# On the vibrations of a string with a concentrated mass and rapidly oscillating density 

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## A R T I C L E I N F O

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#### Abstract

The paper is devoted to the vibrations of a string $I$ with a concentrated mass $\varepsilon^{-1} \varkappa Q\left(\varepsilon^{-1} x\right)$ and rapidly oscillating density $q\left(x, \mu^{-1} x\right)$, where $q(x, \zeta)$ is a 1-periodic in $\zeta$ function, $Q(\xi)$ is a function with compact support, the integral of which is equal to one, $0 \in I, \mu, \varepsilon$ are small positive parameters, $x \in \mathbb{R}$. By combining homogenization and the method of matched asymptotic expansions, we construct solutions to the problems up to $O(\varepsilon+\mu)$.


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## 1. Introduction

Studies on the vibrations of a string with inhomogeneous rapidly oscillating density have been conducted for a long time (see, for instance, [1-3]). In these monographs, using the homogenization method, the authors constructed asymptotics with respect to a small parameter (the period of the oscillations) of solutions to the boundary value problem.

The vibration of a string with a concentration of mass on a small set has been studied by means of other methods for a long time, too. The convergence of solutions to the respective boundary value problem was investigated in [4], and the asymptotics with respect to a small parameter was constructed in [5] by means of the method of matched asymptotic expansions [6]. The analogous spectral problem for the Laplace operator in 3D space was considered in [7]. Note that a close mathematical problem on the convergence of the Schrödinger operator with $\delta$-type potential on the axis was studied in [8].

In the present paper, using the combination of homogenization and the method of matching of asymptotic expansions [9], we study the case when the string with a concentrated mass has an inhomogeneous rapidly oscillating density. We get the homogenized (limit) solution to a boundary value problem, up to $O(\mu+\varepsilon)$, where $\mu$ is a period of rapid oscillations of the density, and $\varepsilon$ is the order of the length of the small part of the string, on which the concentrated mass is located.

## 2. Settings of the problem and main result

Let $I$ be $(a, b),\{0\} \in I$. Consider the problem

$$
\begin{align*}
& \mathcal{L}_{\mu, \varepsilon} u_{\mu, \varepsilon}:=-\frac{\mathrm{d}^{2} u_{\mu, \varepsilon}}{\mathrm{d} x^{2}}+\left(q\left(x, \frac{x}{\mu}\right)+\varepsilon^{-1} \varkappa Q\left(\frac{x}{\varepsilon}\right)\right) u_{\mu, \varepsilon}=f(x), \quad x \in I \\
& l_{a} u_{\mu, \varepsilon}:=h_{a} u_{\mu, \varepsilon}(a)-H_{a} u_{\mu, \varepsilon}^{\prime}(a)=0 \quad l_{b} u_{\mu, \varepsilon}:=h_{b} u_{\mu, \varepsilon}(b)+H_{b} u_{\mu, \varepsilon}^{\prime}(b)=0 \tag{1}
\end{align*}
$$

[^0]where $q(x, \zeta)$ is a 1-periodic in $\zeta$ function belonging to $C^{2,0}(\bar{I} \times(-\infty, \infty)), q(x, \zeta)>0, Q(\xi) \in C_{0}(-\infty, \infty)$,
$$
\int_{-\infty}^{\infty} Q(\xi) \mathrm{d} \xi=1
$$
$0<\mu, \varepsilon \ll 1, h_{a}, h_{b}, H_{a}, H_{b} \geqslant 0, h_{a}+H_{a}>0, h_{b}+H_{b}>0, f \in L^{2}(I)$. Although, for the concentrated mass, $x>0$ and $Q(\xi) \geqslant 0$, we consider a more general case, when $x \in \mathbb{R}$ and $Q(\xi)$ can take negative values as well. We will assume that $\operatorname{supp} Q(\xi) \subset[-1,1]$.

For functions $G(x)$, we use the notation $\{G\}(0)=G(+0)-G(-0)$, and for 1-periodic in $\zeta$ functions $g(x, \zeta)$ we denote

$$
[g](x):=\int_{0}^{1} g(x, \zeta) \mathrm{d} \zeta
$$

Since $q>0$, the boundary value problems

$$
\begin{aligned}
\mathcal{L} u_{0} & :=-\frac{\mathrm{d}^{2} u_{0}}{\mathrm{~d} x^{2}}+[q](x) u_{0}=f, \quad x \in I, \quad & & l_{a} u_{0}=l_{b} u_{0}=0 \\
\mathcal{L} y & =0, \quad x \in I \backslash\{0\}, \quad l_{a} y=l_{b} y=0, & & y(0)=1
\end{aligned}
$$

have unique solutions in $H^{2}(I)$ and $H^{1}(I) \cap H^{2}(a, 0) \cap H^{2}(0, b)$, respectively. It is easy to see that if $x \neq \varkappa_{0}:=\left\{y^{\prime}\right\}(0)$, then the function

$$
\begin{equation*}
u^{0}(x):=u_{0}(x)+\frac{\varkappa u_{0}(0)}{\varkappa_{0}-\varkappa} y(x) \tag{2}
\end{equation*}
$$

solves uniquely the boundary value problem

$$
\begin{equation*}
\mathcal{L} u^{0}=f, \quad x \in I \backslash\{0\}, \quad l_{a} u^{0}=l_{b} u^{0}=0, \quad\left\{\left(u^{0}\right)^{\prime}\right\}(0)=\varkappa u^{0}(0) \tag{3}
\end{equation*}
$$

for any $f \in L^{2}(I)$ in $H^{1}(I) \cap H^{2}(a, 0) \cap H^{2}(0, b)$.
The main result of the paper is the proof of the following proposition.
Theorem 2.1. Assume that $\varkappa \neq \varkappa_{0}$. Then for sufficiently small $\varepsilon$ and $\mu$, the solution to boundary value problem (1) is uniquely determined in $H^{2}(I)$ and satisfies the uniform estimate

$$
\begin{equation*}
\left\|u_{\mu, \varepsilon}-u^{0}\right\|_{L^{2}(I)} \leqslant(\mu+\varepsilon) C\|f\|_{L^{2}(I)} \tag{4}
\end{equation*}
$$

where $u^{0}(x)$ is a solution to boundary value problem (3).

## 3. Preliminaries and auxiliary assertions

Consider the following boundary value problems:

$$
\begin{align*}
& \mathcal{L}_{\mu} u_{\mu}:=-\frac{\mathrm{d}^{2} u_{\mu}}{\mathrm{d} x^{2}}+q\left(x, \frac{x}{\mu}\right) u_{\mu}=f, \quad x \in I, \quad l_{a} u_{\mu}=l_{b} u_{\mu}=0  \tag{5}\\
& \mathcal{L}_{\mu} y_{\mu}=0, \quad x \in I \backslash\{0\}, \quad l_{a} y_{\mu}=l_{b} y_{\mu}=0, \quad y_{\mu}(0)=1  \tag{6}\\
& \mathcal{L}_{\mu} u^{\mu}=f, \quad x \in I \backslash\{0\}, \quad l_{a} u^{\mu}=l_{b} u^{\mu}=0, \quad\left\{\left(u^{\mu}\right)^{\prime}\right\}(0)=x u^{\mu}(0) \tag{7}
\end{align*}
$$

Since $q>0$, the boundary value problems (5) and (6) have unique solutions from $H^{2}(I)$ and $H^{1}(I) \cap C^{2}[a, 0] \cap C^{2}[0, b]$, respectively, and the estimate

$$
\begin{equation*}
\left\|u_{\mu}\right\|_{H^{2}(I)} \leqslant C\|f\|_{L^{2}(I)} \tag{8}
\end{equation*}
$$

holds true.
It is easy to see that if $\varkappa \neq \varkappa_{\mu}:=\left\{y_{\mu}^{\prime}\right\}(0)$, then the function

$$
\begin{equation*}
u^{\mu}(x):=u_{\mu}(x)+\frac{\varkappa u_{\mu}(0)}{\varkappa_{\mu}-\varkappa} y_{\mu}(x) \in H^{1}(I) \cap H^{2}(a, 0) \cap H^{2}(0, b) \tag{9}
\end{equation*}
$$

is a unique solution to problem (7) for any $f \in L^{2}(I)$, in addition, because of (8) and the embedding of $H^{1}(I)$ in $C(\bar{I})$ (see, for instance, [10, Chapter III, §6]), which satisfies

$$
\begin{equation*}
\left\|u^{\mu}\right\|_{H^{2}(a, 0)}+\left\|u^{\mu}\right\|_{H^{2}(0, b)} \leqslant C \frac{\|f\|_{L^{2}(I)}}{\left|\varkappa_{\mu}-\varkappa\right|} \tag{10}
\end{equation*}
$$

Since $q(x, \zeta) \in C^{2,0}(\bar{I} \times(-\infty, \infty))$ and $q(x, \zeta)>0$, then the construction of asymptotics of the functions $u_{\mu}(x)$ and $y_{\mu}(x)$ by the homogenization method leads to:

$$
\begin{equation*}
\left\|u_{\mu}-u_{0}\right\|_{H^{1}(I)} \leqslant C \mu\|f\|_{L^{2}(I)}, \quad\left\|y_{\mu}-y\right\|_{C^{1}[a, 0]}+\left\|y_{\mu}-y\right\|_{C^{1}[0, b]}=O(\mu) \tag{11}
\end{equation*}
$$

From (11) and the embedding theorems, it follows that

$$
\begin{equation*}
\left|u_{\mu}(0)-u_{0}(0)\right| \leqslant \mu C\|f\|_{L^{2}(I)}, \quad x_{\mu}=\varkappa_{0}+O(\mu), \quad\left\|y_{\mu}-y\right\|_{H^{1}(I)}=O(\mu) \tag{12}
\end{equation*}
$$

Due to (2), (9), (11) and (12), we derive for $\varkappa \neq \varkappa_{0}$ and sufficiently small $\mu$ the estimate

$$
\begin{equation*}
\left\|u^{\mu}-u^{0}\right\|_{H^{1}(I)} \leqslant C \mu\|f\|_{L^{2}(I)} \tag{13}
\end{equation*}
$$

## 4. Existence and uniqueness of a solution to boundary value problem (1)

Let $\widetilde{h}_{a}$ be equal to $h_{a} H_{a}^{-1}$, if $H_{a} \neq 0$, and $\widetilde{h}_{a}=0$, if $H_{a}=0$. We define in an analogous way $\widetilde{h}_{b}$. The quadratic forms of the boundary value problems (5), (7) and (1) read as

$$
\begin{aligned}
\mathfrak{h}_{\mu}[u] & =\int_{I}\left(\left(u^{\prime}(x)\right)^{2}+q\left(x, \frac{x}{\mu}\right) u^{2}(x)\right) \mathrm{d} x+\widetilde{h}_{a} u^{2}(a)+\widetilde{h}_{b} u^{2}(b), \\
\mathfrak{h}^{\mu}[u] & =\mathfrak{h}_{\mu}[u]+\varkappa u^{2}(0), \\
\mathfrak{h}_{\mu, \varepsilon}[u] & =\mathfrak{h}^{\mu}[u]+\varkappa \varepsilon^{-1} \int_{I} Q\left(\frac{x}{\varepsilon}\right) u^{2}(x) \mathrm{d} x-\varkappa u^{2}(0)
\end{aligned}
$$

correspondingly, where $u(a)=0$, if $H_{a}=0$, and $u(b)=0$, if $H_{b}=0$.
Let us show the validity of the following estimate

$$
\begin{equation*}
\left|\mathfrak{h}_{\mu, \varepsilon}[u]-\mathfrak{h}^{\mu}[u]\right| \leqslant M_{1} \varepsilon^{1 / 2}\|u\|_{L^{2}(I)}+M_{2} \varepsilon^{1 / 2} \mathfrak{h}^{\mu}[u] \tag{14}
\end{equation*}
$$

Denote $v(x)=u(x)-u(0)$. Then $v(0)=0, v^{\prime}(x)=u^{\prime}(x)$,

$$
\int_{I} \varepsilon^{-1} Q\left(\frac{x}{\varepsilon}\right) u^{2}(x) \mathrm{d} x-u^{2}(0)=2 \int_{I} \varepsilon^{-1} Q\left(\frac{x}{\varepsilon}\right) u(0) v(x) \mathrm{d} x+\int_{I} \varepsilon^{-1} Q\left(\frac{x}{\varepsilon}\right) v^{2}(x) \mathrm{d} x
$$

The Friedrichs-Steklov inequality leads to the estimate

$$
\begin{aligned}
2\left|\int_{I} \varepsilon^{-1} Q\left(\frac{x}{\varepsilon}\right) u(0) v(x) \mathrm{d} x\right| & \leqslant 2 \int_{-\varepsilon}^{\varepsilon} \varepsilon^{-1} \max _{\mathbb{R}}|Q| \cdot|u(0) v(x)| \mathrm{d} x \\
& \leqslant \int_{-\varepsilon}^{\varepsilon} \varepsilon^{-1} \max _{\mathbb{R}}|Q|\left(\varepsilon^{1 / 2} u^{2}(0)+\varepsilon^{-1 / 2} v^{2}(x)\right) \mathrm{d} x \\
& \leqslant \max _{\mathbb{R}}|Q| \varepsilon^{1 / 2}\left(2 u^{2}(0)+\int_{-\varepsilon}^{\varepsilon}\left(v^{\prime}\right)^{2}(x) \mathrm{d} x\right)
\end{aligned}
$$

So, we derive

$$
\begin{align*}
\left|\varepsilon^{-1} \int_{I} Q\left(\frac{x}{\varepsilon}\right) u^{2}(x) \mathrm{d} x-u^{2}(0)\right| & \leqslant 2 \max _{\mathbb{R}}|Q| \varepsilon^{1 / 2} u^{2}(0)+\max _{\mathbb{R}}|Q|\left(\varepsilon^{1 / 2}+\varepsilon\right) \int_{-\varepsilon}^{\varepsilon}\left(u^{\prime}\right)^{2}(x) \mathrm{d} x \\
& \leqslant C \varepsilon^{1 / 2}\|u\|_{H^{1}(I)}^{2}, \quad u \in H^{1}(I) \tag{15}
\end{align*}
$$

Due to the definition of the quadratic forms and estimate (15), we get that the forms $\mathfrak{h}_{\mu, \varepsilon}, \mathfrak{h}^{\mu}$ are bounded from below, and for $x \geqslant 0$, inequality (14) holds.

Let $x<0$. Consider the case $H_{a}+H_{b} \neq 0$. It is known (see, for instance, [10, Chapter, $\left.\S 4\right]$ ) that for any fixed interval $I^{\prime}=\left(a^{\prime}, b^{\prime}\right)$ such that $\bar{I} \subset I^{\prime}$ and any $u \in H^{1}(I)$, there exists a continuation $U \in H^{1}\left(I^{\prime}\right)$ such that

$$
\begin{equation*}
U\left(a^{\prime}\right)=U\left(b^{\prime}\right)=0, \quad\|U\|_{H^{1}\left(I^{\prime}\right)} \leqslant M\|u\|_{H^{1}(I)}, \quad\|U\|_{L^{2}\left(I^{\prime}\right)} \leqslant M\|u\|_{L^{2}(I)} \tag{16}
\end{equation*}
$$

Since $\alpha\|u\|_{H^{1}(I)}^{2} \leqslant \mathfrak{h}_{\mu}[u], \alpha>0$, due to (16) we derive

$$
\begin{aligned}
h^{\mu}[u] & =h_{\mu}[u]+\varkappa u^{2}(0) \geqslant \alpha\|u\|_{H^{1}(I)}^{2}+\varkappa u^{2}(0) \geqslant \frac{\alpha}{M}\|U\|_{H^{1}\left(I^{\prime}\right)}^{2}+\varkappa U^{2}(0) \\
& =\frac{\alpha}{M}\|U\|_{H^{1}\left(I^{\prime}\right)}^{2}+2 \varkappa \int_{a^{\prime}}^{0} U^{\prime}(x) U(x) \mathrm{d} x \geqslant \frac{\alpha}{M}\|U\|_{H^{1}\left(I^{\prime}\right)}^{2}+\varkappa\left(\gamma\|U\|_{L^{2}\left(I^{\prime}\right)}^{2}+\gamma^{-1}\left\|U^{\prime}\right\|_{L^{2}\left(I^{\prime}\right)}^{2}\right) \\
& \geqslant\left(\frac{\alpha}{M}+\varkappa \gamma^{-1}\right)\|U\|_{H^{1}\left(I^{\prime}\right)}^{2}+\varkappa \gamma\|U\|_{L^{2}\left(I^{\prime}\right)}^{2} \geqslant\left(\frac{\alpha}{M}+\varkappa \gamma^{-1}\right)\|u\|_{H^{1}(I)}^{2}+\varkappa \gamma M\|u\|_{L^{2}(I)}^{2}
\end{aligned}
$$

So, for large $\gamma>0$ we have

$$
\|u\|_{H^{1}(I)}^{2} \leqslant K_{1}\|u\|_{L^{2}(I)}^{2}+K_{2} h^{\mu}[u], \quad K_{i}>0
$$

This estimate and inequality (15) imply estimate (14). In case $\varkappa<0, H_{a}=H_{b}=0$ the proof is similar and more simple.
Since $\varkappa_{\mu} \underset{\mu \rightarrow 0}{\rightarrow} \varkappa_{0}$ (see, (12)), then estimate (14) by means of [11, Chapter VI, Theorem 3.9] shows that, for $\varkappa \neq \varkappa_{0}$ and sufficiently small $\mu$, boundary value problem (1) has a unique solution in $H^{2}(I)$. Moreover, for solutions to boundary value problems (1) and (7), the following uniform estimate in $\mu$ :

$$
\begin{equation*}
\left\|u_{\mu, \varepsilon}-u^{\mu}\right\|_{L^{2}(I)} \leqslant C \varepsilon^{1 / 2}\|f\|_{L^{2}(I)} \tag{17}
\end{equation*}
$$

holds true. From this estimate and (13) we derive

$$
\begin{equation*}
\left\|u_{\mu, \varepsilon}-u^{0}\right\|_{L^{2}(I)} \leqslant C\left(\varepsilon^{1 / 2}+\mu\right)\|f\|_{L^{2}(I)} \tag{18}
\end{equation*}
$$

Thus, to complete the proof of Theorem 2.1, it is sufficient to improve estimate (18) to get (4).
From inequalities (17) and (10), it follows that, for $\varkappa \neq \varkappa_{0}$ and sufficiently small $\mu$, the solutions to (1) satisfy the uniform in $\mu$ and $\varepsilon$ estimate $\left\|u_{\mu, \varepsilon}\right\|_{L^{2}(I)} \leqslant C\|f\|_{L^{2}(I)}$. On the other hand, this inequality leads to the estimate

$$
\begin{equation*}
\|w\|_{L^{2}(I)} \leqslant C\left(\|F\|_{L^{2}(I)}+|\widetilde{A}|+|\widetilde{B}|\right) \tag{19}
\end{equation*}
$$

for solutions to the boundary value problem

$$
\mathcal{L}_{\mu, \varepsilon} w=F, \quad x \in I, \quad l_{a} w=\widetilde{A} \quad l_{b} w=\widetilde{B}
$$

outside a vicinity of the point $\varkappa_{0}$ for sufficiently small $\mu$.

## 5. Derivation of estimate (4)

The construction of a formal approximation $Z_{\mu}(x, \varepsilon)$ of solutions to the boundary value problem (1) is based on the well-known method of matched asymptotic expansions [9,12]. Hence, we omit trivial explanations on getting the structure of $Z_{\mu}(x, \varepsilon)$.

Denote

$$
\begin{align*}
v_{1, \mu}(\xi) & =\varkappa u^{\mu}(0)\left(\xi \int_{-\infty}^{\xi} Q(\tau) \mathrm{d} \tau-\int_{-\infty}^{\xi} \tau Q(\tau) \mathrm{d} \tau\right)+\left(u^{\mu}\right)^{\prime}(-0) \xi \\
\widetilde{v}_{1, \mu}(\xi) & =v_{1, \mu}(\xi)-\left(u^{\mu}\right)^{\prime}(\mp 0) \xi, \quad \mp \xi>0 \\
Z_{\mu}(x, \varepsilon) & =u^{\mu}(x)+\varepsilon \widetilde{v}_{1, \mu}\left(\frac{x}{\varepsilon}\right) \tag{20}
\end{align*}
$$

Then, $Z_{\mu} \in H^{2}(I)$ due to (9) and since $v_{1, \mu}^{\prime \prime}(\xi)=\varkappa u^{\mu}(0) Q(\xi)$, then due to (7) we have

$$
\begin{align*}
\mathcal{L}_{\mu, \varepsilon} Z_{\mu}(x, \varepsilon) & =f(x)+\mathcal{I}_{\mu}^{1}(x, \varepsilon)+\mathcal{I}_{\mu}^{2}(x, \varepsilon)+\mathcal{J}_{\mu}(x, \varepsilon), \quad x \in I \\
l_{a} Z_{\mu} & =0, \quad l_{b} Z_{\mu}=-\varepsilon u^{\mu}(0) h_{b} 火 \int_{-\infty}^{\infty} \tau Q(\tau) \mathrm{d} \tau \tag{21}
\end{align*}
$$

where

$$
\begin{aligned}
& \mathcal{I}_{\mu}^{1}(x, \varepsilon)=\varepsilon^{-1} \varkappa Q\left(\frac{x}{\varepsilon}\right)\left(u^{\mu}(x)-u^{\mu}(0)-\left(u^{\mu}\right)^{\prime}(\mp 0) x\right), \quad \mp x>0 \\
& \mathcal{I}_{\mu}^{2}(x, \varepsilon)=\varepsilon q\left(x, \frac{x}{\mu}\right) \widetilde{v}_{1, \mu}\left(\frac{x}{\varepsilon}\right), \quad \mathcal{J}_{\mu}(x, \varepsilon)=\varkappa Q\left(\frac{x}{\varepsilon}\right) v_{1, \mu}\left(\frac{x}{\varepsilon}\right)
\end{aligned}
$$

From (20), (10) and the embedding theorems it follows that

$$
\begin{equation*}
\left\|\mathcal{I}_{\mu}^{2}\right\|_{L^{2}(I)}+\left|l_{b} Z_{\mu}\right| \leqslant \varepsilon C\|f\|_{L^{2}(I)} \tag{22}
\end{equation*}
$$

If $f \in C(\bar{I})$, then, using for $\mathcal{I}_{\mu}^{1}(x, \varepsilon)$ the formula for the remainder term of the Taylor series in the integral form and the equation from (7), we get

$$
\mathcal{I}_{\mu}^{1}(x, \varepsilon)=\varepsilon^{-1} \varkappa Q\left(\frac{x}{\varepsilon}\right) \int_{0}^{x}(x-t)\left(q\left(t, \frac{t}{\mu}\right) u^{\mu}(t)-f(t)\right) \mathrm{d} t
$$

Then, keeping in mind the Cauchy-Bunjakovski-Schwarz inequality and estimate (10), we derive

$$
\begin{equation*}
\left\|\mathcal{I}_{\mu}^{1}\right\|_{L^{2}(I)} \leqslant \varepsilon C\|f\|_{L^{2}(I)} \tag{23}
\end{equation*}
$$

uniformly in $\mu$. Since the set $C(\bar{I})$ is dense in $L^{2}(I)$, then this estimate is valid for any $f \in L^{2}(I)$.
However, it easy to see that $\left\|J_{\mu}\right\|_{L^{2}(I)}=O\left(\varepsilon^{1 / 2}\right)$. Consequently we need a corrector term for the function $Z_{\mu}(x, \varepsilon)$. Let us define that

$$
\begin{align*}
\widetilde{v}_{2, \mu}(\xi) & :=\varkappa\left(\xi \int_{-\infty}^{\xi} v_{1, \mu}(\tau) Q(\tau) \mathrm{d} \tau-\int_{-\infty}^{\xi} \tau v_{1, \mu}(\tau) Q(\tau) \mathrm{d} \tau\right)  \tag{24}\\
W_{\mu}(x, \varepsilon) & :=Z_{\mu}(x, \varepsilon)+\varepsilon^{2} \widetilde{v}_{2, \mu}\left(\frac{x}{\varepsilon}\right) \tag{25}
\end{align*}
$$

Then $W_{\mu} \in H^{2}(I)$ and since $v_{2, \mu}^{\prime \prime}(\xi)=v_{1, \mu}(\xi) Q(\xi)$, then by means of (21) we have

$$
\begin{align*}
& \mathcal{L}_{\mu, \varepsilon} W_{\mu}(x, \varepsilon)=f(x)+\mathcal{I}_{\mu}^{1}(x, \varepsilon)+\mathcal{I}_{\mu}^{2}(x, \varepsilon)+\mathcal{I}_{\mu}^{3}(x, \varepsilon), \quad x \in I \\
& l_{a} W_{\mu}=0, \quad l_{b} W_{\mu}=l_{b} Z_{\mu}+\varepsilon x \int_{-\infty}^{\infty} v_{1, \mu}(\tau) Q(\tau)\left(h_{b}(b-\varepsilon \tau)+H_{b}\right) \mathrm{d} \tau \tag{26}
\end{align*}
$$

where

$$
\begin{equation*}
\mathcal{I}_{\mu}^{3}(x, \varepsilon)=\varepsilon \varkappa Q\left(\frac{x}{\varepsilon}\right) \widetilde{v}_{2, \mu}\left(\frac{x}{\varepsilon}\right)+\varepsilon^{2} q\left(x, \frac{x}{\mu}\right) \widetilde{v}_{2, \mu}\left(\frac{x}{\varepsilon}\right) \tag{27}
\end{equation*}
$$

From (27), (24), (20), (22), (23) and (10) we get

$$
\left\|\mathcal{I}_{\mu}^{1}\right\|_{L^{2}(I)}+\left\|\mathcal{I}_{\mu}^{2}\right\|_{L^{2}(I)}+\left\|\mathcal{I}_{\mu}^{3}\right\|_{L^{2}(I)}+\left|l_{b} W_{\mu}\right| \leqslant \varepsilon C\|f\|_{L^{2}(I)}
$$

Then, due to the problems (1) and (26) we get

$$
\left\|\mathcal{L}_{\mu, \varepsilon}\left(u_{\mu, \varepsilon}-W_{\mu}\right)\right\|_{L^{2}(I)}+\left|l_{a}\left(u_{\mu, \varepsilon}-W_{\mu}\right)\right|+\left|l_{b}\left(u_{\mu, \varepsilon}-W_{\mu}\right)\right| \leqslant \varepsilon C\|f\|_{L^{2}(I)}
$$

uniformly in $\mu$; and using (19), we derive

$$
\begin{equation*}
\left\|u_{\mu, \varepsilon}-W_{\mu}\right\|_{L^{2}(I)} \leqslant \varepsilon C\|f\|_{L^{2}(I)} \tag{28}
\end{equation*}
$$

From (20) and (10), it follows that

$$
\begin{equation*}
\left\|u^{\mu}-Z_{\mu}\right\|_{L^{2}(I)} \leqslant \varepsilon C\|f\|_{L^{2}(I)} \tag{29}
\end{equation*}
$$

In an analogous way, using (25), (24), (20) and (10), we get

$$
\begin{equation*}
\left\|W_{\mu}-Z_{\mu}\right\|_{L^{2}(I)} \leqslant \varepsilon C\|f\|_{L^{2}(I)} \tag{30}
\end{equation*}
$$

The inequalities (28), (29), (30) and (13) lead to estimate (4).

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## References

[1] N.S. Bakhvalov, G.P. Panasenko, Homogenization: Averaging Processes in Periodic Media, Kluwer, London, 1989.
[2] O.A. Oleinik, A.S. Shamaev, G.A. Yosifian, Mathematical Problems in Elasticity and Homogenization, North-Holland, Amsterdam, 1992.
[3] G.A. Chechkin, A.L. Piatnitski, A.S. Shamaev, Homogenization: Methods and Applications, Amer. Math. Soc., Providence, 2007.
[4] O.A. Oleinik, Homogenization problems in elasticity: spectrum of singularly perturbed operators, in: Non-Classical Continuum Mechanics, in: Lecture Note Series, vol. 122, 1987, pp. 188-205.
[5] Yu.D. Golovaty, S.A. Nazarov, O.A. Oleinik, T.S. Soboleva, Eigenoscillations of a string with an additional mass, Sib. Math. J. 29 (5) (1988) 744-760.
[6] V. Maz'ya, S. Nazarov, B. Plamenevsky, Asymptotic Theory of Elliptic Boundary Value Problems in Singularly Perturbed Domains, vols. 1, 2, Birkhäuser Verlag, Basel, Switzeland, 2000.
[7] E. Sánchez-Palencia, Perturbation of eigenvalues in thermo-elasticity and vibration of systems with concentrated masses, in: Trends and Applications of Pure Mathematics to Mechanics, in: Lecture Notes in Physics, Springer-Verlag, 1984, pp. 346-368.
[8] S. Albeverio, F. Gesztezy, R. Høegh-Krohn, H. Holden, On point interactions in one dimension, J. Oper. Theory 12 (1984) 101-126.
[9] A.M. Il'in, Matching of Asymptotic Expansions of Solutions of Boundary-Value Problems, American Mathematical Society, Providence, RI, USA, 1992.
[10] V.P. Mikhailov, Partial Differential Equations, Mir Publishers, Moscow, 1978.
[11] T. Kato, Perturbation Theory for Linear Operators, Springer-Verlag, Heidelberg, Germany, 1966.
[12] R.R. Gadyl'shin, I.Kh. Khusnullin, Perturbation of the Shrödinger operator by a narrow potential, Ufa Math. J. 3 (3) (2011) 54-64; translated from: Ufimsk. Mat. Zh. 3 (3) (2011) 55-66.


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