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# $L^{2}$ estimates and existence theorems for $\bar{\partial}_{b}$ on Lipschitz boundaries of $Q$-pseudoconvex domains 

Sayed Saber ${ }^{a}$<br>${ }^{a}$ Department of Mathematics and Computer Science, Faculty of Science, Beni-Suef University, Egypt.<br>E-mail: sayedkay@yahoo.com.

This research is dedicated to the memory of Professor Osama Abdelkader


#### Abstract

On a bounded $q$-pseudoconvex domain $\Omega$ in $\mathbb{C}^{n}$ with Lipschitz boundary $b \Omega$, we prove the $L^{2}$ existence theorems of the $\bar{\partial}_{b}$-operator on $b \Omega$. This yields the closed range property of $\bar{\partial}_{b}$ and its adjoint $\bar{\partial}_{b}^{*}$. As an application, we establish the $L^{2}$-existence theorems and regularity theorems for the $\bar{\partial}_{b}$-Neumann operator.


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## Introduction and main results

Let $\Omega$ be a bounded domain in $\mathbb{C}^{n}$ with smooth boundary $b \Omega$. The Cauchy-Riemann operators $\bar{\partial}$ on $\mathbb{C}^{n}$ induce the tangential Cauchy-Riemann complex or $\bar{\partial}_{b}$ complex on $b \Omega$. On the boundaries of smooth bounded domains, there are several equivalent ways of defining the $\bar{\partial}_{b}$ complex. The $\bar{\partial}_{b}$ complex was first formulated by J. J. Kohn and H. Rossi in [26] for smooth boundaries to understand the holomorphic extension of CR-functions from the boundaries of complex manifolds. On a strictly pseudoconvex domain with smooth boundary in $\mathbb{C}^{n}$, the $\bar{\partial}_{b}$-complex has been studied in several articles (cf. [2,5, $8,17,18,27]$ ). In the case of a weakly pseudoconvex domain with smooth boundary in $\mathbb{C}^{n}$, the $L^{2}$ and Sobolev estimates for $\bar{\partial}_{b}$ have been obtained by M.-C. Shaw in [33] for $1 \leqslant q<n-1$ and by H. P. Boas and M.-C. Shaw in [3] for $q=n-1$ (see also J. J. Kohn [25]). On the boundary of a weakly pseudoconvex domain, it was pointed out by J. P. Rosay in [32] that one can combine the results of J. J. Kohn and H. Rossi in [26] with those of J. J. Kohn in [24] to prove that the global solutions to the equation $\bar{\partial}_{b} u=f$ exists. Other results in this direction see Andreea C. Nicoara [31] and Phillip S. Harrington and Andrew Raich [14].

When the boundary is only Lipschitz, not every definition can be appropriately extended. On a Lipschitz boundary of a bounded domain in $\mathbb{C}^{n}$, the complex normal vector is defined almost everywhere on $b \Omega$. It was pointed out by D. Sullivan in [39] (see also N. Teleman in [40]) that
on a real Lipschitz manifold, $q$-forms with $L^{2}$ coefficients and the de Rham complex are still well defined. Thus one can still define ( $p, q$ ) -forms with $L^{2}(b \Omega)$ coefficients, denoted by $L_{p, q}^{2}(b \Omega)$. The $\bar{\partial}_{b}$ complex is then well defined as a closed densely defined operator from $L_{p, q-1}^{2}(b \Omega)$ to $L_{p, q}^{2}(b \Omega)$. In [13], Phillip S. Harrington has constructed a compact solution operator to the $\bar{\partial}_{b}$-operator on a pseudoconvex domain with Lipschitz boundary. On the same domain, the $L^{2}$ existence theorems of the $\bar{\partial}_{b}$-operator was established by Mei-Chi Shaw in [37]. The first purpose of the paper is to extend this result to Lipschitz boundaries of $q$-pseudoconvex domains. Our first main result is the following:

Theorem 1. Let $\Omega \Subset \mathbb{C}^{n}$ be a bounded $q$-pseudoconvex domain with Lipschitz boundary $b \Omega$ and let $1 \leqslant q \leqslant n$. For every $\alpha \in L_{p, q}^{2}(b \Omega)$, where $0 \leqslant p \leqslant n, 1 \leqslant q<n-1, n \geqslant 2$ such that

$$
\bar{\partial}_{b} \alpha=0 \text { on } b \Omega,
$$

there exists $a u \in L_{p, q-1}^{2}(b \Omega)$ satisfying $\bar{\partial}_{b} u=\alpha$ in the distribution sense in $b \Omega$. Moreover, there exists a constant $C$ depending only on the diameter and the Lipschitz constant of $\Omega$ but is independent of $\alpha$ such that

$$
\|u\|_{b \Omega} \leqslant C\|\alpha\|_{b \Omega} .
$$

When $q=n-1$, for every $\alpha \in L_{p, n-1}^{2}(b \Omega)$ satisfies

$$
\int_{b \Omega} \alpha \wedge \phi \mathrm{~d} S=0, \text { for any } \phi \in C_{n-p, 0}^{\infty}(\bar{\Omega}) \cap \operatorname{ker} \bar{\partial}
$$

the same conclusion holds.
The proof of the main theorem consists of three parts: first we prove the existence and the boundedness of the $\bar{\partial}$-Neumann operator $N$ on Sobolev spaces $W^{m}(\Omega)$ for $-\frac{1}{2} \leqslant m \leqslant \frac{1}{2}$. This yields that the operators $\bar{\partial} N$ and $\bar{\partial}^{*} N$ and the Bergman projection $P$ are bounded operators on $W^{m}(\Omega)$. Second, we study the solvability of the $\bar{\partial}$-problem in the Sobolev space $W^{m}(\Omega)$ with prescribed support in $\bar{\Omega}$, for $-\frac{1}{2} \leqslant m \leqslant \frac{1}{2}$. Third, by using the jump formula derived from the Bochner-Martinelli-Koppelman kernel, the main result follows.

The closed range property is related to existence and regularity results for $\bar{\partial}_{b}$. Independently, when $b \Omega$ is smooth and weakly pseudoconvex in $\mathbb{C}^{n}$, Mei-Chi Shaw in [33] and H. P. Boas and MeiChi Shaw in [3] proved that the range of $\bar{\partial}_{b}$ was closed on $(p, q)$-forms of degrees $1 \leqslant q<n-1$ and $q=n-1$, respectively. On a boundary of strongly pseudoconvex domain, the range of $\bar{\partial}_{b}$ is closed follows from J. J. Kohn and H. Rossi [26]. If $\Omega$ is Lipschitz pseudoconvex in $\mathbb{C}^{n}$ and if there exists a plurisubharmonic defining function in a neighborhood of $\bar{\Omega}$, the range of $\bar{\partial}_{b}$ is closed follows by Mei-Chi Shaw [37]. Other results in this direction see [31]. In [15], Phillip S. Harrington and Andrew Raich established sufficient conditions for the closed range of $\bar{\partial}$ (and $\bar{\partial}_{b}$ ) on not necessarily pseudoconvex domains (and their boundaries) in Stein manifolds. Also, Phillip S. Harrington and Andrew Raich established sufficient conditions for the closed range of $\bar{\partial}$ (and $\bar{\partial}_{b}$ ) on domains neither boundedness nor pseudoconvexity in $\mathbb{C}^{n}$ (see [16]).

As an application of Theorem 1, we prove that the ranges of $\bar{\partial}_{b}$ and its adjoint $\bar{\partial}_{b}^{*}$ are closed for Lipschitz boundaries of $q$-pseudoconvex domains.
Theorem 2. Let $\Omega \Subset \mathbb{C}^{n}$ be a bounded $q$-pseudoconvex domain with Lipschitz boundary $b \Omega$ and let $1 \leqslant q \leqslant n$. Then, one obtains
(i) $\bar{\partial}_{b}$ and $\bar{\partial}_{b}^{*}$ acting on $L_{p, q}^{2}(b \Omega)$ have closed range for every $0 \leqslant p \leqslant n, 1 \leqslant q \leqslant n-1, n \geqslant 2$.
(ii) The space of harmonic forms on the boundary $b \Omega$ vanishes, i.e.,

$$
\mathscr{H}_{b}^{p, q}(b \Omega)=\{0\}, \quad \text { for } 0 \leqslant p \leqslant n \quad 1 \leqslant q<n-1 .
$$

When the unbounded operator is the $\bar{\partial}_{b}$ operator, the Hilbert space approach has been established by J.J. Kohn in [23] for strongly pseudoconvex domains and by L. Hörmander in [21]
for pseudoconvex domain in a Stein manifold. When the boundary of a pseudoconvex domain is smooth, the Hodge decomposition on $b \Omega$ has been obtained by Mei-Chi Shaw in [36] for $1 \leqslant q<n-1$ and by H. P. Boas and Mei-Chi Shaw in [3] for $q=n-1$ (See also Mei-Chi Shaw in [37] for $C^{1}$ or Lipschitz boundaries).

In the end of the paper, we will prove that the $\bar{\partial}_{b}$-Laplacian, or Kohn Laplacian, $\square_{b}=\bar{\partial}_{b} \bar{\partial}_{b}^{*}+$ $\bar{\partial}_{b}^{*} \bar{\partial}_{b}$ has closed range for $(p, q)$-forms when $0 \leqslant p \leqslant n, 1 \leqslant q \leqslant n-1, n \geqslant 2$. Thus there exists a bounded inverse operator for $\square_{b}$, the $\bar{\partial}_{b}$-Neumann operator $N_{b}$, and we have the decomposition for $\bar{\partial}_{b}$ on $b \Omega$ : $\alpha=\bar{\partial}_{b} \bar{\partial}_{b}^{*} N_{b} \alpha+\bar{\partial}_{b}^{*} \bar{\partial}_{b} N_{b} \alpha$ for any $(p, q)$-forms $\alpha$ with $L^{2}(b \Omega)$ coefficients.

Theorem 3. Let $\Omega \Subset \mathbb{C}^{n}$ be a bounded $q$-pseudoconvex domain with Lipschitz boundary $b \Omega$ and let $1 \leqslant q \leqslant n$. Then, for each $0 \leqslant p \leqslant n, 1 \leqslant q \leqslant n-1, n \geqslant 2$, there exists a bounded linear boundary operator $N_{b}: L_{p, q}^{2}(b \Omega) \longrightarrow L_{p, q}^{2}(b \Omega)$ such that
(i) $\mathscr{R} \operatorname{ang}\left(N_{b}\right) \Subset \operatorname{Dom} \square_{b}$ and $\square_{b} N_{b}=N_{b} \square_{b}=I$ on $\operatorname{Dom} \square_{b}$.
(ii) For $\alpha \in L_{p, q}^{2}(b \Omega)$, we have $\alpha=\bar{\partial}_{b} \bar{\partial}_{b}^{*} N_{b} \alpha \oplus \bar{\partial}_{b}^{*} \bar{\partial}_{b} N_{b} \alpha$.
(iii) $\bar{\partial}_{b} N_{b}=N_{b} \bar{\partial}_{b}$ on $\operatorname{Dom} \bar{\partial}_{b}$, for $1 \leqslant q \leqslant n-1$.
(iv) $\bar{\partial}_{b}^{*} N_{b}=N_{b} \bar{\partial}_{b}^{*}$ on $\operatorname{Dom} \bar{\partial}_{b}^{*}$, for $2 \leqslant q \leqslant n$.
(v) If $\alpha \in L_{p, q}^{2}(b \Omega)$ and $\bar{\partial}_{b} \alpha=0$, then $u=\bar{\partial}_{b}^{*} N_{b} \alpha$ is the unique solution to the equation $\bar{\partial}_{b} u=\alpha$ which is orthogonal to $\operatorname{ker} \bar{\partial}_{b}$.

## 1. Notation and preliminaries

### 1.1. Morrey-Kohn-Hörmander

Let $\Omega$ be a bounded domain in $\mathbb{C}^{n}$ with $C^{2}$ boundary $b \Omega$ and defining function $\rho$ so that $|\partial \rho|=1$ on $b \Omega$. Let $\left(z_{1}, \ldots, z_{n}\right)$ be the complex coordinates for $\mathbb{C}^{n}$. Any $(p, q)$-form $\alpha$ on $\bar{\Omega}$ can be expressed as follows:

$$
\begin{equation*}
\alpha=\sum_{I, J}^{\prime} \alpha_{I, J} \mathrm{~d} z^{I} \wedge \mathrm{~d} \bar{z}^{J} \tag{1}
\end{equation*}
$$

where $I=\left(i_{1}, \ldots, i_{p}\right)$ and $J=\left(j_{1}, \ldots, j_{q}\right)$ are multiindices and $\mathrm{d} z^{I}=\mathrm{d} z_{i_{1}} \wedge \cdots \wedge \mathrm{~d} z_{i_{p}}, \mathrm{~d} \bar{z}^{J}=\mathrm{d} \bar{z}_{j_{1}} \wedge$ $\cdots \wedge \mathrm{d} \bar{z}_{j_{q}}$. The notation $\Sigma^{\prime}$ means the summation over strictly increasing multiindices. Denote by $C^{\infty}\left(\mathbb{C}^{n}\right)$ the space of complex-valued $C^{\infty}$ functions on $\mathbb{C}^{n}$ and $C_{p, q}^{\infty}\left(\mathbb{C}^{n}\right)$ the space of complexvalued differential forms of class $C^{\infty}$ and of type $(p, q)$ on $\mathbb{C}^{n}$, where $0 \leqslant p \leqslant n, 0 \leqslant q \leqslant n$. Let

$$
C_{p, q}^{\infty}(\bar{\Omega})=\left\{\left.u\right|_{\bar{\Omega}} \mid u \in C_{p, q}^{\infty}\left(\mathbb{C}^{n}\right)\right\} .
$$

Denote $\mathscr{D}\left(\mathbb{C}^{n}\right)$, the space of $C^{\infty}$-functions with compact support in $\mathbb{C}^{n}$. A form $u \in C_{p, q}^{\infty}\left(\mathbb{C}^{n}\right)$ is said to be has compact support in $\mathbb{C}^{n}$ if its coefficients belongs to $\mathscr{D}\left(\mathbb{C}^{n}\right)$. The subspace of $C_{p, q}^{\infty}\left(\mathbb{C}^{n}\right)$ which has compact support in $\mathbb{C}^{n}$ is denoted by $\mathscr{D}_{p, q}\left(\mathbb{C}^{n}\right)$. For $u, \alpha \in C_{p, q}^{\infty}\left(\mathbb{C}^{n}\right)$, the local inner product ( $u, \alpha$ ) is denoted by:

$$
(u, \alpha)=\sum_{I, J}^{\prime} u_{I, J} \bar{\alpha}_{I, J}
$$

and $(u, u)$ is defined by

$$
(u, u)=|u|^{2}=\sum_{I, J}^{\prime}\left|u_{I, J}\right|^{2} .
$$

The Cauchy-Riemann operator $\bar{\partial}: C_{p, q-1}^{\infty}(\Omega) \longrightarrow C_{p, q}^{\infty}(\Omega)$ is defined by

$$
\bar{\partial} \alpha=\sum_{I, J}^{\prime} \sum_{k=1}^{n} \frac{\partial \alpha_{I, J}}{\partial \bar{z}^{k}} \mathrm{~d} \bar{z}^{k} \wedge \mathrm{~d} z^{I} \wedge \mathrm{~d} \bar{z}^{J}
$$

Recall that $L^{2}(\Omega)$ is the space of square-integrable functions on $\Omega$ with respect to the Lebesgue measure in $\mathbb{C}^{n}$ and $L_{p, q}^{2}(\Omega)$ is the space of $(p, q)$-forms with coefficients in $L^{2}(\Omega)$. If $u, \alpha \in L_{p, q}^{2}(\Omega)$, the $L^{2}$-inner product $\langle u, \alpha\rangle_{\Omega}$ and norm $\|u\|_{\Omega}$ are defined by

$$
\langle u, \alpha\rangle_{\Omega}=\langle u, \alpha\rangle_{L_{p, q}^{2}(\Omega)}=\int_{\Omega}(u, \alpha) \mathrm{d} V=\int_{\Omega} u \wedge \star \bar{\alpha}
$$

and

$$
\|u\|_{\Omega}^{2}=\|u\|_{L_{p, q}^{2}(\Omega)}^{2}=\langle u, u\rangle_{\Omega},
$$

where $\mathrm{d} V$ is the volume element induced by the Hermitian metric and $\star: C_{p, q}^{\infty}\left(\mathbb{C}^{n}\right) \longrightarrow$ $C_{n-q, n-p}^{\infty}\left(\mathbb{C}^{n}\right)$ is the Hodge star operator such that $\overline{\star u}=\star \bar{u}$ (that is $\star$ is a real operator) and $\star \star u=(-1)^{p+q} u$. For $u \in C_{p, q}^{\infty}(\Omega)$ and $\alpha \in \mathscr{D}_{p, q-1}(\Omega)$, the formal adjoint operator $\vartheta$ of $\bar{\partial}$ : $C_{p, q-1}^{\infty}(\Omega) \longrightarrow C_{p, q}^{\infty}(\Omega)$, with respect to $\langle\cdot \cdot \cdot\rangle_{\Omega}$, is defined by

$$
\langle\bar{\partial} \alpha, u\rangle_{\Omega}=\langle\alpha, \vartheta u\rangle_{\Omega} .
$$

Thus $\vartheta$ can be expressed explicitly by

$$
\begin{equation*}
\vartheta u=(-1)^{p-1} \sum_{\substack{|I|=p \\|K|=q-1}}^{\prime} \sum_{j=1}^{n} \frac{\partial u_{I, j K}}{\partial z^{j}} \mathrm{~d} z^{I} \wedge \mathrm{~d} \bar{z}^{K} . \tag{2}
\end{equation*}
$$

The operator $\vartheta$ defined in (2) satisfies

$$
\begin{equation*}
\vartheta=-\star \partial \star . \tag{3}
\end{equation*}
$$

Let $\bar{\partial}: \operatorname{Dom} \bar{\partial} \subset L_{p, q}^{2}(\Omega) \longrightarrow L_{p, q+1}^{2}(\Omega)$ be the maximal closed extensions of the original $\bar{\partial}$ and $\bar{\partial}^{*}: \operatorname{Dom} \bar{\partial}^{*} \subset L_{p, q}^{2}(\Omega) \longrightarrow L_{p, q-1}^{2}(\Omega)$ be the Hilbert space adjoint of $\bar{\partial}$. Let $\operatorname{ker} \bar{\partial}=\{\alpha \in \operatorname{Dom} \bar{\partial}$ : $\bar{\partial} \alpha=0\}$ and $\mathscr{R}$ ang $\bar{\partial}=\{\bar{\partial} \alpha: \alpha \in \operatorname{Dom} \bar{\partial}\}$, be the kernel and the range of $\bar{\partial}$, respectively. The complex Laplacian $\square$ is defined by $\square=\bar{\partial} \bar{\partial}^{*}+\bar{\partial}^{*} \bar{\partial}: L_{p, q}^{2}(\Omega) \longrightarrow L_{p, q}^{2}(\Omega)$ on Dom $\square=\left\{\alpha \in \operatorname{Dom} \bar{\partial} \cap \operatorname{Dom} \bar{\partial}^{*}:\right.$ $\bar{\partial} \alpha \in \operatorname{Dom} \bar{\partial}^{*}$ and $\left.\bar{\partial}^{*} \alpha \in \operatorname{Dom} \bar{\partial}\right\}$. The space of harmonic forms $\mathscr{H}^{p, q}(\Omega)$ is defined by

$$
\mathscr{H}^{p, q}(\Omega)=\left\{\alpha \in L_{p, q}^{2}(\Omega) \cap \operatorname{Dom} \bar{\partial} \cap \operatorname{Dom} \bar{\partial}^{*}: \bar{\partial} \alpha=\bar{\partial}^{*} \alpha=0\right\} .
$$

Let $\mathbb{H}: L_{p, q}^{2}(\Omega) \longrightarrow \operatorname{ker} \square$ be the orthogonal projection from the space $L_{p, q}^{2}(\Omega)$ onto the space ker $\square$. The $\bar{\partial}$-Neumann operator

$$
N: L_{p, q}^{2}(\Omega) \longrightarrow L_{p, q}^{2}(\Omega)
$$

is defined as the inverse of the restriction of $\square$ to $\left(\mathscr{C}^{p, q}(\Omega)\right)^{\perp}$, i.e.,

$$
N \alpha= \begin{cases}0 & \text { if } \alpha \in \mathscr{H}^{p, q}(\Omega), \\ u & \text { if } \alpha=\square u, \text { and } u \perp \mathscr{H}^{p, q}(\Omega) .\end{cases}
$$

In other words, $N \alpha$ is the unique solution $u$ to the equations $\Vdash u=0, \square u=\alpha-\sharp \alpha$. The Bergman projection operator $P: L_{p, q}^{2}(\Omega) \longrightarrow \operatorname{ker} \bar{\partial}$ is the orthogonal projection of $L_{p, q}^{2}(\Omega)$ onto $\operatorname{ker} \bar{\partial}$. For any $0 \leqslant p \leqslant n$ and $1 \leqslant q \leqslant n, P$ is represented in terms of $N$ by the Kohn's formula

$$
\begin{equation*}
P=I-\bar{\partial}^{*} \bar{\partial} N . \tag{4}
\end{equation*}
$$

Let $a=\left(a_{1}, \ldots, a_{n}\right)$ be a multiindices, that is, $a_{1}, \ldots, a_{n}$ are nonnegative integers. For $x \in \mathbb{R}^{n}$, one defines $x^{a}=x_{1}^{a_{1}} \ldots x_{n}^{a_{n}}$ and $D^{a}$ is the operator

$$
D^{a}=\left(\frac{1}{i} \frac{\partial}{\partial x_{1}}\right)^{a_{1}} \cdots\left(\frac{1}{i} \frac{\partial}{\partial x_{n}}\right)^{a_{n}} .
$$

Denote by $\mathscr{S}$ the Schwartz space of rapidly decreasing smooth functions on $\mathbb{R}^{n}$, that is, $\mathscr{S}$ consists of all functions $u$ which are smooth on $\mathbb{R}^{n}$ with $\sup _{x \in \mathbb{R}^{n}}\left|x^{a} D^{b} u(x)\right|<\infty$ for all multiindices $a, b$. The Fourier transform $\widehat{u}$ of a function $u \in \mathscr{S}$ is defined by

$$
\widehat{u}(\xi)=(2 \pi)^{-n / 2} \int_{\mathbb{R}^{n}} u(x) e^{-i x \cdot \xi} \mathrm{~d} x,
$$

where $x . \xi=\sum_{v=1}^{n} x_{v} \xi_{v}$ and $\mathrm{d} x=\mathrm{d} x_{1} \wedge \cdots \wedge \mathrm{~d} x_{n}$ with $x=\left(x_{1}, \ldots, x_{n}\right)$ and $\xi=\left(\xi_{1}, \ldots, \xi_{n}\right)$. If $u \in \mathscr{S}$, then $\widehat{u} \in \mathscr{S}$. The Sobolev space $W^{m}\left(\mathbb{R}^{n}\right), m \in \mathbb{R}$, is the completion of $\mathscr{S}$ under the Sobolev norm

$$
\|u\|_{W^{m}\left(\mathbb{R}^{n}\right)}^{2}=\int_{\mathbb{R}^{n}}\left(1+|\xi|^{2}\right)^{m}|\widehat{u}|^{2} \mathrm{~d} \xi
$$

Denote by $W^{m}(\Omega), m \geqslant 0$, the space of the restriction of all functions $u \in W^{m}\left(\mathbb{C}^{n}\right)=W^{m}\left(\mathbb{R}^{2 n}\right)$ to $\Omega$ and

$$
\|u\|_{W^{m}(\Omega)}=\inf \left\{\|\alpha\|_{W^{m}\left(\mathbb{C}^{n}\right)}, \alpha \in W^{m}\left(\mathbb{C}^{n}\right),\left.\alpha\right|_{\Omega}=u\right\}
$$

is the $W^{m}(\Omega)$-norm. Let $W_{0}^{m}(\Omega)$ be the completion of $\mathscr{D}(\Omega)$ under the $W^{m}(\Omega)$-norm. If $\Omega$ is a Lipschitz domain, $C^{\infty}(\bar{\Omega})$ is dense in $W^{m}(\Omega)$ with respect to the $W^{m}(\Omega)$-norm. If $0 \leqslant m \leqslant \frac{1}{2}$, $\mathscr{D}(\Omega)$ is dense in $W^{m}(\Omega)$ (cf. [11, Theorem 1.4.2.4]). Thus

$$
W^{m}(\Omega)=W_{0}^{m}(\Omega), \text { for } 0 \leqslant m \leqslant \frac{1}{2}
$$

For $m>0$, one defines $W^{-m}(\Omega)$ to be the dual of $W_{0}^{m}(\Omega)$ and the norm of $W^{-m}(\Omega)$ is defined by

$$
\|u\|_{W^{-m}(\Omega)}=\sup _{0 \neq \alpha \in W_{0}^{m}(\Omega)} \frac{\left|\langle u, \alpha\rangle_{\Omega}\right|}{\|\alpha\|_{W^{m}(\Omega)}}
$$

Denote by $W_{p, q}^{m}(\Omega), m \in \mathbb{R}$, the Hilbert spaces of $(p, q)$-forms with $W^{m}(\Omega)$-coefficients and their norms are denoted by $\|u\|_{W^{m}(\Omega)}$. Noting that, for a bounded domain $\Omega$, the generalized Schwartz inequality, for $u \in W^{m}(\Omega)$ and $\alpha \in W^{-m}(\Omega)$,

$$
\begin{equation*}
\left|\langle u, \alpha\rangle_{\Omega}\right| \leqslant\|u\|_{W^{m}(\Omega)}\|\alpha\|_{W^{-m}(\Omega)} \tag{5}
\end{equation*}
$$

holds when $-\frac{1}{2} \leqslant m \leqslant \frac{1}{2}$.
Lemma 4 ([4]). Let $\Omega \Subset \mathbb{C}^{n}$ be a bounded domain with $C^{2}$ boundary and $\rho$ be a $C^{2}$ defining function of $\Omega$. Let $\varphi \in C^{2}(\bar{\Omega})$ with $\varphi \geqslant 0$. Then, for $\alpha \in C_{p, q}^{\infty}(\bar{\Omega}) \cap \operatorname{Dom}_{\bar{\partial}}{ }^{*}$ with $1 \leqslant q \leqslant n-1$, one obtains

$$
\begin{align*}
&\|\sqrt{\varphi} \bar{\partial} \alpha\|_{\Omega}^{2}+\left\|\sqrt{\varphi} \bar{\partial}^{*} \alpha\right\|_{\Omega}^{2}=\sum_{I, J}^{\prime} \sum_{j, k=1}^{n} \int_{b \Omega} \varphi \frac{\partial^{2} \rho}{\partial z^{j} \partial \bar{z}^{k}} \alpha_{I, j K} \bar{\alpha}_{I, k K} \mathrm{~d} S \\
&+\sum_{I, J}^{\prime} \sum_{k=1}^{n} \int_{\Omega} \varphi\left|\frac{\partial \alpha_{I, J}}{\partial \bar{z}^{k}}\right|^{2} \mathrm{~d} V+2 \operatorname{Re}\left(\sum_{I, K}^{\prime} \sum_{j=1}^{n} \frac{\partial \varphi}{\partial z^{j}} \alpha_{I, j K} \mathrm{~d} \bar{z}_{K}, \bar{\partial}^{*} \alpha\right) \\
&-\sum_{I, K}^{\prime} \sum_{j, k=1}^{n} \int_{\Omega} \frac{\partial^{2} \varphi}{\partial z^{j} \partial \bar{z}^{k}} \alpha_{I, j K} \bar{\alpha}_{I, k K} \mathrm{~d} V . \tag{6}
\end{align*}
$$

The case of $\varphi \equiv 1$ is the classical Kohn-Morrey formula see [21,23].

### 1.2. The $\bar{\partial}_{b}$ complex on Lipschitz domains

In this subsection, we introduce square-integrable ( $p, q$ )-forms on a Lipschitz boundary $b \Omega$ of a bounded domain $\Omega$ in $\mathbb{C}^{n}$ with distance function $\rho$. We equip $b \Omega$ with the induced metric from $\mathbb{C}^{n}$. A boundary $b \Omega$ of a bounded domain $\Omega \Subset \mathbb{C}^{n}$ is called Lipschitz if locally the boundary $b \Omega$ is the graph of a Lipschitz function. Let $\psi: \mathbb{R}^{2 n-1} \longrightarrow \mathbb{R}$ be a function that satisfies the Lipschitz condition

$$
\begin{equation*}
\left|\psi(x)-\psi\left(x^{\prime}\right)\right| \leqslant M\left|x-x^{\prime}\right|, \text { for all } x, x^{\prime} \in \mathbb{R}^{2 n-1} \tag{7}
\end{equation*}
$$

The smallest $M>0$ in which (7) holds is called the bound of the Lipschitz constant. A boundary $b \Omega$ of a bounded domain $\Omega \Subset \mathbb{C}^{n}$ is called Lipschitz if near every boundary point $p \in b \Omega$ there exists a neighborhood $U$ of $p$ such that, after a rotation,

$$
\Omega \cap U=\left\{\left(x, x_{2 n}\right) \in U \mid x_{2 n}>\psi(x)\right\}
$$

for some Lipschitz function $\psi$. By choosing finitely many balls $\left\{U_{i}\right\}$ covering $b \Omega$, the Lipschitz constant for a Lipschitz domain is the smallest $M$ such that the Lipschitz constant is bounded by $M$ in every ball $U_{i}$. A Lipschitz function is almost everywhere differentiable (see [7]).
Definition 5. A bounded domain $\Omega$ with Lipschitz boundary b $\Omega$ in $\mathbb{C}^{n}$ is said to have a global Lipschitz defining function if there exists a Lipschitz function $\rho: \mathbb{C}^{n} \longrightarrow \mathbb{R}$ such that $\rho<0$ in $\Omega$, $\rho>0$ outside $\bar{\Omega}$ and

$$
\begin{equation*}
c_{1}<|\mathrm{d} \rho|<c_{2} \text { a.e. on } b \Omega, \tag{8}
\end{equation*}
$$

where $c_{1}, c_{2}$ are positive constants.
We cover $b \Omega$ by finitely many boundary coordinate patches $U_{i}$ where $i=1, \ldots, k$. Let $r_{i}$ be a local defining function on $U_{i}$ which is locally a Lipschitz graph. Let $\chi_{i} \in C_{0}^{\infty}\left(U_{i}\right)$ be a partition of unity such that $\sum_{i} \chi_{i}=1$ in a neighborhood of $b \Omega$. We define $\rho=\sum_{i} \chi_{i} r_{i}$. Then $\rho$ is a defining function for $\Omega$.

Lemma 6 ([37]). Let $\Omega \Subset \mathbb{C}^{n}$ be a bounded domain with Lipschitz boundary $b \Omega$. Then $\Omega$ has a global Lipschitz defining function $\rho$. Furthermore, the distance function to the boundary is comparable to $|\rho|$ for any global Lipschitz defining function $\rho$ near the boundary.

Let $C^{\infty}(b \Omega)$ be the space of the restriction of all smooth functions in $\mathbb{C}^{n}$ to $b \Omega$. For each $m$ with $1 \leqslant m \leqslant \infty$, one defines $\widetilde{L}_{p, q}^{m}(b \Omega)$ to be the space of $(p, q)$-forms in $\mathbb{C}^{n}$ such that each coefficient of $\alpha$, when restricted to $b \Omega$, is in $L^{m}(b \Omega)$. Write $\alpha$ as in (1), then $\alpha \in \widetilde{L}_{p, q}^{m}(b \Omega)$ if and only if $\left.\alpha_{I, J}\right|_{b \Omega} \in L^{m}(b \Omega)$ for each $I$, $J$. Let $\vee$ be the interior product which is the dual of the wedge product. Since the boundary is Lipschitz, the normal vector is defined almost everywhere and satisfies (8). If we fix $p \in b \Omega$, then for some neighborhood $U$ of $p$ we may locally choose an orthonormal coordinate patch $\left\{\mathrm{d} z_{1}, \ldots, \mathrm{~d} z_{n}\right\}$ defined almost everywhere in $U \cap \bar{\Omega}$ such that $\mathrm{d} \bar{z}_{n}=\bar{\partial} \rho$ (note that $|\bar{\partial} \rho|=\frac{1}{2}$ because we are using the metric where $\left|\mathrm{d} z^{j}\right|=1$, which is half the size induced by the usual Euclidean metric on $\left.\mathbb{R}^{n}\right)$. We define $L_{p, q}^{m}(b \Omega) \subset \widetilde{L}_{p, q}^{m}(b \Omega)$ as the space of all $\widetilde{L}_{p, q}^{m}$ such that d $\bar{z}_{n} \vee \alpha=0$ almost everywhere on $b \Omega$.

Locally, if $\alpha \in \widetilde{L}_{p, q}^{m}(b \Omega \cap U)$, one can express

$$
\alpha=\sum_{\substack{I, J \\ n \notin J}}^{\prime} \alpha_{I, J} \mathrm{~d} z^{I} \wedge \mathrm{~d} \bar{z}^{J}+\sum_{\substack{I, J \\ n \in J}}^{\prime} \alpha_{I, J} \mathrm{~d} z^{I} \wedge \mathrm{~d} \bar{z}^{J}
$$

where $\alpha_{I, J}$ 's are $L^{m}(b \Omega \cap U)$ functions. Let $\tau$ denote the projection map

$$
\tau: \widetilde{L}_{p, q}^{m}(b \Omega) \longrightarrow L_{p, q}^{m}(b \Omega)
$$

defined by

$$
\begin{equation*}
\tau \alpha=\sum_{\substack{I, J \\ n \notin J}}^{\prime} \alpha_{I, J} \mathrm{~d} z^{I} \wedge \mathrm{~d} \bar{z}^{J} \tag{9}
\end{equation*}
$$

Since changing basis will result in multiplication by $L^{\infty}(b \Omega \cap U)$ functions, the projection $\tau$ is well-defined since it is independent of the choice of $\left\{\mathrm{d} \bar{z}_{1}, \ldots, \mathrm{~d} \bar{z}_{n-1}\right\}$.

Let $\Lambda_{p, q}(b \Omega)$ denote the restriction of $C_{p, q}^{\infty}\left(\mathbb{C}^{n}\right)$ to $b \Omega$. Define $\mathscr{B}_{p, q}(b \Omega)$ to be the subspace of $L_{p, q}^{\infty}(b \Omega)$ such that $\alpha \in \mathscr{B}_{p, q}(b \Omega)$ if and only if there exists $\widetilde{\alpha} \in \Lambda_{p, q}(b \Omega)$ such that $f=\tau \widetilde{\alpha}$. In other words, we have $\tau\left(\Lambda_{p, q}\left(\mathbb{C}^{n}\right)\right)=\mathscr{B}_{p, q}(b \Omega)$. Obviously, $\mathscr{B}_{p, q}(b \Omega) \Subset L_{p, q}^{\infty}(b \Omega) \Subset L_{p, q}^{2}(b \Omega)$. Denote by $W_{p, q}^{m}(b \Omega), 0 \leqslant m \leqslant 1$, the space of forms that are the completion of $\mathscr{B}_{p, q}(b \Omega)$-forms with $W^{m}(b \Omega)$-norms. This is well defined also for Lipschitz domains since on $b \Omega, W^{1}(b \Omega)$ is well defined and the boundary value of any function in $W^{1}(\Omega)$ to the boundary belongs to $W^{\frac{1}{2}}(b \Omega)$ (see [22]).

Lemma 7 ([30, Lemma 1.4]). Let $\Omega$ be a bounded domain with Lipschitz boundary b $\Omega$ in $\mathbb{R}^{n}$. Then $C^{\infty}(b \Omega)$ is dense in $L^{2}(b \Omega)$ and $\Lambda_{p, q}(b \Omega)$ is dense in $\widetilde{L}_{p, q}^{2}(b \Omega)$ for every $0 \leqslant p \leqslant n, 0 \leqslant q \leqslant n$. Also $\mathscr{B}_{p, q}(b \Omega)$ is a dense subset in $L_{p, q}^{2}(b \Omega)$ for every $0 \leqslant p \leqslant n, 0 \leqslant q \leqslant n-1$.

The Bochner-Martinelli-Koppelman transform on $(p, q)$-forms is defined as follows. Let

$$
\begin{gathered}
(\zeta-z)=\left(\zeta_{1}-z_{1}, \ldots, \zeta_{n}-z_{n}\right) \\
\mathrm{d} \zeta=\left(\mathrm{d} \zeta_{1}, \ldots, \mathrm{~d} \zeta_{n}\right)
\end{gathered}
$$

Define

$$
\begin{aligned}
\langle\bar{\zeta}-\bar{z}, \mathrm{~d} \zeta\rangle & =\sum_{j=1}^{n}\left(\bar{\zeta}_{j}-\bar{z}_{j}\right) \mathrm{d} \zeta_{j}, \\
\langle\mathrm{~d} \bar{\zeta}-\mathrm{d} \bar{z}, \mathrm{~d} \zeta\rangle & =\sum_{j=1}^{n}\left(\mathrm{~d} \bar{\zeta}_{j}-\mathrm{d} \bar{z}_{j}\right) \mathrm{d} \zeta_{j} .
\end{aligned}
$$

The Bochner-Martinelli-Koppelman kernel $K(\zeta, z)$ is defined by

$$
\begin{aligned}
K(\zeta, z) & =\frac{1}{(2 \pi i)^{n}} \frac{\langle\bar{\zeta}-\bar{z}, \mathrm{~d} \zeta\rangle}{|\zeta-z|^{2}} \wedge\left(\frac{\langle\mathrm{~d} \bar{\zeta}-\mathrm{d} \bar{z}, \mathrm{~d} \zeta\rangle}{|\zeta-z|^{2}}\right)^{n-1} \\
& =\sum_{q=0}^{n-1} K_{q}(\zeta, z)
\end{aligned}
$$

where $K_{q}(\zeta, z)$ denote the component of $K(\zeta, z)$ that is a $(p, q)$-form in $z$ and an $(n-p, n-q-1)$ form in $\zeta$.

When $n=1$,

$$
K(\zeta, z)=\frac{1}{2 \pi i} \frac{\mathrm{~d} \zeta}{\zeta-z}
$$

is the Cauchy kernel. As in the Cauchy integral case, for any $f \in L_{p, q}^{2}(b \Omega)$ the Cauchy principal value integral $K_{b} f$ is defined as follows:

$$
K_{b} f(z)=\lim _{\epsilon \longrightarrow 0^{+}} \int_{\substack{b \Omega \\|\zeta-z|>\epsilon}} K_{q}(\zeta, z) \wedge f(\zeta)
$$

whenever the limit exists. Denote by $v_{z}$ the outward unit normal to $b \Omega$ at $z$. Since $b \Omega$ is Lipschitz, $v_{z}$ exists almost everywhere on $b \Omega$. Then, for $z \in b \Omega$, one defines

$$
\begin{aligned}
& K_{b}^{-} f(z)=\lim _{\epsilon \longrightarrow 0^{+}} \int_{b \Omega} K_{q}\left(\cdot, z-\epsilon v_{z}\right) \wedge f \\
& K_{b}^{+} f(z)=\lim _{\epsilon \longrightarrow 0^{+}} \int_{b \Omega} K_{q}\left(\cdot, z+\epsilon v_{z}\right) \wedge f
\end{aligned}
$$

Proposition 8 ([37]). Let $\Omega$ be a bounded domain in $\mathbb{C}^{n}$ with Lipschitz boundary. For $\alpha \in$ $\mathscr{B}_{p, q}(b \Omega), 0 \leqslant q \leqslant n-1$, the following formula holds for almost every $z \in b \Omega$ :

$$
\begin{equation*}
\alpha(z)=\tau \lim _{\epsilon \longrightarrow 0^{+}}\left(\int_{b \Omega} K_{q}\left(\cdot, z-\epsilon v_{z}\right) \wedge \alpha-\int_{b \Omega} K_{q}\left(\cdot, z+\epsilon v_{z}\right) \wedge \alpha\right) \tag{10}
\end{equation*}
$$

The $\bar{\partial}_{b}$-operator is defined distributionally as follows: for any $u \in L_{p, q}^{2}(b \Omega)$ and $\alpha \in L_{p, q+1}^{2}(b \Omega)$ we say that $u$ is in $\operatorname{Dom} \bar{\partial}_{b}$ and $\bar{\partial}_{b} u=\alpha$ if and only if:

$$
\int_{b \Omega} u \wedge \bar{\partial} \phi \mathrm{~d} S=(-1)^{p+q} \int_{b \Omega} \alpha \wedge \phi \mathrm{~d} S, \text { for every } \phi \in C_{n-p, n-q-1}^{\infty}\left(\mathbb{C}^{n}\right)
$$

Since $\bar{\partial}^{2}=0$, it follows that $\bar{\partial}_{b}^{2}=0$. Thus $\bar{\partial}_{b}$ is a complex and we have the following:

$$
0 \longrightarrow L_{p, 0}^{2}(b \Omega) \xrightarrow{\bar{\partial}_{b}} L_{p, 1}^{2}(b \Omega) \xrightarrow{\bar{\partial}_{b}} L_{p, 2}^{2}(b \Omega) \xrightarrow{\bar{\partial}_{b}} \ldots \xrightarrow{\bar{\partial}_{b}} L_{p, n-1}^{2}(b \Omega) \longrightarrow 0
$$

Proposition 9 ([35]). Let $\Omega$ be a bounded domain in $\mathbb{C}^{n}$ with Lipschitz boundary b $\Omega$. The $\bar{\partial}_{b}$ operator is a closed, densely defined, linear operator from $L_{p, q-1}^{2}(b \Omega)$ to $L_{p, q}^{2}(b \Omega)$, where $0 \leqslant p \leqslant n$, $1 \leqslant q \leqslant n-1$.

We need to define $\bar{\partial}_{b}^{*}$, the $L^{2}$ adjoint of $\bar{\partial}_{b}$. Again, we first define its domain:

Definition 10. $\operatorname{Dom} \bar{\partial}_{b}^{*}$ is the subset of $L_{p, q}^{2}(b \Omega)$ composed of all forms $\alpha$ for which there exists $a$ constant $c>0$ such that

$$
\left|\left\langle\alpha, \bar{\partial}_{b} u\right\rangle_{b \Omega}\right| \leqslant C\|u\|_{b \Omega}
$$

for all $u \in \operatorname{Dom} \bar{\partial}_{b}$.
For all $\alpha \in \operatorname{Dom} \bar{\partial}_{b}^{*}$, let $\bar{\partial}_{b}^{*} \alpha$ be the unique form in $L_{p, q}^{2}(b \Omega)$ satisfying

$$
\left\langle\bar{\partial}_{b}^{*} \alpha, u\right\rangle_{b \Omega}=\left\langle\alpha, \bar{\partial}_{b} u\right\rangle_{b \Omega}
$$

for all $u \in \operatorname{Dom} \bar{\partial}_{b}$.
Definition 11. Let $\square_{b}=\bar{\partial}_{b} \bar{\partial}_{b}^{*}+\bar{\partial}_{b}^{*} \bar{\partial}_{b}: \operatorname{Dom} \square_{b} \longrightarrow L_{p, q}^{2}(b \Omega)$ the $\bar{\partial}_{b}$-Laplacian operator defined on $\operatorname{Dom} \square_{b}=\left\{\alpha \in L_{p, q}^{2}(b \Omega): \alpha \in \operatorname{Dom} \bar{\partial}_{b} \cap \operatorname{Dom} \bar{\partial}_{b}^{*}: \bar{\partial}_{b} \alpha \in \operatorname{Dom} \bar{\partial}_{b}^{*}\right.$ and $\left.\bar{\partial}_{b}^{*} \alpha \in \operatorname{Dom} \bar{\partial}_{b}\right\}$.
Proposition 12 ([35, Proposition 1.3]). The $\bar{\partial}_{b}$-Laplacian operator is a closed, densely defined self-adjoint operator.

Let $\mathscr{C}_{b}^{p, q}(b \Omega)$ denote the space of harmonic forms on the boundary $b \Omega$, i.e.,

$$
\mathscr{H}_{b}^{p, q}(b \Omega)=\left\{\alpha \in \operatorname{Dom} \square_{b}: \bar{\partial}_{b} \alpha=\bar{\partial}_{b}^{*} \alpha=0\right\}
$$

The space $\mathscr{H}_{b}^{p, q}(b \Omega)$ is a closed subspace of $\operatorname{Dom} \square_{b}$ since $\square_{b}$ is a closed operator. One defines the boundary operator or the $\bar{\partial}_{b}$-Neumann operator

$$
N_{b}: L_{p, q}^{2}(b \Omega) \longrightarrow L_{p, q}^{2}(b \Omega)
$$

as the inverse of the restriction of $\square_{b}$ to $\left(\mathscr{H}_{b}^{p, q}(b \Omega)\right)^{\perp}$, i.e.,

$$
N_{b} \alpha= \begin{cases}0 & \text { if } \alpha \in \mathscr{H}_{b}^{p, q}(b \Omega), \\ u & \text { if } \alpha=\square_{b} u, \text { and } u \perp \mathscr{H}_{b}^{p, q}(b \Omega) .\end{cases}
$$

In other words, $N_{b} \alpha$ is the unique solution $u$ to the equations $\alpha=\square_{b} u$ with $\alpha \perp \mathscr{H}_{b}^{p, q}(b \Omega)$ and we extend $N_{b}$ by linearity.

## 1.3. $q$-pseudoconvex domains

In this subsection, we recall the following definition of $q$-subharmonic functions which has been introduced by H. Ahn and N. Q. Dieu in [1] (also see Lop-Hing Ho [20]). For a real valued $C^{2}$ function $u$ defined on $U \subseteq \mathbb{C}^{n}$, Lop-Hing Ho [20] first defined $q$-subharmonicity of $u$ on $U$ and using this $q$-subharmonic function, he introduce the notion of weak $q$-convexity for domains with smooth boundaries. As Theorem 1.4 in [20], Ahn and Dieu [1] investigated a natural extension of these notions to the class of upper semicontinuous functions and $q$-pseudoconvex domains with non-smooth boundaries.

Definition 13 ([1]). Let $u$ be an upper semicontinuous function on $\Omega$. Then we say that $u$ is $q$ subharmonic on $\Omega$ iffor every $q$-complex dimension space $H$ and for every compact set $D \Subset H \cap \Omega$, the following holds: if $h$ is a continuous harmonic function on $D$ and $h \geqslant u$ on the boundary of $D$, then $h \geqslant u$ on $D$.

Definition 14 ([1]). The function $u$ is called strictly $q$-subharmonic iffor every $U \Subset \Omega$ there exists a constant $C_{U}>0$ such that $u-C_{U}|z|^{2}$ is $q$-subharmonic.

The following results gives some basic properties of $q$-subharmonic functions, follows the same lines as for plurisubharmonic functions, that will be used later (see [1,20]).

Proposition 15. Let $\Omega$ be a domain in $\mathbb{C}^{n}$ and let $1 \leqslant q \leqslant n$. Then we have
(i) If $\left\{u_{v}\right\}_{v=1}^{\infty}$ is a decreasing sequence of $q$-subharmonic functions then $u=\lim _{v \rightarrow+\infty} u_{v}$ is a $q$-subharmonic function.
(ii) Let $\left.\chi(z) \in C_{0}^{\infty} \mathbb{C}^{n}\right)$ be a function such that $\chi \geqslant 0, \int_{\mathbb{C}^{n}} \chi(z) \mathrm{d} V=1, \chi(z)$ depends only on $|z|$ and vanishes when $|z|>1$. Set $\chi_{\varepsilon_{v}}(z)=\varepsilon_{v}^{-2 n} \chi\left(z / \varepsilon_{v}\right)$ for $\varepsilon_{v} \downarrow 0$. If $u$ is a $q$-subharmonic function, then

$$
u_{\varepsilon_{v}}(z)=u * \chi_{\varepsilon_{v}}(z)=\int_{\Omega} u(\zeta) \chi_{\varepsilon_{v}}(z-\zeta) \mathrm{d} V(\zeta)
$$

is smooth $q$-subharmonic on $\Omega_{\varepsilon_{v}}=\left\{z \in \Omega: d(z, b \Omega)>\varepsilon_{v}\right\}$. Moreover, $u_{\varepsilon_{v}} \downarrow u$ as $\varepsilon_{v} \downarrow 0$.
(iii) If $u \in C^{2}(\Omega)$ such that $\frac{\partial^{2} u}{\partial z^{j} \partial \bar{z}^{k}}(z)=0$ for all $j \neq k$ and $z \in \Omega$. Then $u$ is $q$-subharmonic if and only if $\sum_{j, k \in J} \frac{\partial^{2} u}{\partial z^{j} \partial \bar{z}^{k}}(z) \geqslant 0$, for all $|J|=q$ and for all $z \in \Omega$.
Proof. The proof of this proposition follows from properties of subharmonic functions. The proof of (ii) is exact as Proposition 1.2 in [1]. Similarly, it is easy to see that (i) and (iii) hold because these properties are true for subharmonic functions.
Proposition 16. Let $\Omega$ be a bounded domain in $\mathbb{C}^{n}$ and let $q$ be an integer with $1 \leqslant q \leqslant n$. Let $u \in C^{2}(\Omega)$. Then the $q$-subharmonicity of $u$ is equivalent to

$$
\sum_{|K|=q-1}^{\prime} \sum_{j, k=1}^{n} \frac{\partial^{2} u}{\partial z^{j} \partial \bar{z}^{k}} \alpha_{j K} \bar{\alpha}_{k K} \geqslant 0
$$

for all $(0, q)$-forms $\alpha=\sum_{|J|=q} \alpha_{J} \mathrm{~d} \bar{z}^{J}$.
Proof. By Theorem 1.4 and Lemma 1.2 in [20], it is easy to see that this fact is true if $u \in C^{2}(\Omega)$. In the case $u$ is arbitrary we note that the assertion is true for $u_{\varepsilon}$. Let $\varepsilon \searrow 0$ we obtain the assertion for $u$ and the proof follows.

The following examples of a $q$-subharmonic function which is not plurisubharmonic.
Example 17. One of the most typical examples of $q$-subharmonic function which is not plurisubharmonic is

$$
u(z)=-\sum_{v=1}^{q-1}\left|z_{v}\right|^{2}+(q-1) \sum_{v=q}^{n}\left|z_{v}\right|^{2} .
$$

Indeed, for $q=2$, we have

$$
\left(\frac{\partial^{2} u}{\partial z^{j} \partial \bar{z}^{k}}\right)=\left[\begin{array}{ccccc}
-1 & 0 & \cdots & 0 & 0 \\
0 & 1 & \cdots & 0 & 0 \\
\cdots & \cdots & \cdots & \cdots & \cdots \\
0 & 0 & \cdots & 0 & 1
\end{array}\right]
$$

Then, following Proposition 15 (iii), $u$ is 2-subharmonic. More precisely

$$
\sum_{l=1}^{n} \sum_{j, k=1}^{n} \frac{\partial^{2} u}{\partial z^{j} \partial \bar{z}^{k}} \alpha_{j l} \bar{\alpha}_{k l}=\sum_{l=1}^{n}\left(-\left|\alpha_{1 l}\right|^{2}+\sum_{j=2}^{n}\left|\alpha_{j l}\right|^{2}\right)=\sum_{l=2}^{n} \sum_{j=2}^{n}\left|\alpha_{j l}\right|^{2} \geqslant 0
$$

for all (0,2)-forms $\alpha=\sum_{|J|=2}^{\prime} \alpha_{J} \mathrm{~d} \bar{z}^{J}$ (because $\alpha_{l l}=0, l=1,2, \ldots, n$ ). Thus, $u$ is 2 -subharmonic.
Example 18. Let $d>1$ and $1<q \leqslant d$. Consider the function

$$
u(z)=|z|^{2}-q\left|z_{1}\right|^{2}=\sum_{j=1}^{d}\left|z_{j}\right|^{2}-q\left|z_{1}\right|^{2}, \quad z \in \mathbb{C}^{d}
$$

It is easy to see that $\sum_{j=1}^{q} \frac{\partial^{2} u}{\partial z^{j} \partial \bar{z}^{j}}(z)=0$ and by Proposition 15 it follows that $u$ is $q$-subharmonic. However, $u$ is not plurisubharmonic. Indeed, let $L=\left\{\left(z_{1}, 0, \ldots, 0\right)\right\} \subset \mathbb{C}^{d}$ be a complex line. Then $\left.u\right|_{L}=(1-q)\left|z_{1}\right|^{2}$ is not subharmonic, and the desired conclusion follows.

Example 19 ([12]). Let $q$ be an integer with $1 \leqslant q \leqslant n$ and let $\lambda_{k}, k=1,2, \ldots, q$, are complex numbers such that $\sum_{k=1}^{q}\left|\lambda_{k}\right|^{2}>0$ and $\lambda_{k} \neq 0$, for $k=q+1, \ldots, n$. Then the function

$$
u(z)=2\left\{\operatorname{Re}\left(\sum_{k=1}^{q} \lambda_{k} z_{k}\right)\right\}^{2}+\sum_{k=q+1}^{n}\left|\lambda_{k} z_{k}\right|^{2}
$$

is 1 -subharmonic and strictly $q$-subharmonic on $\mathbb{C}^{n}$. Moreover, if $q>1$ then $u$ is not strictly ( $q-1$ )-subharmonic on any open set of $\mathbb{C}^{n}$.

According to Lop-Hing Ho (see [20, Definition 2.1]), a smoothly bounded domain $\Omega$ is called weakly $q$-convex if $\Omega$ has a defining function $\rho$ such that for every $z \in b \Omega$ one obtains

$$
\sum_{|K|}^{\prime} \sum_{j, k} \frac{\partial^{2} \rho}{\partial z^{j} \partial \bar{z}^{k}} \alpha_{j K} \bar{\alpha}_{k K} \geqslant 0,
$$

for every $(0, q)$-forms $\alpha=\sum_{|J|=q} \alpha_{J} \mathrm{~d} \bar{z}^{J}$ such that

$$
\sum_{j=1}^{n} \frac{\partial \rho}{\partial z^{j}} \alpha_{j K}=0 \text { for all }|K|=q-1 .
$$

Definition 20. A (Lipschitz) domain $\Omega \Subset \mathbb{C}^{n}$ is said to be $q$-pseudoconvex if there exists $q$ subharmonic exhaustion (Lipschitz) function on $\Omega$. Moreover, a $C^{2}$ smooth bounded domain $\Omega$ is called strictly $q$-convex if it admits a $C^{2}$ smooth defining function which is strictly $q$-subharmonic on a neighbourhood of $\bar{\Omega}$.

## Remark 21.

(i) Note that $\Omega$ is pseudoconvex if and only if it is 1 -pseudoconvex, since 1 -subharmonic function is just plurisubharmonic.
(ii) By Theorem 2.4 in Lop-Hing Ho [20], every weakly $q$-convex domain (see [20, Definition 2.1]) is $q$-pseudoconvex.
(iii) If $\Omega \Subset \mathbb{C}^{n}$ is a $q$-pseudoconvex domain, $1 \leqslant q \leqslant n$, and if $b \Omega$ is of class $C^{2}$, then $\Omega$ is weakly $q$-convex (see Lop-Hing Ho [20]).
(iv) Let $\Omega \Subset \mathbb{C}^{n}$ be a bounded domain satisfy the $Z(q)$ condition. Thus $\Omega$ is strictly $q$ pseudoconvex.
(v) Every $n$-dimensional connected non compact complex manifold has a strongly subharmonic exhaustion function with respect to any hermitian metric $\omega$. Thus, every open set in $\mathbb{C}^{n}$ is $n$-pseudoconvex (see Greene and $\mathrm{Wu}[10]$ ).

## 2. Existence of the $\bar{\partial}$-Neumann operator

This section deals with the existence of the $\bar{\partial}$-Neumann operator $N$ on $q$-pseudoconvex domains in $\mathbb{C}^{n}$.

Lemma 22 ( $\left[\mathbf{1}, \mathbf{1 2 ] ) . ~ L e t ~} \Omega \Subset \mathbb{C}^{n}\right.$ be a $q$-pseudoconvex domain, $1 \leqslant q \leqslant n$. Then $\Omega$ has a $C^{\infty}$-smooth strictly $q$-subharmonic exhaustion function. More precisely, there are strictly $q$ pseudoconvex domains, $\Omega_{v}$ 's, $v=1,2, \ldots$, with smooth boundary satisfying

$$
\Omega=\cup_{v=1}^{\infty} \Omega_{v}, \Omega_{v} \Subset \Omega_{v+1} \Subset \Omega \text { for all } v .
$$

Proof. Let $u$ be $q$-subharmonic exhaustion function for $\Omega$. By induction one can choose a sequence $\left\{a_{v}\right\}_{v \geqslant 1} \uparrow \infty$ such that the open sets $U_{v}:=\left\{u<a_{v}\right\}$ satisfy $U_{v} \Subset U_{v+1} \Subset \Omega$. Next, for each $v$ we choose $\varepsilon_{v}>0$ so small that

$$
d\left(U_{v-1}, b U_{v}\right)>\varepsilon_{v}, \quad d\left(U_{v}, b U_{v+1}\right)>\varepsilon_{v}, \quad U_{v+1} \Subset \Omega_{\varepsilon_{v}},
$$

where $\Omega_{\varepsilon_{v}}=\left\{z: d(z, b \Omega)>\varepsilon_{v}\right\}$. Put

$$
u_{v}(z)=\int_{\Omega} u(\zeta) \chi_{\varepsilon_{v}}(z-\zeta) \mathrm{d} V(\zeta), \forall z \in \Omega_{\varepsilon_{v}}
$$

Since $u<u_{v}$ on $\Omega_{\varepsilon_{v}}$ we deduce that

$$
\Omega_{v}=\left\{z \in \Omega_{\varepsilon_{v}}: u_{v}(z)>a_{v}\right\} \subset U_{v} \Subset U_{v+1} \Subset \Omega_{\varepsilon_{v}} \subset \Omega
$$

We claim that $U_{v-1} \subset \Omega_{v}$. Indeed, let $z \in U_{v-1}$. Then we have $\mathbb{B}\left(z, \varepsilon_{v}\right) \subset U_{v}$ and, hence,

$$
u_{v}(z)=\int_{\mathbb{B}\left(z, \varepsilon_{v}\right)} u(\zeta) \chi_{\varepsilon_{v}}(z-\zeta) \mathrm{d} V(\zeta)<a_{v}
$$

This proves the claim, and therefore $\Omega=\cup_{v=1}^{\infty} \Omega_{v}$, Finally, it is easy to see that $\Omega_{v}$ is a $q$ pseudoconvex domain with the smooth strictly $q$-subharmonic exhaustion function $\varphi_{v}(z):=$ $\frac{1}{a_{v}-u_{v}(z)}+|z|^{2}$.

Following Lemmas 6 and 22, as Lemma 2.1 in [29], we prove the following lemma:
Lemma 23. Let $\Omega \Subset \mathbb{C}^{n}$ be a bounded $q$-pseudoconvex domain with Lipschitz boundary $b \Omega$. There exists an exhaustion $\left\{\Omega_{v}\right\}$ of $\Omega$ such that
(i) there exists a Lipschitz function $\rho: \mathbb{C}^{n} \longrightarrow \mathbb{R}$ such that $\rho<0$ in $\Omega, \rho>0$ outside $\bar{\Omega}$ and satisfies (8).
(ii) $\left\{\Omega_{v}\right\}$ is an increasing sequence of relatively compact subsets of $\Omega$ and $\Omega=\cup_{v} \Omega_{v}$.
(iii) Each $\Omega_{v}, v=1,2, \ldots$, is strictly $q$-pseudoconvex domains, i.e., each $\Omega_{v}$ has a $C^{\infty}$ strictly $q$-subharmonic defining function $\rho_{v}$ on a neighbourhood of $\bar{\Omega}$, such that

$$
\sum_{I,|K|}^{\prime} \sum_{j, k} \frac{\partial^{2} \rho_{v}}{\partial z^{j} \partial \bar{z}^{k}} \alpha_{I, j K} \bar{\alpha}_{I, k K} \geqslant c_{0}|\alpha|^{2}
$$

$\alpha \in C_{p, q}^{\infty}\left(\bar{\Omega}_{v}\right) \cap \operatorname{Dom} \bar{\partial}_{v}^{*}$ with $q \geqslant 1$ and $c_{0}$ is independent of $v$.
(iv) There exist positive constants $c_{1}, c_{2}$ such that $c_{1} \leqslant\left|\nabla \eta_{v}\right| \leqslant c_{2}$ on $b \Omega_{v}$, where $c_{1}, c_{2}$ are independent of $v$.

Proof. Using Lemma 6, there exists a global Lipschitz defining function $\rho$ for $\Omega$ satisfying (8) and $\rho$ is obtained as above by a partition of unity of defining functions $r_{i}$ which is a Lipschitz graph. Choose a function $\chi(z) \in C_{0}^{\infty}\left(\mathbb{C}^{n}\right)$ as in Proposition 15. Let $\delta_{v}$ be a sequence of small numbers with $\delta_{v} \searrow 0$. For each $\delta_{v}$, one defines

$$
\Omega_{\delta_{v}}=\left\{z \in \Omega \mid \rho(z)<-\delta_{v}\right\}
$$

Then $\Omega_{\delta_{v}}$ is a sequence of relatively compact open subsets of $\Omega$ with union equal to $\Omega$. For each $\delta_{v}$, one defines, for $0<\epsilon<\delta_{v}$ and $z \in \Omega_{\delta_{v}}$,

$$
\rho_{\epsilon}(z)=\rho * \chi_{\epsilon}(z)=\int \rho(z-\epsilon \zeta) \chi(\zeta) \mathrm{d} V(\zeta)
$$

Then, $\rho_{\epsilon} \in C^{\infty}\left(\Omega_{\delta_{v}}\right)$ and $\rho_{\epsilon} \searrow \rho$ on $\Omega_{\delta_{v}}$. Since $\rho$ is $q$-subharmonic, it follows that, from Proposition 15 , that $\rho_{\epsilon}$ is $q$-subharmonic.

Each $\rho_{\epsilon_{v}}$ is well defined if $0<\epsilon_{v}<\delta_{v+1}$ for $z \in \Omega_{\delta_{v+1}}$. Let $c_{2}=\sup _{\Omega}|\nabla \rho|$, then for $\epsilon_{v}$ sufficiently small, one obtains

$$
\rho(z)<\rho_{\epsilon_{v}}(z)<\rho(z)+c_{2} \epsilon_{v} \text { on } \Omega_{\delta_{v+1}}
$$

For each $v$ we choose

$$
\epsilon_{v}=\frac{1}{2 c_{2}}\left(\delta_{v-1}-\delta_{v}\right) \text { and } t_{v} \in\left(\delta_{v+1}-\delta_{v}\right)
$$

Define

$$
\Omega_{v}=\left\{z \in \mathbb{C}^{n} \mid \rho_{\epsilon_{v}}<-t_{v}\right\}
$$

Since $\rho(z)<\rho_{\epsilon_{v}}(z)<-t_{v}<-\delta_{v+1}$, we have that $\Omega_{\delta_{v+1}} \supset \Omega_{v}$. Also, if $z \in \Omega_{\delta_{v-1}}$, then $\rho_{\epsilon_{v}}(z)<$ $\rho(z)+c_{2} \epsilon_{v}<-\delta_{v}<-t_{v}$. Thus we have $\Omega_{\delta_{v+1}} \supset \Omega_{v} \supset \Omega_{\delta_{v-1}}$ and (ii) is satisfied.

Each $\Omega_{v}$ is defined by $\eta_{v}=\rho_{\epsilon_{v}}+t_{v}$ which is strictly $q$-subharmonic in $\Omega_{v}$ and (iii) is satisfied. That the subdomain $\Omega_{v}$ has smooth boundary will follow from (iv).

To prove (iv), it is easy to see that $\left|\nabla \eta_{v}\right| \leqslant c_{2}$ in $b \Omega_{v}$. To show that $\left|\nabla \eta_{v}\right|$ is uniformly bounded from below, we note $b \Omega$ satisfies the uniform interior cone property. Then there exists a conic neighborhood $\Gamma$ with vertex $0 \in \mathbb{C}^{n}$ such that for any unit vector $\xi \in \Gamma+\{p\},-\langle\nabla \rho, \xi\rangle_{p}>c_{0}$ a.e. in $U \cap b \Omega$, where $c_{0}$ is a positive constant independent of $p$ if $U$ is sufficiently small. There exist a finite covering $\left\{V_{\mu}\right\}_{1 \leqslant \mu \leqslant K}$ of $b \Omega$, a finite set of unit vectors $\left\{\xi_{\mu}\right\}_{1 \leqslant \mu \leqslant K}$ and $c_{1}>0$ such that the inner product $\left\langle\nabla \rho, \xi_{\mu}\right\rangle \geqslant c_{1}>0$ a.e. for $z \in V_{\mu}, 1 \leqslant \mu \leqslant K$. Since this is preserved by convolution, (iv) is proved. Thus the proof follows.

Theorem 24. Let $\Omega \Subset \mathbb{C}^{n}$ be a $q$-pseudoconvex domain, $1 \leqslant q \leqslant n$. Then, for any $1 \leqslant q \leqslant n$, there exists a bounded linear operator $N: L_{p, q}^{2}(\Omega) \longrightarrow L_{p, q}^{2}(\Omega)$ which satisfies the following properties:
(i) $\mathscr{R} a n g N \Subset \operatorname{Dom} \square, N \square=I$ on $\operatorname{Dom} \square$.
(ii) For any $\alpha \in L_{p, q}^{2}(\Omega)$, we have $\alpha=\bar{\partial} \bar{\partial}^{*} N \alpha \oplus \bar{\partial}^{*} \bar{\partial} N \alpha$.
(iii) $\bar{\partial} N=N \bar{\partial}$ on $\operatorname{Dom} \bar{\partial}, 1 \leqslant q \leqslant n-1, n \geqslant 2$.
(iv) $\bar{\partial}^{*} N=N \bar{\partial}^{*}$ on $\operatorname{Dom} \bar{\partial}^{*}, 2 \leqslant q \leqslant n$.
(v) If $\delta$ is the diameter of $\Omega$, we have the following estimates:

$$
\begin{gather*}
\|N \alpha\|_{\Omega} \leqslant \frac{e \delta^{2}}{q}\|\alpha\|_{\Omega} \\
\|\bar{\partial} N \alpha\|_{\Omega} \leqslant \sqrt{\frac{e \delta^{2}}{q}}\|\alpha\|_{\Omega}  \tag{11}\\
\left\|\bar{\partial}^{*} N \alpha\right\|_{\Omega} \leqslant \sqrt{\frac{e \delta^{2}}{q}}\|\alpha\|_{\Omega} .
\end{gather*}
$$

Proof. We first prove the theorem for $\Omega$ with $C^{2}$ boundary. Let $\rho$ be the defining function of $\Omega$. Following Definitions 14 and 20, there exists $\alpha \in C_{p, q}^{\infty}(\bar{\Omega}) \cap \operatorname{Dom} \overline{\bar{\partial}}^{*}$ with $q \geqslant 1$, such that

$$
\sum_{I, K}^{\prime} \sum_{j, k=1}^{n} \int_{b \Omega} \frac{\partial^{2} \rho}{\partial z^{j} \partial \bar{z}^{k}} \alpha_{I, j K} \bar{\alpha}_{I, k K} \mathrm{~d} S \geqslant 0 .
$$

If we replace $\varphi$ by $1-e^{\psi}$, where $\psi$ is an arbitrary twice continuously differentiable non-positive function, and after applying the Cauchy-Schwarz inequality (5), for $m=0$, to the term in (6) involving first derivatives of $\varphi$, we find

$$
\|\sqrt{\varphi} \bar{\partial} \alpha\|_{\Omega}^{2}+\left\|\sqrt{\varphi} \bar{\partial}^{*} \alpha\right\|_{\Omega}^{2} \geqslant \sum_{I, K}^{\prime} \sum_{j, k=1}^{n} \int_{\Omega} e^{\psi} \frac{\partial^{2} \psi}{\partial z^{j} \partial \bar{z}^{k}} \alpha_{I, j K} \bar{\alpha}_{I, k K} \mathrm{~d} V-\left\|e^{\psi / 2} \bar{\partial}^{*} \alpha\right\|_{\Omega} .
$$

Since $\varphi+e^{\psi}=1$ and $\varphi \leqslant 1$, it follows that

$$
\|\bar{\partial} \alpha\|_{\Omega}^{2}+\left\|\bar{\partial}^{*} \alpha\right\|_{\Omega}^{2} \geqslant \sum_{I, K}^{\prime} \sum_{j, k=1}^{n} \int_{\Omega} e^{\psi} \frac{\partial^{2} \psi}{\partial z^{j} \partial \bar{z}^{k}} \alpha_{I, j K} \bar{\alpha}_{I, k K} \mathrm{~d} V,
$$

for all $\alpha \in C_{p, q}^{\infty}(\bar{\Omega}) \cap \operatorname{Dom} \bar{\partial}^{*}$ and for $q \geqslant 1$. If $z_{0} \in \Omega$, and $\psi(z)=-1+\left|z-z_{0}\right|^{2} / \delta^{2}$, where $\delta=\sup _{z, z^{\prime} \in \Omega}\left|z-z^{\prime}\right|$ is the diameter of the bounded domain $\Omega$, then the preceding inequality implies

$$
\|\alpha\|_{\Omega}^{2} \leqslant\left(\frac{e \delta^{2}}{q}\right)\left(\|\bar{\partial} \alpha\|_{\Omega}^{2}+\left\|\bar{\partial}^{*} \alpha\right\|_{\Omega}^{2}\right) .
$$

This estimate was derived when $\alpha$ is continuous and differential on $\bar{\Omega}$. Thus, by density it holds for all square-integrable forms $\alpha \in \operatorname{Dom} \bar{\partial} \cap \operatorname{Dom} \overline{\bar{\partial}}^{*}$. Thus, for $q \geqslant 1$, one obtains

$$
\begin{equation*}
\|\alpha\|_{\Omega} \leqslant\left(\frac{e \delta^{2}}{q}\right)\|\square \alpha\|_{\Omega} . \tag{12}
\end{equation*}
$$

For a general $q$-pseudoconvex domain, from Lemma 22 , one can exhaust $\Omega$ by a sequence of strictly $q$-pseudoconvex domains with $C^{\infty}$ boundary $b \Omega$. We write $\Omega=\cup_{v} \Omega_{v}$, where each $\Omega_{v}$ is a bounded strictly $q$-pseudoconvex domains with $C^{\infty}$ boundary and $\Omega_{v} \Subset \Omega_{v+1} \Subset \Omega$ for each $v$. Let $\delta_{v}$ be the diameter of $\Omega_{v}$ and let $\square_{v}$ be the complex Laplacian on each $\Omega_{v}$. Thus (12) holds on each $\Omega_{v}$. That is, there exists a $\alpha_{v} \in C_{p, q}^{\infty}\left(\bar{\Omega}_{v}\right) \cap$ Dom $\bar{\partial}_{v}^{*}$ with $q \geqslant 1$, such that

$$
\left\|\alpha_{v}\right\|_{\Omega_{v}} \leqslant\left(\frac{e \delta_{v}^{2}}{q}\right)\left\|\square_{v} \alpha_{v}\right\|_{\Omega_{v}} \leqslant\left(\frac{e \delta^{2}}{q}\right)\|\square \alpha\|_{\Omega} .
$$

We can choose a subsequence of $\alpha_{v}$, still denoted by $\alpha_{v}$, such that $\alpha_{v} \longrightarrow \alpha$ weakly in $L_{p, q}^{2}(\Omega)$. Furthermore, $\alpha$ satisfies the estimate

$$
\begin{equation*}
\|\alpha\|_{\Omega} \leqslant \liminf \left(\frac{e \delta_{v}^{2}}{q}\right)\left\|\square_{v} \alpha_{v}\right\|_{\Omega_{v}} \leqslant\left(\frac{e \delta^{2}}{q}\right)\|\square \alpha\|_{\Omega} . \tag{13}
\end{equation*}
$$

Since $\square$ is a linear closed densely defined operator, then, from Theorem 1.1.1 in [21], $\mathscr{R}$ ang $(\square)$ is closed. Thus, from (1.1.1) in [21] and the fact that $\square$ is self adjoint, one obtains

$$
L_{p, q}^{2}(\Omega)=\bar{\partial} \bar{\partial}^{*} \operatorname{Dom} \square \oplus \bar{\partial}^{*} \bar{\partial} \operatorname{Dom} \square .
$$

Since $\square: \operatorname{Dom} \square \longrightarrow \mathscr{R}$ ang $(\square)=L_{p, q}^{2}(\Omega)$ is one to one on $\operatorname{Dom} \square$ from (13), there exists a unique bounded inverse operator $N: \mathscr{R}$ ang $(\square) \longrightarrow$ Dom $\square$ such that $N \square \alpha=\alpha$ on Dom $\square$. Also, from the definition of $N$, one obtains $\square N=I$ on $L_{p, q}^{2}(\Omega)$. Thus (i) and (ii) are satisfied. To show that $\bar{\partial} N=N \bar{\partial}$ on $\operatorname{Dom} \bar{\partial}$, by using (ii), we have $\bar{\partial} \alpha=\frac{\bar{\partial}^{\prime}}{}{ }^{*} \bar{\partial} N \alpha$, for $\alpha \in \operatorname{Dom} \bar{\partial}$. Thus

$$
N \bar{\partial} \alpha=N \bar{\partial}_{\partial} \bar{\partial}^{*} \bar{\partial} N \alpha=N\left(\bar{\partial}^{*} \bar{\partial}+\bar{\partial} \bar{\partial}^{*}\right) \bar{\partial} N \alpha=\bar{\partial} N \alpha .
$$

And, by similar argument, $\bar{\partial}^{*} N=N \bar{\partial}^{*}$ on $\operatorname{Dom} \bar{\partial}^{*}$ for $2 \leqslant q \leqslant n$. By using (iii) and the condition on $\alpha, \overline{\bar{\partial}} \alpha=0$, we have $\bar{\partial} N \alpha=N \bar{\partial} \alpha=0$. Then, by using (ii), one obtains $\alpha=\bar{\partial} \bar{\partial}^{*} N \alpha$. Thus the form $u=\bar{\partial}^{*} N \alpha$ satisfies the equation $\bar{\partial} u=\alpha$. Since $\mathscr{R}$ ang $N \subset \operatorname{Dom} \square$, then by applying (13) to $N \alpha$ instead of $\alpha$, (11) follows. Thus the proof follows.

## 3. Sobolev estimates for the $\bar{\alpha}$-Neumann problem

Let $\Omega$ be a domain with $C^{1}$-boundary $b \Omega$, and let $\rho$ be a $C^{1}$-defining function of $\Omega$. Assume that $\mathscr{E}_{p, q}(b \Omega)$ is the space of the restriction to $b \Omega$ of all $(p, q)$-forms with $C^{1}(\bar{\Omega})$-coefficients which are pointwise orthogonal to the ideal generated by $\bar{\partial} \rho$ and $\digamma_{p, q}(b \Omega)$ is the space of the restriction to $b \Omega$ of all $(p, q)$-forms that are multiples of $\bar{\partial} \rho$. Denote by $\rceil_{1}$, the projection $\rceil_{1}: C_{p, q}^{1}(\bar{\Omega}) \longrightarrow$ $\mathscr{E}_{p, q}(b \Omega)$ and $\rceil_{2}$, the projection $\rceil_{2}: C_{p, q}^{1}(\bar{\Omega}) \longrightarrow \digamma_{p, q}(b \Omega)$. In particular, $\left.\left.\rceil_{1} \oplus\right\rceil_{2}=\right\rceil$, where $\rceil$ is simply the restriction map from $C_{p, q}^{1}(\bar{\Omega})$ to the boundary. If $\Omega$ has only Lipschitz boundary $b \Omega$, the operators $T_{1}$ and $T_{2}$ are also defined almost everywhere on $b \Omega$.
Lemma 25 ([30]). Let $\Omega$ be a domain with $C^{1}$-boundary $b \Omega$, and let $\rho$ be a $C^{1}$-defining function of $\Omega$. For any $f \in L_{p, q}^{2}(\Omega)$, the restriction maps $\left.\urcorner(\bar{\partial} N),\right\urcorner P$ and $\urcorner\left(\bar{\partial}^{*} N\right)$ can be extended as bounded operators from $L_{p, q}^{2}(\Omega)$ to $W_{p, q+1}^{-\frac{1}{2}}(b \Omega), W_{p, q}^{-\frac{1}{2}}(b \Omega)$ and $W_{p, q-1}^{-\frac{1}{2}}(b \Omega)$, respectively.

Denote by $\mathscr{R}_{1}$ and $\mathscr{R}_{2}$, the restriction maps of $\bar{\partial}^{*} N$ and $\bar{\partial} N$ to $b \Omega$, respectively. For any $\alpha \in L_{p, q}^{2}(\Omega), \mathscr{R}_{1} f=7\left(\bar{\partial}^{*} N \alpha\right)$ and $\mathscr{R}_{2} f=\tau(\bar{\partial} N \alpha)$. From Lemma 25, one obtains

$$
\mathscr{R}_{1}: L_{p, q}^{2}(\Omega) \longrightarrow W_{p, q-1}^{-\frac{1}{2}}(b \Omega)
$$

and

$$
\mathscr{R}_{2}: L_{p, q}^{2}(\Omega) \longrightarrow W_{p, q+1}^{-\frac{1}{2}}(b \Omega)
$$

Let $T_{1}: W_{p, q-1}^{\frac{1}{2}}(b \Omega) \longrightarrow L_{p, q}^{2}(\Omega)$ be the dual of $\mathscr{R}_{1}$ and be defined as follows: For a fixed $\alpha \in$ $W_{p, q-1}^{\frac{1}{2}}(b \Omega)$ and for any $u \in L_{p, q}^{2}(\Omega)$, we have, using Lemma 25 , that

$$
\left|\int_{b \Omega}\left(\bar{\partial}^{*} N u, \alpha\right) \mathrm{d} S\right| \leqslant C\left\|\mathscr{R}_{1} u\right\|_{W^{-\frac{1}{2}}(b \Omega)}\|\alpha\|_{W^{\frac{1}{2}}(b \Omega)} \leqslant C\|u\|_{\Omega}
$$

where $C$ depends on $\alpha$. Thus there exists an element $g=T_{1} \alpha \in L_{p, q}^{2}(\Omega)$ such that

$$
\int_{b \Omega}\left(\bar{\partial}^{*} N u, \alpha\right) \mathrm{d} S=\left\langle u, T_{1} \alpha\right\rangle_{\Omega}, \text { for any } u \in L_{p, q}^{2}(\Omega)
$$

Let $T_{2}: W_{p, q+1}^{\frac{1}{2}}(b \Omega) \longrightarrow L_{p, q}^{2}(\Omega)$ be the dual of $\mathscr{R}_{2}$, such that for any $\alpha \in W_{p, q+1}^{\frac{1}{2}}(b \Omega)$

$$
\int_{b \Omega}(\bar{\partial} N u, \alpha) \mathrm{d} S=\left\langle u, T_{2} \alpha\right\rangle_{\Omega}, \text { for any } u \in L_{p, q}^{2}(\Omega)
$$

Lemma 26 ([30, Lemma 4.3]). Let $\Omega$ be a domain with $C^{\infty}$-boundary such that $\Omega$ has a $C^{\infty}{ }_{-}$ plurisubharmonic defining function $\rho$. For any $\alpha \in W_{p, q-1}^{\frac{1}{2}}(b \Omega), 0 \leqslant p \leqslant n, 1 \leqslant q \leqslant n$, we have the following estimate:

$$
\begin{equation*}
\int_{\Omega}(-\rho)\left|T_{1} \alpha\right|^{2} \mathrm{~d} V \leqslant C \int_{b \Omega}|\alpha|^{2} \mathrm{~d} S \tag{14}
\end{equation*}
$$

Also for any $\alpha \in W_{p, q+1}^{\frac{1}{2}}(b \Omega), 0 \leqslant p \leqslant n, 1 \leqslant q \leqslant n$, one obtains

$$
\begin{equation*}
\int_{\Omega}(-\rho)\left|T_{2} \alpha\right|^{2} \mathrm{~d} V \leqslant C \int_{b \Omega}|\alpha|^{2} \mathrm{~d} S \tag{15}
\end{equation*}
$$

where $C$ is a constant depending only on the Lipschitz constant of $\Omega$ and the diameter of $\Omega$.
Theorem 27. Let $\Omega \Subset \mathbb{C}^{n}$ be a bounded $q$-pseudoconvex domain with a Lipschitz boundary $b \Omega$ and let $1 \leqslant q \leqslant n$. Then, for $0 \leqslant p \leqslant n, 1 \leqslant q \leqslant n$ and $-\frac{1}{2} \leqslant m \leqslant \frac{1}{2}$, the operators $N, \bar{\partial} N, \bar{\partial}^{*} N$ and the Bergman projection $P$ are bounded on $W_{p, q}^{m}(\Omega)$ and satisfies the following estimates: there exists $C>0$ such that for any $\alpha \in W_{p, q}^{m}(\Omega)$,

$$
\begin{gather*}
\|N \alpha\|_{W_{p, q}^{m}(\Omega)} \leqslant C\|\alpha\|_{W_{p, q}^{m}(\Omega)} \\
\|\bar{\partial} N \alpha\|_{W_{p, q+1}^{m}(\Omega)}+\left\|\bar{\partial}^{*} N \alpha\right\|_{W_{p, q-1}^{m}(\Omega)} \leqslant C\|\alpha\|_{W_{p, q}^{m}(\Omega)}  \tag{16}\\
\|P \alpha\|_{W_{p, q}^{m}(\Omega)}^{m} \leqslant C\|\alpha\|_{W_{p, q}^{m}(\Omega)}
\end{gather*}
$$

where $C$ depends only on the Lipschitz constant and the diameter of $\Omega$, but is independent of $\alpha$.
Proof. To prove (16) for $m=\frac{1}{2}$, we approximate $\Omega$ as Lemma 23 by a sequence of subdomains $\Omega_{v}=\left\{\rho<-\epsilon_{v}\right\}$ such that each $\Omega_{v}$ is strictly $q$-pseudoconvex domains with $C^{\infty}$ smooth boundary, i.e, each $\Omega_{v}$ has a $C^{\infty}$ strictly $q$-subharmonic defining function $\rho_{v}$ such that (ii) and (iii) in Lemma 23. Thus, we can apply Lemma 26 on each $\Omega_{v}$. We use $T_{1}^{v}, T_{2}^{v}, \mathscr{R}_{1}^{v}, \mathscr{R}_{2}^{v}$, and $N^{v}$, to denote the corresponding operators on each $\Omega_{v}$. By applying (14) and (15) to $T_{1}^{v}, T_{2}^{v}$, for any $\alpha \in W_{p, q-1}^{\frac{1}{2}}\left(b \Omega_{v}\right)$, where $0 \leqslant p \leqslant n, 1 \leqslant q \leqslant n$, we have

$$
\begin{equation*}
\int_{\Omega_{v}}\left(-\rho_{v}\right)\left|T_{1}^{v} \alpha\right|^{2} \mathrm{~d} V_{v} \leqslant C \int_{b \Omega_{v}}|\alpha|^{2} \mathrm{~d} S_{v} \tag{17}
\end{equation*}
$$

where $C$ can be chosen independent of $v$. Also for any $\alpha \in W_{p, q+1}^{\frac{1}{2}}\left(b \Omega_{v}\right)$, where $0 \leqslant p \leqslant n$, $1 \leqslant q \leqslant n$, one obtains

$$
\begin{equation*}
\int_{\Omega_{v}}\left(-\rho_{v}\right)\left|T_{2}^{v} \alpha\right|^{2} \mathrm{~d} V_{v} \leqslant C \int_{b \Omega_{v}}|\alpha|^{2} \mathrm{~d} S_{v} \tag{18}
\end{equation*}
$$

Using Lemma 26, $T_{1}^{v}$ is bounded from $L_{p, q-1}^{2}\left(b \Omega_{v}\right)$ to $W_{p, q-1}^{-\frac{1}{2}}\left(\Omega_{v}\right)$ and $T_{2}^{v}$ is bounded from $L_{p, q+1}^{2}\left(b \Omega_{v}\right)$ to $W_{p, q}^{-\frac{1}{2}}\left(\Omega_{v}\right)$. Also, the bounds depend only on the Lipschitz constant and the diameter of the domain. Thus, from (17) and (18) we have from duality, for any $\alpha \in W_{p, q}^{\frac{1}{2}}\left(\Omega_{v}\right)$,

$$
\begin{equation*}
\left\|\mathscr{R}_{1}^{v} \alpha\right\|_{L_{p, q-1}^{2}\left(b \Omega_{v}\right)} \leqslant C\|\alpha\|_{W_{p, q}\left(\Omega_{v}\right)} \tag{19}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|\mathscr{R}_{1}^{v} \alpha\right\|_{L_{p, q+1}^{2}\left(b \Omega_{v}\right)} \leqslant C\|\alpha\|_{W_{p, q}\left(\Omega_{v}\right)} \tag{20}
\end{equation*}
$$

where $C$ is a constant independent of $v$. For any $\alpha \in W_{p, q}^{\frac{1}{2}}\left(\Omega_{v}\right)$ and by using the trace theorem for elliptic equations (cf. [22,38]), from (19) and (20), one obtains

$$
\begin{equation*}
\left\|\bar{\partial}^{*} N^{v} \alpha\right\|_{W_{p, q-1}^{\frac{1}{2}\left(\Omega_{v}\right)}} \leqslant C\|\alpha\|_{W_{p, q}^{\frac{1}{2}}\left(\Omega_{v}\right)} \tag{21}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|\bar{\partial} N^{v} \alpha\right\|_{W_{p, q+1}^{\frac{1}{2}}\left(\Omega_{v}\right)} \leqslant C\|\alpha\|_{W_{p, q}^{\frac{1}{2}}\left(\Omega_{v}\right)} \tag{22}
\end{equation*}
$$

where $C$ is a constant independent of $v$. Passing to the limit, one obtains from (21) and (22) that

$$
\begin{equation*}
\left\|\bar{\partial}^{*} N \alpha\right\|_{W_{p, q-1}^{\frac{1}{2}}(\Omega)} \leqslant C\|\alpha\|_{W_{p, q}^{\frac{1}{2}}(\Omega)}, \tag{23}
\end{equation*}
$$

and

$$
\begin{equation*}
\|\bar{\partial} N \alpha\|_{W_{p, q+1}^{\frac{1}{2}}(\Omega)} \leqslant C\|\alpha\|_{W_{p, q}^{\frac{1}{2}(\Omega)}} \tag{24}
\end{equation*}
$$

Using Theorem 24 (iii) and (iv), one can write

$$
N=\left(\bar{\partial}^{*} \bar{\partial}^{*}+\bar{\partial}^{*} \bar{\partial}\right) N^{2}=(\bar{\partial} N)\left(\bar{\partial}^{*} N\right)+\left(\bar{\partial}^{*} N\right)(\bar{\partial} N) .
$$

It follows from (23) and (24) that

$$
\begin{equation*}
\|N \alpha\|_{W_{p, q}^{\frac{1}{2}}(\Omega)} \leqslant C\|\alpha\|_{W_{p, q}^{\frac{1}{2}}(\Omega)} \tag{25}
\end{equation*}
$$

By virtue of Kohn's formula (4), we have

$$
\begin{equation*}
\|P \alpha\|_{W_{p, q}^{\frac{1}{2}}(\Omega)} \leqslant C\|f\|_{W_{p, q}}^{\frac{1}{2}(\Omega)}, \tag{26}
\end{equation*}
$$

This proves the continuity of $P$ in $W_{p, q}^{\frac{1}{2}}(\Omega)$. Thus (16) follows for $m=\frac{1}{2}$. Using duality, the estimates (23) to (26) also hold for $m=-\frac{1}{2}$. The other cases follow from interpolation.

## 4. The $L^{2} \bar{\partial}$-Cauchy problem

Using the duality relations pertaining to the $\bar{\partial}$-Neumann problem, one can solve the Cauchy problem for $\bar{\partial}$ on $q$-pseudoconvex domains. This method was first used in [26] for smooth forms on strongly pseudo-convex domains. As Theorem 9.1.2 in [6] (cf. [34, Proposition 2.7]), the following result is proved:

Theorem 28. Let $\Omega \Subset \mathbb{C}^{n}$ be a bounded $q$-pseudoconvex domain with a Lipschitz boundary $b \Omega$ and let $1 \leqslant q \leqslant n$. Then, for every $\alpha \in L_{p, q}^{2}\left(\mathbb{C}^{n}\right)$ which is supported in $\bar{\Omega}$ such that

$$
\begin{equation*}
\bar{\partial} \alpha=0, \text { for } 1 \leqslant q \leqslant n-1, \tag{27}
\end{equation*}
$$

one can find $u \in L_{p, q-1}^{2}\left(\mathbb{C}^{n}\right)$ such that $\bar{\partial} u=\alpha$ in the distribution sense in $\mathbb{C}^{n}$ with $u$ supported in $\bar{\Omega}$ and

$$
\|u\|_{\mathbb{C}^{n}} \leqslant C\|\alpha\|_{\mathbb{C}^{n}}
$$

where $C$ depends only on the Lipschitz constant and the diameter of $\Omega$, but is independent of $\alpha$.

Proof. Let $\alpha \in L_{p, q}^{2}\left(\mathbb{C}^{n}\right)$ which is supported in $\bar{\Omega}$, then $\alpha \in L_{p, q}^{2}(\Omega)$. From Theorem 24, the $\bar{\partial}$ Neumann operator $N_{n-p, n-q}$ exists for $n-q \geqslant 1$. Since $N_{n-p, n-q}=\square_{n-p, n-q}^{-1}$ on $\mathscr{R}$ ang $\square_{n-p, n-q}$ and $\mathscr{R}$ ang $N_{n-p, n-q} \subset \operatorname{Dom} \square_{n-p, n-q}$, then $N_{n-p, n-q} \star \bar{\alpha} \in \operatorname{Dom} \square_{n-p, n-q} \subset L_{n-p, n-q}^{2}(\Omega)$, for $q \leqslant n-1$. Thus, one can define $u \in L_{p, q-1}^{2}(\Omega)$ by

$$
\begin{equation*}
u=-\star \overline{\bar{\partial} N_{n-p, n-q} \star \bar{\alpha}} \tag{28}
\end{equation*}
$$

Extending $u$ to $\mathbb{C}^{n}$ by defining $u=0$ in $\mathbb{C}^{n} \backslash \bar{\Omega}$. We want to prove that the extended form $u$ satisfies the equation $\bar{\partial} u=\alpha$ in the distribution sense in $\mathbb{C}^{n}$. To do that we need first clear that

$$
\bar{\partial}^{*}(\star \bar{\alpha})=0 \text { on } \Omega .
$$

For $\eta \in \operatorname{Dom} \bar{\partial} \subset L_{n-p, n-q-1}^{2}(\Omega)$, one obtains

$$
\langle\bar{\partial} \eta, \star \bar{\alpha}\rangle_{\Omega}=\int_{\Omega} \bar{\partial} \eta \wedge \star \star \alpha=(-1)^{(p+q)(p+q-1)} \int_{\Omega} \alpha \wedge \bar{\partial} \eta=(-1)^{p+q}\langle\alpha, \star \partial \bar{\eta}\rangle_{\Omega}
$$

Since $\vartheta=\bar{\partial}^{*}$ on $\mathscr{D}_{p, q}(\Omega)$, when $\vartheta$ acts in the distribution sense and $\mathscr{D}_{p, q}(\Omega)$ is dense in $\operatorname{Dom} \bar{\partial} \cap$ Dom $\bar{\partial}^{*}$ in the graph norm (cf. [21]), then from (3) one obtains

$$
\langle\bar{\partial} \eta, \star \bar{\alpha}\rangle_{\Omega}=\left\langle\alpha, \bar{\partial}^{*} \star \bar{\eta}\right\rangle_{\mathbb{C}^{n}}=\langle\bar{\partial} \alpha, \star \bar{\eta}\rangle_{\mathbb{C}^{n}}=0
$$

because $\alpha$ is supported in $\bar{\Omega}$. It follows that $\bar{\partial}^{*}(\star \bar{\alpha})=0$ on $\Omega$. Using Theorem 24 (iv), one obtains

$$
\bar{\partial}^{*} N_{n-p, n-q}(\star \bar{\alpha})=N_{n-p, n-q-1} \bar{\partial}^{*}(\star \bar{\alpha})=0
$$

Thus, from (3) and (28), one obtains

$$
\begin{align*}
\bar{\partial} u & =-\overline{\partial \star \bar{\partial} N_{n-p, n-q} \star \bar{\alpha}} \\
& =(-1)^{p+q+1} \overline{\star \star \partial \star \bar{\partial} N_{n-p, n-q} \star \bar{\alpha}} \\
& =(-1)^{p+q} \overline{\star \bar{\partial}^{*} \bar{\partial} N_{n-p, n-q} \star \bar{\alpha}}  \tag{29}\\
& =(-1)^{p+q} \overline{\star\left(\bar{\partial}^{*} \bar{\partial}+\bar{\partial}^{*} \bar{\partial}^{*}\right) N_{n-p, n-q} \star \bar{\alpha}} \\
& =(-1)^{p+q} \overline{\star \star \bar{\alpha}} \\
& =\alpha
\end{align*}
$$

in the distribution sense in $\Omega$. Since $u=0$ in $\mathbb{C}^{n} \backslash \Omega$, then for $g \in \operatorname{Dom} \bar{\partial}^{*} \subset L_{p, q}^{2}\left(\mathbb{C}^{n}\right)$, and from (3) and (29), one obtains

$$
\begin{aligned}
\left\langle u, \bar{\partial}^{*} g\right\rangle_{\mathbb{C}^{n}} & =\left\langle u, \bar{\partial}^{*} g\right\rangle_{\Omega} \\
& =\left\langle\star \overline{\bar{\partial}}^{*} g, \star \bar{u}\right\rangle_{\Omega} \\
& =-\langle\star \star \overline{ } \quad \\
& =(-1)^{p+q}\langle\star \bar{g}, \star \bar{u}\rangle_{\Omega} \\
& =\langle\bar{\partial} u, g\rangle_{\Omega} \\
& =\langle\alpha, g\rangle_{\Omega} \\
& =\langle\alpha, g\rangle_{\mathbb{C}^{n}} .
\end{aligned}
$$

Thus $\bar{\partial} u=\alpha$ in the distribution sense in $\mathbb{C}^{n}$.
As Proposition 9.1.3 in [6], the following result is proved:
Proposition 29. Let $\Omega \Subset \mathbb{C}^{n}$ be a bounded $q$-pseudoconvex domain with a Lipschitz boundary $b \Omega$ and let $1 \leqslant q \leqslant n$. Then, for $-\frac{1}{2} \leqslant m<\frac{1}{2}$ and for any $\alpha \in L_{p, n}^{2}\left(\mathbb{C}^{n}\right)$ such that $\alpha$ is supported in $\bar{\Omega}$ and

$$
\begin{equation*}
\int_{\Omega} \alpha \wedge g=0 \text { for any } g \in L_{n-p, 0}^{2}(\Omega) \cap \operatorname{ker} \bar{\partial} \tag{30}
\end{equation*}
$$

one can find $u \in L_{p, n-1}^{2}\left(\mathbb{C}^{n}\right)$ such that $\bar{\partial} u=\alpha$ in the distribution sense in $\mathbb{C}^{n}$ with $u$ is supported in $\bar{\Omega}$ and

$$
\|u\|_{\mathbb{C}^{n}} \leqslant C\|\alpha\|_{\mathbb{C}^{n}},
$$

where $C$ depends only on the Lipschitz constant and the diameter of $\Omega$, but is independent of $\alpha$.
Proof. To prove this result we need to prove that condition (30) is equivalent to $\bar{\partial} \alpha=0$. Fist, we see that (30) implies (27). To see that, if we take $g=\bar{\partial} \star \beta$ for some $\beta \in C_{p, n+1}^{\infty}\left(\mathbb{C}^{n}\right)$ in (30). It is clear that $g \in \operatorname{ker} \bar{\partial}$. By (30) and the fact $\bar{\partial} \star \beta \in \operatorname{ker} \bar{\partial}$, we see that

$$
\left\langle\alpha, \bar{\partial}^{*} \beta\right\rangle_{\mathbb{C}^{n}}=\int_{\Omega} \alpha \wedge \star\left(\bar{\partial}^{*} \beta\right)=(-1)^{p+n+1} \int_{\Omega} \alpha \wedge \bar{\partial}(\star \beta)=0
$$

for any $\beta \in C_{p, n+1}^{\infty}\left(\mathbb{C}^{n}\right)$, where we used the equality $\star(\star \alpha)=(-1)^{p+n} \alpha$ for $u=\bar{\partial} \star \beta \in C_{p, n+1}^{\infty}\left(\mathbb{C}^{n}\right)$. This implies that $\bar{\partial} \alpha=0$ in the distribution sense in $\mathbb{C}^{n}$.

If $\Omega$ has Lipschitz boundary, $C_{n-p, 1}^{\infty}(\bar{\Omega})$ is dense in dom $\bar{\partial}$ in the graph norm. This follows essentially from Friedrichs lemma (see [21, Proposition 1.2.3]). From the definition of $\bar{\partial}^{*}$, one obtains that $\star \alpha \in \operatorname{Dom} \bar{\partial}^{*}$ and $\bar{\partial}^{*}(\star \alpha)=0$. For any $g \in L_{n-p, n-q}^{2}(\Omega) \cap$ ker $\bar{\partial}$, using Theorem 28 (since $1 \leqslant q<n$ ), there exist $u \in L_{n-p, n-q-1}^{2}(\Omega)$ such that $\bar{\partial} u=g$ in $\Omega$. This implies that

$$
\overline{\int_{\Omega} g \wedge \alpha}=\langle\star \alpha, g\rangle=\langle\star \alpha, \bar{\partial} u\rangle=\left\langle\bar{\partial}^{*} \star \alpha, u\right\rangle=0
$$

for any $g \in L_{n-p, n-q}^{2}(\Omega) \cap \operatorname{ker} \bar{\partial}$. Thus $\alpha$ satisfies (30). Thus the proof follows.
As Proposition 3.4 in [37], we prove the following result:
Lemma 30. Let $\Omega \Subset \mathbb{C}^{n}$ be a bounded $q$-pseudoconvex domain with Lipschitz boundary $b \Omega$ and let $1 \leqslant q \leqslant n$. The set of $\bar{\alpha}$-closed forms in $L_{p, q}^{2}\left(\mathbb{C}^{n}\right)$ with support in $\bar{\Omega}$ is dense in the set of $\overline{\bar{\alpha}}$-closed forms in $W_{p, q}^{m}\left(\mathbb{C}^{n}\right)$ with support in $\bar{\Omega}$.
Proof. By this we mean that if $f \in W_{p, q}^{m}\left(\mathbb{C}^{n}\right)$ with $f$ is supported in $\bar{\Omega}$ and $\bar{\partial} f=0$ on $\Omega$, one can construct a sequence $f_{v} \in L_{p, q}^{2}\left(\mathbb{C}^{n}\right)$ with support in $\bar{\Omega}$ such that $f_{v} \longrightarrow f$ in $W_{p, q}^{m}\left(\mathbb{C}^{n}\right)$ and $\bar{\partial} f_{v} \longrightarrow \bar{\partial} f$ in $W_{p, q+1}^{m}\left(\mathbb{C}^{n}\right)$. This is possible by the well-known method of Friedrichs (see [9] or [21] or [6, Appendix D]) as follows: first assume that the domain $\Omega$ is star-shaped and $0 \in \Omega$. Since $\Omega$ is Lipschitz, locally it is star-shaped. This can be done from the usual regularization by convolution. We approximate $f$ first by dilation componentwise. Let $\Omega^{\epsilon}=\{(1+\epsilon) z \mid z \in \Omega\}$ and

$$
f^{\epsilon}=f\left(\frac{z}{1+\epsilon}\right)
$$

where the dilation is performed for each component of $f$. Then $\Omega \Subset \Omega^{\epsilon}$ and $f^{\epsilon} \in W_{p, q}^{m}\left(\Omega^{\epsilon}\right)$. Also $\bar{\partial} f^{\epsilon} \longrightarrow \bar{\partial} f$ in $W_{p, q+1}^{m}\left(\mathbb{C}^{n}\right)$. Choose a function $\chi(z) \in C_{0}^{\infty}\left(\mathbb{C}^{n}\right)$ as in Proposition 15. Extend $f \in W_{p, q}^{m}\left(\mathbb{C}^{n}\right) \cap \operatorname{ker} \bar{\partial}$ to be 0 outside $\Omega$. By regularizing $f^{\epsilon}$ componentwise as before, one can find a family of $f_{(\epsilon)} \in W_{p, q+1}^{m}\left(\mathbb{C}^{n}\right)$ defined by

$$
f_{(\epsilon)}(z)=f\left(\frac{z}{1+\epsilon}\right) * \chi_{\delta_{\epsilon}}=\int_{\mathbb{C}^{n}} f(w) \chi_{\delta_{\epsilon}}\left(\frac{z-w}{1+\epsilon}\right) \mathrm{d} V_{w},
$$

where $\delta_{\epsilon} \backslash 0$ as $\epsilon \backslash 0$ and $\delta_{\epsilon}$ is chosen sufficiently small. The convolution is performed on each component of $f$. In the first integral defining $f_{j}$, we can differentiate under the integral sign to show that $f_{j}$ is $C^{\infty}\left(\mathbb{C}^{n}\right)$. From Young's inequality for convolution, we have

$$
\left\|f_{(\epsilon)}\right\| \leqslant\|f\| .
$$

Since $f_{(\epsilon)} \longrightarrow f$ uniformly when $f \in C_{0}^{\infty}\left(\mathbb{C}^{n}\right)$, a dense subset of $W^{m}\left(\mathbb{C}^{n}\right)$, we have that

$$
f_{(\epsilon)} \longrightarrow f \text { in } W^{m}\left(\mathbb{C}^{n}\right) \text { for every } f \in W^{m}\left(\mathbb{C}^{n}\right) .
$$

Obviously, this implies that $f_{(\epsilon)} \longrightarrow f$ in $W_{p, q+1}^{m}\left(\mathbb{C}^{n}\right)$.
Let $\delta_{v}$ be a sequence of small numbers with $\delta_{v} \backslash 0$. For each $\delta_{v}$, we define

$$
\Omega_{\delta_{v}}=\left\{z \in \Omega \mid \rho(z)<-\delta_{v}\right\} .
$$

Then $\Omega_{\delta_{v}}$ is a sequence of relatively compact open subsets of $\Omega$ with union equal to $\Omega$. When the boundary is Lipschitz, one can use a partition of unity $\left\{\zeta_{v}\right\}_{v=1}^{N}$, with each $\zeta_{v}$ supported in an open set $U_{v}$ such that $U_{v} \cap \Omega$ is star-shaped. We then regularize $\zeta_{v} f$ in $U_{v}$ as before. Thus, there exists a sequence $\alpha_{v} \in C_{p, q}^{\infty}(\Omega)$ with compact support in $\Omega$ such that $\alpha_{v} \longrightarrow f$ and $\bar{\partial} \alpha_{v} \longrightarrow 0$ in $W^{m}\left(\mathbb{C}^{n}\right)$. Applying Theorem 28 to each $\bar{\partial} \alpha_{v}$, we obtain $u_{v} \in L_{p, q}^{2}\left(\mathbb{C}^{n}\right)$ which is supported in $\bar{\Omega}$ and $\bar{\partial} u_{v}=\bar{\partial} \alpha_{v}$ in $\mathbb{C}^{n}$. Also using Theorem 28, we have

$$
\left\|u_{v}\right\|_{W^{m}\left(\Omega_{v}\right)} \leqslant C\left\|\bar{\partial} \alpha_{v}\right\|_{W^{m}\left(\Omega_{v}\right)} .
$$

Letting $h_{v}=\alpha_{v}-u_{v}$, we have that $h_{v} \in L_{p, q}^{2}\left(\mathbb{C}^{n}\right)$, which is supported in $\bar{\Omega}, \bar{\partial} h_{v}=0$ in $\mathbb{C}^{n}$ and $h_{v} \longrightarrow f$ in the $W^{m}\left(\mathbb{C}^{n}\right)$ norm. Thus the proof follows.

Theorem 31. Let $\Omega \in \mathbb{C}^{n}$ be a bounded $q$-pseudoconvex domain with Lipschitz boundary $b \Omega$ and let $1 \leqslant q \leqslant n$. Then, for $-\frac{1}{2} \leqslant m<\frac{1}{2}$ and for every $\alpha \in W_{p, q}^{m}\left(\mathbb{C}^{n}\right)$ which is supported in $\bar{\Omega}$ such that $\bar{\partial} \alpha=0$, for $1 \leqslant q \leqslant n-1$, one can find $u \in W_{p, q-1}^{m}\left(\mathbb{C}^{n}\right)$ such that $\bar{\partial} u=\alpha$ in the distribution sense in $\mathbb{C}^{n}$ with $u$ supported in $\bar{\Omega}$ and

$$
\begin{equation*}
\|u\|_{W_{p, q-1}^{m}}\left(\mathbb{C}^{n}\right) \leqslant C\|\alpha\|_{W_{p, q}^{m}\left(\mathbb{C}^{n}\right)}, \tag{31}
\end{equation*}
$$

where $C$ depends only on the Lipschitz constant and the diameter of $\Omega$, but is independent of $\alpha$.
Proof. For $m=0$, Theorem 31 follows from Theorem 28. Since $\Omega$ has Lipschitz boundary, there is a bounded extension operator from $W^{m}(\Omega) \longrightarrow W^{m}\left(\mathbb{C}^{n}\right)$ for $0 \leqslant m<\frac{1}{2}$ (see e.g. [11]). Let $\widetilde{\alpha} \in W_{p, q}^{m}\left(\mathbb{C}^{n}\right)$ be the extension of $\alpha$ so that $\left.\widetilde{\alpha}\right|_{\Omega}=\alpha$ with

$$
\left.\|\widetilde{\alpha}\|_{W^{m}} \mathbb{C}^{n}\right) \leqslant C\|\alpha\|_{W^{m}(\Omega)} .
$$

This is obviously true for $m=0$. Notice that this is not true for $m=\frac{1}{2}$ (see [28]).
We show that there is a bounded extension operator from $W^{-m}(\Omega) \longrightarrow W^{-m}\left(\mathbb{C}^{n}\right)$ for $0<m \leqslant \frac{1}{2}$ by using the fact that $\mathscr{D}(\Omega)$ is dense in $W^{m}(\Omega)$. Thus, from (5), for any $g \in W^{-m}(\Omega)$ and for $0<m \leqslant \frac{1}{2}$, we have

$$
\|\widetilde{\alpha}\|_{W^{-m}\left(\mathbb{C}^{n}\right)}=\sup _{g \in W^{0}(\Omega)} \frac{\left|\langle\widetilde{\alpha}, g\rangle_{\mathbb{C}^{n}}\right|}{\|g\|_{W^{m}\left(\mathbb{C}^{n}\right)}} \leqslant \sup _{g \in W^{0}(\Omega)} \frac{\left|\langle\alpha, g\rangle_{\Omega}\right|}{\|g\|_{W^{m}(\Omega)}}=\|\alpha\|_{W^{-m}(\Omega)} .
$$

Then for any $\alpha \in W^{-m}(\Omega),-\frac{1}{2} \leqslant m<\frac{1}{2}, \widetilde{\alpha}$ can be identified as a distribution in $W^{m}\left(\mathbb{C}^{n}\right)$ by setting $\alpha=0$ outside $\Omega$.

Assume that $\alpha \in W_{p, q}^{m}\left(\mathbb{C}^{n}\right)$ be any $\bar{\partial}$-closed form with support in $\bar{\Omega}$ and $0 \leqslant m<\frac{1}{2}$. Let $u$ be defined by (28). Using Theorem 27, one obtains

$$
\|u\|_{W^{m}(\Omega)} \leqslant C\|\star \alpha\|_{W^{m}(\Omega)}=C\|\alpha\|_{W^{m}(\Omega)}
$$

where $C$ depends only on the Lipschitz constant and the diameter of $\Omega$, but is independent of $\alpha$. Setting $u=0$ outside $\Omega$. It follows that $u$ is in $W^{m}\left(\mathbb{C}^{n}\right)$ and $u$ satisfies (31) for any $\alpha \in W_{p, q}^{m}\left(\mathbb{C}^{n}\right)$ with $0 \leqslant m<\frac{1}{2}$.

To show that $\bar{\partial} u=\alpha$ in $\mathbb{C}^{n}$, one use an approximation argument of Lemma 30. Let $\alpha_{v} \longrightarrow \alpha$ in $W_{p, q}^{m}\left(\mathbb{C}^{n}\right)$ such that the support of $\alpha_{v}$ is in $\bar{\Omega}$. Set

$$
u_{v}=\left\{\begin{array}{cl}
-\star \overline{\bar{\partial} N_{n-p, n-q} \star \bar{\alpha}_{v}} & \text { if } z \in \Omega,  \tag{32}\\
0 & \text { if } z \in \mathbb{C}^{n} \backslash \Omega .
\end{array}\right.
$$

Then, for each $j$, by applying Theorem 27 to obtain a $(p, q-1)$-form $u_{v}$ such that $\bar{\partial} u_{v}=\alpha_{v}$ in $\mathbb{C}^{n}$. From Lemma 30, $u_{v}$ converges to $u$ in $W_{p, q-1}^{m}(\Omega)$. Extending $u$ to be zero outside $\bar{\Omega}, u$ satisfies (31) and $\bar{\partial} u=f$ in the distribution sense in $\mathbb{C}^{n}$. This proves the theorem for $0 \leqslant m<\frac{1}{2}$.

Now, we prove the theorem for any $\bar{\partial}$-closed $\alpha \in W_{p, q}^{m}\left(\mathbb{C}^{n}\right)$ with support in $\bar{\Omega},-\frac{1}{2} \leqslant m<0$, by approximate $\alpha$ by $\bar{\partial}$-closed forms $\alpha_{v}$ in $L_{p, q}^{2}\left(\mathbb{C}^{n}\right)$ such that $\alpha_{v}$ has support in $\bar{\Omega}$ and $\alpha_{v} \longrightarrow \alpha$ in $W_{p, q}^{m}(\Omega)$. Set $u$ and $u_{v}$ as in (28) and (32). For each $j$, applying Theorem 28 to obtain a ( $p, q-1$ )form $u_{v}$ such that $\bar{\partial} u_{v}=\alpha_{v}$ in $\mathbb{C}^{n}$. Using Theorem 27 , one obtains

$$
\|u\|_{W^{m}(\Omega)} \leqslant C\|\star \alpha\|_{W^{m}(\Omega)}=C\|\alpha\|_{W^{m}(\Omega)}
$$

where $C$ depends only on the Lipschitz constant and the diameter of $\Omega$, but is independent of $\alpha$. From Lemma 30, $u_{v}$ converges to $u$ in $W_{p, q-1}^{m}(\Omega)$. Extending $u$ to be zero outside $\bar{\Omega}$. Then $u$ satisfies (31) and $\bar{\partial} u=\alpha$ in the distribution sense in $\mathbb{C}^{n}$. Thus the proof follows.

As Proposition 3.5 in [37], we prove the following result:
Proposition 32. Let $\Omega \Subset \mathbb{C}^{n}$ be a bounded $q$-pseudoconvex domain with Lipschitz boundary $b \Omega$ and let $1 \leqslant q \leqslant n$. Then, for any $\alpha \in W_{p, n}^{m}\left(\mathbb{C}^{n}\right),-\frac{1}{2} \leqslant m<\frac{1}{2}$, such that $\alpha$ is supported in $\bar{\Omega}$ and

$$
\begin{equation*}
\int_{\Omega} \alpha \wedge g=0 \text { for any } g \in C_{n-p, 0}^{\infty}(\bar{\Omega}) \cap \operatorname{ker} \bar{\partial} \tag{33}
\end{equation*}
$$

one can find $u \in W_{p, n-1}^{m}\left(\mathbb{C}^{n}\right)$ such that $\bar{\partial} u=\alpha$ in the distribution sense in $\mathbb{C}^{n}$ with $u$ supported in $\bar{\Omega}$ and

$$
\begin{equation*}
\left.\|u\|_{W^{m}\left(\mathbb{C}^{n}\right)} \leqslant C\|\alpha\|_{W^{m}} \mathbb{C}^{n}\right), \tag{34}
\end{equation*}
$$

where $C$ depends only on the Lipschitz constant and the diameter of $\Omega$, but is independent of $\alpha$.
Proof. Let $\alpha \in W_{p, n}^{m}\left(\mathbb{C}^{n}\right)$ such that $\alpha$ is supported in $\bar{\Omega}$ and $\alpha$ satisfies (33). Define $u$ as in (28) as follows

$$
\begin{equation*}
u=-\star \overline{\bar{\partial} N_{n-p, n-q} \star \bar{\alpha}} \tag{35}
\end{equation*}
$$

Now, we extend $u$ to $\mathbb{C}^{n}$ by defining $u=0$ in $\mathbb{C}^{n} \backslash \bar{\Omega}$. From Theorem 28, $u \in W_{p, n-1}^{m}\left(\mathbb{C}^{n}\right)$ and satisfies (34). From Lemma 30, we have $C_{n-p, 0}^{\infty}(\bar{\Omega}) \cap \operatorname{ker} \bar{\partial}$ is dense in $W_{n-p, 0}^{m}\left(\mathbb{C}^{n}\right) \cap \operatorname{ker}(\bar{\partial})$. Thus, Condition (33) implies that $P_{n-p, 0} \star \bar{\alpha}=0$.

To show that $\bar{\partial} u=\alpha$ in $\mathbb{C}^{n}$, we use an approximation argument as before. Let $\alpha_{v} \longrightarrow \alpha$ in $W_{p, n}^{m}\left(\mathbb{C}^{n}\right)$ such that the support of $\alpha_{v}$ is in $\bar{\Omega}$. Let

$$
h_{v}=\alpha_{v}-\star P_{n-p, 0} \star \bar{\alpha}_{v}
$$

Using Lemma 30, $h_{v}$ converges to $\alpha$ in $W_{n-p, 0}^{m}(\Omega)$. Each $h_{v}$ is in $L_{p, n}^{2}(\Omega)$ and satisfies

$$
\int_{\Omega} h_{v} \wedge g=0 \text { for any } g \in L_{n-p, 0}^{2}(\Omega) \cap \operatorname{ker} \bar{\partial}
$$

Define

$$
u_{v}=\left\{\begin{array}{cl}
-\star \overline{\bar{\partial} N_{n-p, 0} \star \bar{h}_{v}} & \text { if } z \in \Omega \\
0 & \text { if } z \in \mathbb{C}^{n} \backslash \Omega
\end{array}\right.
$$

Then, for each $j$, by applying Proposition 29 to obtain a ( $p, n-1$ )-form $u_{v}$ such that $\bar{\partial} u_{v}=h_{v}$ in $\mathbb{C}^{n}$ and

$$
\left\|u_{v}\right\|_{W^{m}\left(\mathbb{C}^{n}\right)} \leqslant C\left\|h_{v}\right\|_{W^{m}\left(\mathbb{C}^{n}\right)} .
$$

From Lemma 30, $u_{v}$ converges to $u$ in $W_{p, n-1}^{m}(\Omega)$. Extending $u$ to be zero outside $\bar{\Omega}, u$ satisfies (34) and $\bar{\partial} u=\alpha$ in the distribution sense in $\mathbb{C}^{n}$. Thus the proof follows.

## 5. Proof of the main theorems

Following the construction in [37, Lemma 4.1]. Let $B$ be a large ball in $\mathbb{C}^{n}$ such that $\bar{\Omega} \Subset B$. Let $\Omega^{+}=B \backslash \bar{\Omega}$ and $\Omega^{-}=\Omega$. In [19], a Martinelli-Bochner-Koppelman type kernel is constructed, and in [36] it is shown that the transformation induced by this kernel satisfies a jump formula (10). As a result, there exists an integral kernel $K_{q}(\zeta, z)$ of type ( $p, q$ ) in $z$ and $(n-p, n-q-1)$ in $\zeta$ satisfying a Martinelli-Bochner-Koppelman formula

$$
\int_{b \Omega} K_{q}(\zeta, z) \wedge \alpha(\zeta)= \begin{cases}\alpha^{+}(z) & \text { if } z \in \Omega^{+}, \\ \alpha^{-}(z) & \text { if } z \in \Omega^{-},\end{cases}
$$

where $\alpha^{+}(z)=K^{+} \alpha(z)$ if $z \in \Omega^{+}$and $\alpha^{-}(z)=K^{-} \alpha(z)$ if $z \in \Omega$ (see [19, Section 2.3]). Let $\alpha \in$ $L_{p, q}^{2}(b \Omega), 0<q \leqslant n-1$, and $\bar{\partial}_{b} \alpha=0$ in $b \Omega$. If, we extend $\alpha^{+}(z)$ and $\alpha^{-}(z)$ to $b \Omega$ by considering non-tangential limits, we have the jump formula

$$
\begin{equation*}
\tau\left(\alpha^{+}(z)-\alpha^{-}(z)\right)=\alpha(z) \text { on } b \Omega, \tag{36}
\end{equation*}
$$

where $\tau$ is defined in (9). Since $\alpha^{+}$and $\alpha^{-}$have non-tangential boundary values in $L^{2}$, then from Lemma 4.1 in [37], $\alpha^{+} \in W_{p, q}^{\frac{1}{2}}\left(\Omega^{+}\right)$and $\alpha^{-} \in W_{p, q}^{\frac{1}{2}}(\Omega)$ such that $\bar{\partial} \alpha^{+}=0$ in $\Omega^{+}$and $\bar{\partial} \alpha^{-}=0$ in $\Omega$. Furthermore, we have

$$
\begin{equation*}
\left\|\alpha^{+}\right\|_{W^{\frac{1}{2}\left(\Omega^{+}\right)}} \leqslant C\|\alpha\|_{b \Omega}, \tag{37}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|\alpha^{-}\right\|_{W^{\frac{1}{2}(\Omega)}} \leqslant C\|\alpha\|_{b \Omega}, \tag{38}
\end{equation*}
$$

where the constant $C$ depends only on the Lipschitz constant of $\Omega$ but is independent of $\alpha$.
Since $\Omega^{+}$is a bounded Lipschitz domain, there exists a continuous linear operator $E$ from $W^{m}\left(\Omega^{+}\right)$into $W^{m}\left(\mathbb{C}^{n}\right)$, for any $m \geqslant 0$, such that for any $g \in W^{m}\left(\Omega^{+}\right)$,

$$
\left.E g\right|_{\Omega^{+}}=g
$$

(see Stein [38, Chapter 6] or Grisvard [11]).
Extend $\alpha^{+}$from $\Omega^{+}$to $E \alpha^{+}$componentwise on $B$ such that

$$
\left\|E \alpha^{+}\right\|_{W^{\frac{1}{2}(B)}} \leqslant C\left\|\alpha^{+}\right\|_{W^{\frac{1}{2}}\left(\Omega^{+}\right)} .
$$

Since $\bar{\partial} \alpha^{+}=0$ on $\Omega^{+}$, we have $\bar{\partial} E \alpha^{+}=0$ on $\Omega^{+}$. Set $f=\bar{\partial} E \alpha^{+}$in $\Omega$ and zero outside. Then $\bar{\partial} f=0$ and $f$ is supported in $\bar{\Omega}$. Furthermore, we have

$$
\begin{equation*}
\|f\|_{W^{-\frac{1}{2}}(\Omega)} \leqslant C\left\|E \alpha^{+}\right\|_{W^{\frac{1}{2}(B)}} \leqslant C\left\|\alpha^{+}\right\|_{W^{\frac{1}{2}\left(\Omega^{+}\right)}} . \tag{39}
\end{equation*}
$$

When $1 \leqslant q \leqslant n-2$, one defines

$$
h=\left\{\begin{array}{cl}
-\star \overline{\bar{\partial} N_{n-p, n-q} \star \bar{f}} & \text { if } z \in \Omega,  \tag{40}\\
0 & \text { if } z \in \mathbb{C}^{n} \backslash \Omega .
\end{array}\right.
$$

It follows, from Theorem 31, that $\bar{\partial} h=f$ and $h$ is supported in $\bar{\Omega}$.
Also from (16) and (39), one obtains

$$
\begin{equation*}
\|h\|_{W^{-\frac{1}{2}}(\Omega)} \leqslant C\|f\|_{W^{-\frac{1}{2}}(\Omega)} \leqslant C\left\|\alpha^{+}\right\|_{W^{\frac{1}{2}}\left(\Omega^{+}\right)^{\prime}} . \tag{41}
\end{equation*}
$$

Set

$$
\widetilde{\alpha}^{+}=E \alpha^{+}-h \text { in } B .
$$

Then $\widetilde{\alpha}^{+}=\alpha^{+}-h=\alpha^{+}$in $\Omega^{+}$. Since $f=\bar{\partial} E \alpha^{+}$in $\Omega$, we have $\bar{\partial} \widetilde{\alpha}^{+}=0$ in $B$. Then from (41) and Theorem 27, one obtains

$$
\begin{aligned}
\left\|\widetilde{\alpha}^{+}\right\|_{W^{-\frac{1}{2}(B)}} & \leqslant C\left(\left\|E \alpha^{+}\right\|_{W^{-\frac{1}{2}(B)}}+\|h\|_{W^{-\frac{1}{2}}(\Omega)}\right) \\
& \leqslant C\left\|\alpha^{+}\right\|_{W^{\frac{1}{2}}\left(\Omega^{+}\right)^{\prime}}
\end{aligned}
$$

for some constant $C$ independent of $\alpha$. Thus, for $1 \leqslant q \leqslant n-2$, one obtains

$$
\begin{equation*}
\left\|\widetilde{\alpha}^{+}\right\|_{W^{-\frac{1}{2}(B)}} \leqslant C\left\|\alpha^{+}\right\|_{W^{\frac{1}{2}}\left(\Omega^{+}\right)} \tag{42}
\end{equation*}
$$

When $\underline{q}=n-\underline{1}$, notice that $\bar{\partial} \alpha^{-}=0$ in $\Omega$. By using the jump formula (10), for every $\phi \in$ $C_{n-p, 0}^{\infty}(\bar{\Omega}) \cap \operatorname{ker} \bar{\partial}$, one obtains

$$
\begin{aligned}
\int_{\Omega} f \wedge \phi & =\int_{\Omega} \bar{\partial}\left(E \alpha^{+} \wedge \phi\right) \\
& =\int_{b \Omega} K_{b}^{+} \alpha \wedge \phi \\
& =\int_{b \Omega} K_{b}^{-} \alpha \wedge \phi+\int_{b \Omega} \alpha \wedge \phi \\
& =\int_{\Omega} \bar{\partial} \alpha^{-} \wedge \phi+\int_{b \Omega} \alpha \wedge \phi \\
& =\int_{b \Omega} \alpha \wedge \phi \\
& =0
\end{aligned}
$$

This implies that $f=\bar{\partial} E \alpha^{+}$satisfies Condition (33). Defining $h$ as in (40) and using Proposition 32, one obtains that $\bar{\partial} h=f$ in $\mathbb{C}^{n}$ and $h$ is supported in $\bar{\Omega}$. Repeating the arguments as before, thus estimate (42) holds also for $q=n-1$. Thus, for $1 \leqslant q \leqslant n-1$, one obtains

$$
\left\|\widetilde{\alpha}^{+}\right\|_{W^{-\frac{1}{2}}(B)} \leqslant C\left\|\alpha^{+}\right\|_{W^{\frac{1}{2}}\left(\Omega^{+}\right)}
$$

Then, from (37) and for $1 \leqslant q \leqslant n-1$, one obtains

$$
\begin{equation*}
\left\|\widetilde{\alpha}^{+}\right\|_{W^{-\frac{1}{2}}(B)} \leqslant C\|\alpha\|_{b \Omega} . \tag{43}
\end{equation*}
$$

Since $\alpha$ satisfies (36) and $\bar{\partial} \alpha^{+}=0$ in $\Omega^{+}$and $\bar{\partial} \alpha^{-}=0$ in $\Omega$, then $\bar{\partial}_{b} \alpha=0$ in $b \Omega$. Let

$$
u=u^{+}-u^{-} \text {on } b \Omega
$$

First we solve $\bar{\partial} u^{+}=\alpha^{+}$on $\Omega^{+}$. To do that, we use the canonical solution operator to define

$$
u^{+}=\bar{\partial}^{*} N^{B} \widetilde{\alpha}^{+}
$$

Then $\bar{\partial} u^{+}=\widetilde{\alpha}^{+}=0$ in $B$. Using the interior regularity for $N^{B}, u^{+}$gains a derivative on compact subsets of $B$, so $u^{+}$has components in $W^{\frac{1}{2}}(b \Omega)$ on a neighborhood of $\bar{\Omega}$. From (43), we have for any $\xi \in \mathscr{D}(B)$ such that $\xi=1$ in a neighborhood of $\bar{\Omega}$,

$$
\begin{equation*}
\left\|\xi u^{+}\right\|_{W^{\frac{1}{2}(B)}} \leqslant C\left\|\alpha^{+}\right\|_{W^{-\frac{1}{2}}(B)} \leqslant C\|\alpha\|_{b \Omega} \tag{44}
\end{equation*}
$$

for some constant $C$ independent of $\alpha$. Restricting $u^{+}$to $b \Omega$, using (44) and the trace theorem again, one obtains

$$
\begin{equation*}
\left\|u^{+}\right\|_{b \Omega} \leqslant C\left\|\xi \alpha^{+}\right\|_{W^{\frac{1}{2}(B)}} \leqslant C\|\alpha\|_{b \Omega} \tag{45}
\end{equation*}
$$

for some constant $C$ independent of $\alpha$.
To solve $\bar{\partial} u^{-}=\alpha^{-}$on $\Omega$, one defines

$$
u^{-}=\bar{\partial}^{*} N^{\Omega} \alpha^{-}
$$

Then $\bar{\partial} u^{-}=\alpha^{-}$. It follows from (38) and Theorem 27 that

$$
\left\|u^{-}\right\|_{W^{\frac{1}{2}(\Omega)}} \leqslant C\left\|\alpha^{-}\right\|_{W^{\frac{1}{2}(\Omega)}} \leqslant C\|\alpha\|_{b \Omega}
$$

for some constant $C$ independent of $\alpha$.

Since $u^{-}$satisfies a system of elliptic equations, it can be treated like harmonic functions. Thus $u^{-}$has boundary value in $L^{2}(b \Omega)$. Restricting $u^{-}$to $b \Omega$, using the trace theorem for smooth domains (see [22]), we have

$$
\begin{equation*}
\left\|u^{-}\right\|_{b \Omega} \leqslant C\left\|u^{-}\right\|_{W^{\frac{1}{2}(\Omega)}} \leqslant C\|\alpha\|_{b \Omega} \tag{46}
\end{equation*}
$$

Then $\bar{\partial}_{b} u=v$ on $b \Omega$. Also from (45) and (46), we have

$$
\|u\|_{b \Omega} \leqslant C\|\alpha\|_{b \Omega}
$$

for some constant $C$ independent of $\alpha$. Thus the proof of Theorem 1 follows.
Now, we prove Theorems 2 and 3. The proof of (i) in Theorem 2 follows directly from the main theorem and Theorem 1.1.1 in [21]. To prove (ii) in Theorem 2, we need to claim that, for all $0 \leqslant p \leqslant n, 1 \leqslant q \leqslant n-2$,

$$
\begin{equation*}
\mathscr{H}_{b}^{p, q}(b \Omega)=\operatorname{ker} \bar{\partial}_{b} \cap \operatorname{ker} \bar{\partial}_{b}^{*}=\{0\} \tag{47}
\end{equation*}
$$

To prove (47), let $\alpha \in \mathscr{H}_{b}^{p, q}(b \Omega)=\operatorname{ker} \bar{\partial}_{b} \cap \operatorname{ker} \bar{\partial}_{b}^{*}$. Then $\alpha=\bar{\partial}_{b} u$ for some $u \in L_{p, q-1}^{2}(b \Omega)$ by the main theorem. Since $\alpha \in \operatorname{ker} \bar{\partial}_{b}^{*}$, then

$$
\langle\alpha, \alpha\rangle_{b \Omega}=\left\langle\bar{\partial}_{b} u, \alpha\right\rangle_{b \Omega}=\left\langle u, \bar{\partial}_{b}^{*} \alpha\right\rangle_{b \Omega}=0
$$

We have $\mathscr{H}_{b}^{p, q}(b \Omega)=\{0\}$.
From Theorem 2, the range of $\bar{\partial}_{b}$, denoted by $\mathscr{R}$ ang $\bar{\partial}_{b}$, is closed in every degree. Then, one obtains ker $\bar{\partial}_{b}=\mathscr{R}$ ang $\bar{\partial}_{b}^{*}$ and the following orthogonal decomposition:

$$
L_{p, q}^{2}(b \Omega)=\operatorname{ker} \bar{\partial}_{b} \oplus \mathscr{R} \operatorname{ang} \bar{\partial}_{b}^{*}=\mathscr{R} \operatorname{ang} \bar{\partial}_{b}^{*} \oplus \mathscr{R} \operatorname{ang} \bar{\partial}_{b}^{*}
$$

Repeating the arguments of Theorem 8.4.10 in [6], one can prove that for every $\alpha \in \operatorname{Dom} \bar{\partial}_{b} \cap$ Dom $\bar{\partial}_{b}^{*}$,

$$
\begin{aligned}
\|\alpha\|_{b \Omega}^{2} & \leqslant C\left(\left\|\bar{\partial}_{b} \alpha\right\|_{b \Omega}^{2}+\left\|\bar{\partial}_{b}^{*} \alpha\right\|_{b \Omega}^{2}\right) \\
& =C\left\langle\square_{b} \alpha, \alpha\right\rangle_{b \Omega} \\
& \leqslant C\left\|\square_{b} \alpha\right\|\|\alpha\|_{b \Omega}
\end{aligned}
$$

i.e.,

$$
\begin{equation*}
\|\alpha\|_{b \Omega} \leqslant C\left\|\square_{b} \alpha\right\|_{b \Omega} \tag{48}
\end{equation*}
$$

Since $\square_{b}$ is a linear closed densely defined operator, then, from Theorem 1.1.1 in [21], $\mathscr{R}$ ang $\square_{b}$ is closed. Thus, from (1.1.1) in [21] and the fact that $\square_{b}$ is self adjoint, we have the Hodge decomposition

$$
L_{p, q}^{2}(b \Omega)=\mathscr{R} \operatorname{ang} \square_{b} \oplus \mathscr{H}_{b}^{p, q}(b \Omega)=\bar{\partial}_{b} \bar{\partial}_{b}^{*} \operatorname{Dom} \square_{b} \oplus \bar{\partial}_{b}^{*} \bar{\partial}_{b} \operatorname{Dom} \square_{b}
$$

Since $\square_{b}$ is one to one on Dom $\square_{b}$ from (48), then there exists a unique bounded inverse operator

$$
N_{b}: \mathscr{R} \operatorname{ang} \square_{b} \longrightarrow \operatorname{Dom} \square_{b} \cap\left(\mathscr{H}_{b}^{p, q}(b \Omega)\right)^{\perp}
$$

such that $N_{b} \square_{b} \alpha=\alpha$ on $\operatorname{Dom} \square_{b}$. We can write $N_{b} \square_{b}=I$ on $\operatorname{Dom} \square_{b} \cap\left(\not \mathscr{H}_{b}^{p, q}(b \Omega)\right)^{\perp}$. From the definition of $N_{b}$, extend $N_{b}$ to $L_{p, q}^{2}(b \Omega)$ to obtain $\square_{b} N_{b}=I$ on $L_{p, q}^{2}(b \Omega)$. Thus $N_{b}$ satisfies (i) and (ii).

To show that $\bar{\partial}_{b}^{*} N_{b}=N_{b} \bar{\partial}_{b}^{*}$ on $\operatorname{Dom} \bar{\partial}_{b}^{*}$. By using (ii), we have $\bar{\partial}_{b}^{*} \alpha=\bar{\partial}_{b}^{*} \bar{\partial}_{b} \bar{\partial}_{b}^{*} N_{b} \alpha$, for $\alpha \in$ Dom $\bar{\partial}_{b}^{*}$. Thus

$$
N_{b} \bar{\partial}_{b}^{*} \alpha=N_{b} \bar{\partial}_{b}^{*} \bar{\partial}_{b} \bar{\partial}_{b}^{*} N_{b} \alpha=N_{b}\left(\bar{\partial}_{b}^{*} \bar{\partial}_{b}+\bar{\partial}_{b} \bar{\partial}_{b}^{*}\right) \bar{\partial}_{b}^{*} N_{b} \alpha=\bar{\partial}_{b}^{*} N_{b} \alpha
$$

A similar arguments shows that $\bar{\partial}_{b} N_{b}=N_{b} \bar{\partial}_{b}$ on $\operatorname{Dom} \bar{\partial}_{b}$. Thus the proof of Theorem 3 follows.

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