space  $\mathbb{R}^4$ . The Hessian one equation is a special case of special Lagrangian equations [14], *split* special Lagrangian equations [15,19,20], and affine mean curvature equations [2,5,22]. Furthermore, its solutions induce flat surfaces in hyperbolic space  $\mathbb{H}^3$  [21].

1.2. Isometric deformations of helicoidal  $\mathcal{K}^{\frac{1}{4}}$ -translators

**Theorem 1** (Moduli space of  $\mathcal{K}^{\frac{1}{4}}$ -translators with rotational or helicoidal symmetry).

- (A) Any helicoidal  $\mathcal{K}^{\frac{1}{4}}$ -translator  $\Sigma$  of pitch  $\mu$  admits a one-parameter family of isometric helicoidal  $\mathcal{K}^{\frac{1}{4}}$ -translators  $\Sigma^{h}$  with pitch h such that  $\Sigma = \Sigma^{\mu}$  and that  $\Sigma^{0}$  is rotational.
- (B) The cylinder over a circle in the xy-plane is a rotational  $\mathcal{K}^{\frac{1}{4}}$ -translator. Additionally, there exists a one-parameter family of  $\mathcal{K}^{\frac{1}{4}}$ -translators  $\mathcal{H}_c$  invariant under the rotation with z-axis. The profile curve of rotational surface  $\mathcal{H}_c$  is congruent to the graph  $(U, 0, \Lambda_c(U))$ , where the one-parameter family of height functions  $\Lambda_c(U)$  is explicitly given by:

$$\Lambda_{c}(U) = \begin{cases} \frac{1}{2} [U\sqrt{U^{2} + \kappa^{2}} + \kappa^{2} \operatorname{arcsinh}(\frac{U}{\kappa})], & U > 0 \text{ (when } c = 1 + \kappa^{2}, \ \kappa > 0), \\ \frac{1}{2} U^{2}, & U \ge 0 \text{ (when } c = 1), \\ \frac{1}{2} [U\sqrt{U^{2} - \kappa^{2}} - \kappa^{2} \operatorname{arccosh}(\frac{U}{\kappa})], & U > \kappa \text{ (when } c = 1 - \kappa^{2}, \ \kappa > 0). \end{cases}$$

- (C) There exists a two-parameter family of helicoidal  $\mathcal{K}^{\frac{1}{4}}$ -translators  $\mathcal{H}^{h}_{c}$  and the geometric coordinates (U, t) on  $\mathcal{H}^{h}_{c}$  satisfying the following conditions.
  - (C1) The geometric meaning of parameter h is that the surface  $\mathcal{H}_c^h$  is invariant under the helicoidal motion with pitch h. The surface  $\mathcal{H}_c^h$  is invariant under the one-parameter subgroup  $\{\mathbf{S}_T\}$  of the group of rigid motions of  $\mathbb{R}^3 = \mathbb{C} \times \mathbb{R}$  given by:

$$(\zeta, z) \in \mathbb{C} \times \mathbb{R} \mapsto \mathbf{S}_T(\zeta, z) = (e^{iT}\zeta, hT + z) \in \mathbb{C} \times \mathbb{R}.$$

- (C2) There exist the coordinates (U, t) on the helicoidal surface  $\mathcal{H}_c^h$  such that its metric reads  $I_{\mathcal{H}_c^h} = (U^2 + c) dU^2 + U^2 dt^2$ .
- (C3) The geometric meaning of parameter c is the property that the helicoidal surface  $\mathcal{H}_c^h$  is isometric to the rotational surface  $\mathcal{H}_c^0 = \mathcal{H}_c$ .
- (C4) The geometric meaning of coordinate U is the property that the function  $\frac{1}{\sqrt{U^2+c}}$  coincides with the angle function on the surface  $\mathcal{H}_c^h$  up to a sign.

The statement (A) in Theorem 1 is inspired by the 1982 do Carmo–Dajczer theorem [8] that a surface of non-zero constant mean curvature is helicoidal if and only if it lies in the associate family [18] of a Delaunay's rotational surface [10,17] with the same constant mean curvature. In 1998, Haak [12] presented an alternative proof of the do Carmo–Dajczer theorem.

The *mean curvature flow* in  $\mathbb{R}^3$  also admits the translating solitons with helicoidal symmetry. In 1994, Altschuler and Wu [1] showed the existence of the convex, rotational, entire graphical translator. In 2007, Clutterbuck, Schnürer and Schulze [7] constructed the bigraphical translator, which is also rotationally symmetric.

**Open problem.** Prove or disprove that Halldorsson's helicoidal translators [13] for the mean curvature flow admit the isometric deformation from rotational translators.

# 2. Proof of Theorem 1

We first need to revisit Bour's construction [8] with details to specify the behavior of the angle function on his isometric helicoidal surfaces.

**Lemma 2** (Angle function on Bour's helicoidal surfaces). Let  $\Sigma$  be a helicoidal surface with pitch vector  $\mu \mathbf{k} = (0, 0, \mu)$  and the generating curve  $\gamma = (\mathcal{R}, 0, \Lambda)$  in the xz-plane, which admits the parameterization  $(u, \theta) \mapsto (\mathcal{R} \cos \theta, \mathcal{R} \sin \theta, \Lambda + \mu \theta)$ , where u denotes a parameter of the generating curve  $\gamma$ . We then define the Bour coordinate transformation:

$$(u, \theta) \mapsto (s, t) = (s, \theta + \Theta),$$

via the relations

$$ds^{2} = d\mathcal{R}^{2} + \frac{\mathcal{R}^{2}}{\mathcal{R}^{2} + \mu^{2}} d\Lambda^{2}$$
$$d\Theta = \frac{\mu}{\mathcal{R}^{2} + \mu^{2}} d\Lambda,$$

and also introduce the Bour function U using the relation  $U^2 = R^2 + \mu^2$ .

(A) The helicoidal surface  $\Sigma$  admits the reparametrization satisfying (A1), (A2), and (A3):

$$(s,t) \mapsto \mathbf{X}(s,t) = (\mathcal{R}\cos(t-\Theta), \mathcal{R}\sin(t-\Theta), \Lambda + \mu(t-\Theta))$$

- (A1) Its first fundamental form reads  $I_{\Sigma} = ds^2 + U^2 dt^2$ .
- (A2) The parameters R,  $\Lambda$ , and  $\Theta$  can be recovered from the Bour function U explicitly:

$$\begin{cases} \mathcal{R}^{2} = U^{2} - \mu^{2}, \\ d\Lambda^{2} = \frac{U^{2}}{(U^{2} - \mu^{2})^{2}} \left( U^{2} \left( 1 - \left( \frac{dU}{ds} \right)^{2} \right) - h^{2} \right) ds^{2}, \\ d\Theta = \frac{\mu}{U^{2}} d\Lambda. \end{cases}$$

(A3) The angle function  $\mathbf{n}_3$  defined as the third component  $\mathbf{n} \cdot \mathbf{k}$  of the induced unit normal  $\mathbf{n} = \frac{1}{\|\mathbf{X}_s \times \mathbf{X}_t\|} \mathbf{X}_s \times \mathbf{X}_t$  is also determined by the Bour function U.

$$n_3^2 = \left(\frac{\mathrm{d}U}{\mathrm{d}s}\right)^2.$$

(B) We construct a two-parameter family of helicoidal surfaces  $\Sigma^{\lambda,h}$  of pitch h by the patch:

$$\mathbf{X}^{\lambda,h}(s,t) = \left(\mathcal{R}^{\lambda,h}\cos\left(\frac{t}{\lambda} - \Theta^{\lambda,h}\right), \mathcal{R}^{\lambda,h}\sin\left(\frac{t}{\lambda} - \Theta^{\lambda,h}\right), \Lambda^{\lambda,h} + h\left(\frac{t}{\lambda} - \Theta^{\lambda,h}\right)\right),$$

where the geometric datum  $(\mathcal{R}^{\lambda,h}, \Lambda^{\lambda,h}, \Theta^{\lambda,h})$  is explicitly determined by the pair  $(\lambda, h)$  of constants and the Bour function U(s) arising from the reparametrization  $\mathbf{X}(s, t)$  of  $\Sigma$ 

$$\begin{cases} \left(\mathcal{R}^{\lambda,h}\right)^2 = \lambda^2 U^2 - h^2, \\ \left(d\Lambda^{\lambda,h}\right)^2 = \frac{\lambda^2 U^2}{(\lambda^2 U^2 - h^2)^2} \left(\lambda^2 U^2 \left(1 - \lambda^2 \left(\frac{dU}{ds}\right)^2\right) - h^2\right) ds^2, \\ d\Theta^{\lambda,h} = \frac{h}{\lambda^2 U^2} d\Lambda^{\lambda,h}. \end{cases}$$
(2.1)

Then, the helicoidal surface  $\Sigma^{\lambda,h}$  is isometric to the initial surface  $\Sigma$ , and its angle function  $n_3^{\lambda,h} = \mathbf{n}^{\lambda,h} \cdot \mathbf{k}$  is determined by the Bour function U of the initial surface  $\Sigma$ .

$$\left(n_3^{\lambda,h}\right)^2 = \lambda^2 \left(\frac{\mathrm{d}U}{\mathrm{d}s}\right)^2.$$

(C) Furthermore, the helicoidal surface  $\Sigma^{1,\mu}$  coincides with the initial surface  $\Sigma$ .

**Proof.** (A) The definitions of the Bour coordinate (s, t) and the Bour function U yield:

$$\begin{split} I_{\Sigma} &= \left( \mathrm{d}R^2 + \mathrm{d}\Lambda^2 \right) + 2\mu \, \mathrm{d}\Lambda \, \mathrm{d}\theta + \left( \mathcal{R}^2 + \mu^2 \right) \mathrm{d}\theta^2 \\ &= \left( \mathrm{d}\mathcal{R}^2 + \frac{\mathcal{R}^2}{\mathcal{R}^2 + \mu^2} \, \mathrm{d}\Lambda^2 \right) + \left( \mathcal{R}^2 + \mu^2 \right) \left( \mathrm{d}\theta + \frac{\mu}{\mathcal{R}^2 + \mu^2} \, \mathrm{d}\Lambda \right)^2 \\ &= \mathrm{d}s^2 + U^2 \, \mathrm{d}t^2. \end{split}$$

Noticing that the definition  $U^2 = R^2 + \mu^2$  implies  $d\mathcal{R}^2 = \frac{U^2}{U^2 - \mu^2} dU^2$ , we can recover the function  $\dot{A} = \frac{dA}{ds}$  from the Bour function U(s) explicitly:

$$ds^{2} = d\mathcal{R}^{2} + \frac{\mathcal{R}^{2}}{\mathcal{R}^{2} + \mu^{2}} d\Lambda^{2} = \frac{U^{2}}{U^{2} - \mu^{2}} dU^{2} + \frac{U^{2} - \mu^{2}}{U^{2}} d\Lambda^{2}$$

and

$$d\Lambda^{2} = \frac{U^{2}}{U^{2} - \mu^{2}} \left( ds^{2} - \frac{U^{2}}{U^{2} - \mu^{2}} dU^{2} \right) = \frac{U^{2}}{(U^{2} - \mu^{2})^{2}} \left( U^{2} \left( 1 - \left( \frac{dU}{ds} \right)^{2} \right) - h^{2} \right) ds^{2}.$$

Adopting the symbol  $\dot{} = \frac{d}{ds}$  again, we obtain:

$$\mathbf{X}_{s} \times \mathbf{X}_{t} = (\mu \dot{\mathcal{R}} \sin \theta - \mathcal{R} \dot{\Lambda} \cos \theta, -\mu \dot{\mathcal{R}} \cos \theta - \mathcal{R} \dot{\Lambda} \sin \theta, \mathcal{R} \dot{\mathcal{R}}).$$

After setting  $I_{\Sigma} := E ds^2 + 2F ds dt + G dt^2 = ds^2 + U^2 dt^2$ , we immediately see that:  $\|\mathbf{X}_s \times \mathbf{X}_t\|^2 = EG - F^2 = U^2$ . It thus follows that:

$$n_3^2 = \frac{\left(\mathcal{R}\dot{\mathcal{R}}\right)^2}{U^2} = \dot{U}^2 = \left(\frac{\mathrm{d}U}{\mathrm{d}s}\right)^2$$

(B) We first show that the surface  $\Sigma^{\lambda,h}$  is isometric to the initial surface  $\Sigma$ . Let us write:  $I_{\Sigma^{\lambda,h}} = E^{\lambda,h} ds^2 + 2F^{\lambda,h} ds dt + G^{\lambda,h} dt^2$ . Adopting the symbol  $\dot{} = \frac{d}{ds}$  and using (2.1), we have:

$$\begin{split} E^{\lambda,h} &= \left(\dot{\mathcal{R}}^{\lambda,h}\right)^2 + \mathcal{R}^2 \left(\dot{\Theta}^{\lambda,h}\right)^2 + \left(\dot{A}^{\lambda,h} - h\dot{\Theta}^{\lambda,h}\right)^2 = \left(\dot{\mathcal{R}}^{\lambda,h}\right)^2 + \frac{\lambda^2 U^2 - h^2}{\lambda^2 U^2} \left(\dot{\Theta}^{\lambda,h}\right)^2 \\ &= \frac{\lambda^2 U^2 \dot{U}^2}{\lambda^2 U^2 - h^2} + \frac{\lambda^2 U^2 - h^2}{\lambda^2 U^2} \cdot \frac{\lambda^2 U^2 [\lambda^2 U^2 (1 - \lambda^2 \dot{U}^2) - h^2]}{(\lambda^2 U^2 - h^2)^2} = 1. \end{split}$$

We also deduce:

$$F^{\lambda,h} = -\frac{1}{\lambda} \left[ \left( \left( \mathcal{R}^{\lambda,h} \right)^2 + h^2 \right) \dot{\Theta} - h \dot{A} \right] = -\frac{1}{\lambda} \left[ \lambda^2 U^2 \dot{\Theta} - h \dot{A} \right] = 0,$$

and

$$G^{\lambda,h} = \frac{1}{\lambda^2} \left[ \left( \mathcal{R}^{\lambda,h} \right)^2 + h^2 \right] = U^2$$

Combining these, we meet  $I_{\Sigma^{\lambda,h}} = E^{\lambda,h} ds^2 + 2F^{\lambda,h} ds dt + G^{\lambda,h} dt^2 = ds^2 + U^2 dt^2 = I_{\Sigma}$ . Now, it remains to determine the angle function of the surface  $\Sigma^{\lambda,h}$ . Adopting the new variable  $\theta = \frac{t}{\lambda} - \Theta^{\lambda,h}$  for simplicity, we write:

$$\mathbf{X}_{s}^{\lambda,h} \times \mathbf{X}_{t}^{\lambda,h} = \frac{1}{\lambda} \left( h \dot{\mathcal{R}}^{\lambda,h} \sin \theta - \mathcal{R}^{\lambda,h} \dot{\Lambda} \cos \theta, -h \dot{\mathcal{R}}^{\lambda,h} \cos \theta - \mathcal{R}^{\lambda,h} \dot{\Lambda} \sin \theta, \mathcal{R}^{\lambda,h} \dot{\mathcal{R}}^{\lambda,h} \right)$$

Taking account into this and the equality  $\|\mathbf{X}_{s}^{\lambda,h} \times \mathbf{X}_{t}^{\lambda,h}\|^{2} = E^{\lambda,h}G^{\lambda,h} - (F^{\lambda,h})^{2} = U^{2}$ , we meet

$$(n_3^{\lambda,h})^2 = (\mathbf{n}^{\lambda,h} \cdot \mathbf{k})^2 = \frac{1}{U^2} \cdot \frac{(\mathcal{R}^{\lambda,h})^2 (\dot{\mathcal{R}}^{\lambda,h})^2}{\lambda^2} = \lambda^2 \dot{U}^2 = \lambda^2 \left(\frac{\mathrm{d}U}{\mathrm{d}s}\right)^2.$$

(C) The datum  $(\mathcal{R}^{1,\mu}, \Lambda^{1,\mu}, \Theta^{1,\mu})$  of  $\Sigma^{1,\mu}$  coincides with the datum  $(\mathcal{R}, \Lambda, \Theta)$  of  $\Sigma$ .  $\Box$ 

We briefly sketch the geometric ingredients in our construction in Theorem 1. For given a helicoidal  $\mathcal{K}^{\frac{1}{4}}$ -translator, we prove that there exists a sub-family chosen from the two-parameter family of Bour's isometric helicoidal surfaces, so that each member of this sub-family is a  $\mathcal{K}^{\frac{1}{4}}$ -translator and that one member is rotationally symmetric.

Our one-parameter family of  $\mathcal{K}^{\frac{1}{4}}$ -translators admits the parametrizations by so-called the Bour coordinate (s, t) and the Bour function U = U(s). The trick to obtain the explicit construction in (C) of Theorem 1 is to perform the coordinate transformation  $s \mapsto U$  to have the geometric coordinate (U, t) on our one-parameter family of  $\mathcal{K}^{\frac{1}{4}}$ -translators.

**Lemma 3** (Existence of helicoidal  $\mathcal{K}^{\frac{1}{4}}$ -translators of pitch h). Let h be a given constant. Then, any non-cylindrical helicoidal  $\mathcal{K}^{\frac{1}{4}}$ -translator with pitch h admits the parameterization:

$$(U,t)\mapsto \big(\mathcal{R}(U)\cos\big(t-\Theta(U)\big),\ \mathcal{R}(U)\sin\big(t-\Theta(U)\big),\ \Lambda(U)+h\big(t-\Theta(U)\big)\big),$$

where the geometric datum  $(\mathcal{R}(U), \Lambda(U), \Theta(U))$  can be obtained from the relation:

$$\begin{cases} \mathcal{R}(U)^{2} = U^{2} - h^{2}, \\ \left(\frac{dA}{dU}\right)^{2} = \frac{U^{2}}{\left(U^{2} - h^{2}\right)^{2}} \left[U^{4} + \left(c - 1 - h^{2}\right)U^{2} - h^{2}c\right], \\ \left(\frac{d\Theta}{dU}\right)^{2} = \frac{h^{2}}{U^{2}\left(U^{2} - h^{2}\right)^{2}} \left[U^{4} + \left(c - 1 - h^{2}\right)U^{2} - h^{2}c\right], \end{cases}$$
(2.2)

where  $c \in \mathbb{R}$  is a constant.

**Proof.** Taking  $\lambda = 1$  in Lemma 2, we construct a helicoidal surface  $\Sigma$  with pitch *h*:

$$(s,t) \mapsto \mathbf{X}^{1,h}(s,t) = (\mathcal{R}\cos(t-\Theta), \mathcal{R}\sin(t-\Theta), \Lambda + h(t-\Theta)),$$

where the geometric datum  $(\mathcal{R}, \Lambda, \Theta) = (\mathcal{R}(s), \Lambda(s), \Theta(s))$  is given by the relation:

$$\begin{cases} \mathcal{R}^{2} = U^{2} - h^{2}, \\ (d\Lambda)^{2} = \frac{U^{2}}{(U^{2} - h^{2})^{2}} \left( U^{2} \left( 1 - \left( \frac{dU}{ds} \right)^{2} \right) - h^{2} \right) ds^{2}, \\ d\Theta = \frac{h}{U^{2}} d\Lambda. \end{cases}$$
(2.3)

The key point is to take the Bour function U as the new parameter on our helicoidal surface  $\Sigma$ . According to Lemma 2 again, we see that the induced metric on  $\Sigma$  reads  $I_{\Sigma} = ds^2 + U^2 dt^2$ , that its Gaussian curvature K is equal to  $K = -\frac{1}{U} \frac{d^2 U}{ds^2}$ , and that its angle function reads  $n_3^2 = (\frac{dU}{ds})^2$ . Thus, the condition that the helicoidal surface  $\Sigma$  becomes a  $\mathcal{K}^{\frac{1}{4}}$ -translator implies that  $K = n_3^4$ , which means the ordinary differential equation:

$$-\frac{1}{U}\frac{\mathrm{d}^2 U}{\mathrm{d}s^2} = \left(\frac{\mathrm{d}U}{\mathrm{d}s}\right)^4.$$

In the case when  $\frac{dU}{ds}$  vanishes locally, our surface  $\Sigma$  becomes the cylinder over a circle in the *xy*-plane. When  $\frac{dU}{ds}$  does not vanish, we are able to make a coordinate transformation  $s \mapsto U$  and can rewrite the above ODE as:  $0 = \frac{d}{ds} (1/(\frac{dU}{ds})^2 - U^2)$ . Hence its first integral is explicitly given by, for some constant  $c \in \mathbb{R}$ ,  $ds^2 = (U^2 + c) dU^2$ . We now can employ this to perform the coordinate transformation  $(s, t) \mapsto (U, t)$  on  $\Sigma$ . Rewriting (2.3) in terms of the new variable U gives indeed the relation in (2.2).  $\Box$ 

**Proof of Theorem 1.** We first prove (B). Taking h = 0 in Lemma 3, we see that any rotational  $\mathcal{K}^{\frac{1}{4}}$ -translator admits the patch:

$$(U,t)\mapsto \big(\mathcal{R}(U)\cos(t-\Theta(U)),\ \mathcal{R}(U)\sin(t-\Theta(U)),\ \Lambda(U)+h(t-\Theta(U))\big),$$

where the geometric datum ( $\mathcal{R}(U), \Lambda(U), \Theta(U)$ ) satisfies the relation:

$$\left(\mathcal{R}(U)\right)^2 = U^2, \qquad \left(\frac{\mathrm{d}\Lambda}{\mathrm{d}U}\right)^2 = U^2 + (c-1), \qquad \left(\frac{\mathrm{d}\Theta}{\mathrm{d}U}\right)^2 = 0$$

for some constant  $c \in \mathbb{R}$ . The condition that the helicoidal surface  $\Sigma$  becomes a  $\mathcal{K}^{\frac{1}{4}}$ -translator implies the ordinary differential equation:

$$-\frac{1}{U}\frac{\mathrm{d}^2 U}{\mathrm{d}s^2} = \left(\frac{\mathrm{d}U}{\mathrm{d}s}\right)^4$$

When  $\frac{dU}{ds}$  vanishes locally, our surface  $\Sigma$  becomes the cylinder over a circle in the *xy*-plane. In the case when  $\frac{dU}{ds}$  does not vanish, we can introduce a coordinate transformation  $s \mapsto U$ . Since  $\frac{d\Theta}{dU}$  vanishes, without loss of generality, after a translation of the coordinate *t*, we may take  $\Theta = 0$  in the above patch as follows:

$$(U,t) \mapsto (U\cos t, U\sin t, \Lambda(U)).$$

As in the proof of Lemma 3,  $\Lambda(U)$  solves the ordinary differential equation:  $\frac{d\Lambda}{dU} = \pm \sqrt{U^2 + (c-1)}$ . Considering the sign of the constant c - 1, we meet the explicit solution  $\Lambda_c(U) = \Lambda(U)$  (up to the sign) as follows:

$$\Lambda(U) = \begin{cases} \frac{1}{2} [U\sqrt{U^2 + \kappa^2} + \kappa^2 \operatorname{arcsinh}(\frac{U}{\kappa})] & \text{(when } c = 1 + \kappa^2, \ \kappa > 0), \\ \frac{1}{2} U^2 & \text{(when } c = 1), \\ \frac{1}{2} [U\sqrt{U^2 - \kappa^2} - \kappa^2 \operatorname{arccosh}(\frac{U}{\kappa})] & \text{(when } c = 1 - \kappa^2, \ \kappa > 0). \end{cases}$$

We next prove (A). Using Lemma 2, we see that, for a given helicoidal  $\mathcal{K}^{\frac{1}{4}}$ -translator  $\Sigma$ , we are able to introduce the Bour coordinate (s, t) and the Bour function U(s) on the surface  $\Sigma$  so that  $I_{\Sigma} = ds^2 + U(s)^2 dt^2$ . The condition that  $\Sigma$  is a  $\mathcal{K}^{\frac{1}{4}}$ -translator says:

$$-\frac{1}{U}\frac{\mathrm{d}^2 U}{\mathrm{d}s^2} = \left(\frac{\mathrm{d}U}{\mathrm{d}s}\right)^4,\tag{2.4}$$

just as we saw in the proof of Lemma 3. Next, by Lemma 2 again, we can associate a one-parameter family of isometric helicoidal surfaces  $\Sigma^h$  satisfying three conditions:  $\Sigma = \Sigma^{\mu}$ ,  $I_{\Sigma^h} = I_{\Sigma}$ , and the angle function on  $\Sigma^h$  coincides with the one on  $\Sigma$ . Hence, as we saw in the proof of Lemma 3, the above ordinary differential equation in (2.4) guarantees that any helicoidal surface  $\Sigma^h$  becomes indeed a  $\mathcal{K}^{\frac{1}{4}}$ -translator.

It now remains to show (C). The statement (C1) is obvious by the construction in Lemma 3. Next, the equality  $ds^2 = (U^2 + c) dU^2$  proved in Lemma 3 implies that the induced metric of the helicoidal surface constructed in Lemma 3 reads:  $ds^2 + U^2 dt^2 = (U^2 + c) dU^2 + U^2 dt^2$  (which implies (C2) and (C3)), and that the angle function is given by, up to a sign:

$$\frac{\mathrm{d}U}{\mathrm{d}s} = \frac{1}{\frac{\mathrm{d}s}{\mathrm{d}U}} = \frac{1}{\sqrt{U^2 + c}},$$

which is (C4). This completes the proof of our description of the moduli space of helicoidal  $\mathcal{K}^{\frac{1}{4}}$ -translators in Theorem 1.

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