# On a singularly perturbed Steklov problem in a domain perforated along the boundary 

# Le problème Steklov singulièrement perturbé dans un domaine perforé le long de la frontière 

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#### Abstract

We study the asymptotic behavior of solutions and eigenelements to a boundary value problem for the Laplace equation in a domain perforated along part of the boundary. On the boundary of holes, we set the homogeneous Dirichlet boundary condition and the Steklov spectral condition on the mentioned part of the outer boundary of the domain. Assuming that the boundary microstructure is periodic, we construct the limit problem and prove the homogenization theorem. © 2015 Académie des sciences. Published by Elsevier Masson SAS. All rights reserved.


## Ré S U M É

Nous étudions le comportement asymptotique des solutions et des éléments propres à un problème aux limites pour l'équation de Laplace dans un domaine perforé le long d'une partie de la frontière. Sur la frontière de trous, nous posons la condition de Dirichlet homogène et la condition spectrale de Steklov sur la part mentionnée de la frontière extérieure du domaine. En supposant que la microstructure de la frontière est périodique, nous construisons le problème aux limites et prouvons le théorème d'homogénéisation.
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Fig. 1. Structure of the domain $\Omega_{\varepsilon, a}$ and the cell of periodicity.

## 1. Introduction

Steklov's spectral problem is well known and it has been studied for several decades (see, for instance [1-5]). There are numerous papers dealing with homogenization of problems in domains perforated along the boundary (e.g., see [6-15]).

In this paper, we study spectral problems with Steklov-type boundary conditions in 2D and 3D domains periodically perforated along part of the boundary, which are of interest in connection with the applications. Under the assumption that the ratio between the diameter of the cavities and the distance between them tends to zero, we present different cases of the limiting behavior of eigenpairs.

## 2. Setting of the problem and main results

Let $\Omega$ be a bounded domain in $\mathbb{R}^{d}, d=2,3$, situated in the semi-plane $x_{2}>0$ for $d=2$ and in the semi-space $x_{3}>0$ for $d=3$. Its boundary $\Gamma$ consists of two parts: $\Gamma=\Gamma_{1} \cup \Gamma_{2}$, where $\Gamma_{1}$ is the segment $[0,1]$ in the axis $x_{2}=0$ for $d=2$ and square $[0,1]^{2}$ in the plane $x_{3}=0$ for $d=3$. For $d=2$, the part $\Gamma_{2}$ is infinitely differentiable and in a neighborhood of the points $(0,0)$ and $(1,0)$ coincides with the lines $x_{1}=0$ and $x_{1}=1$, respectively. For $d=3$, the part $\Gamma_{2}$ coincides with the lateral faces of the cube $[0,1]^{3}$ in a small neighborhood of the plane $x_{3}=0$ and, in addition, it is infinitely differentiable everywhere, except the vertical edges.

Then assume that $B$ is an arbitrary bounded domain with Lipschitz boundary. Denote $B_{a}=\left\{x:\left(a^{-1}\left(x_{1}-b_{1}\right)\right.\right.$, $\left.\left.a^{-1}\left(x_{2}-c\right)\right) \in B\right\}$ for $d=2, B_{a}=\left\{x:\left(a^{-1}\left(x_{1}-b_{1}\right), a^{-1}\left(x_{2}-b_{2}\right), a^{-1}\left(x_{3}-c\right)\right) \in B\right\}, j=1,2$, for $d=3$, where $0<b_{j}<1$, $c>0$ are arbitrary fixed numbers, $a$ is a sufficiently small positive parameter, such that $\overline{B_{a}}$ lies in the semi-strip $\Pi=(0,1) \times(0, \infty)$ for $d=2$ and in the semi-infinite parallelepiped $\Pi=(0,1)^{2} \times(0, \infty)$ for $d=3$.

Denote $B_{\varepsilon, a}^{\mathbf{k}}=\left\{x:\left(\varepsilon^{-1} x_{1}-\mathbf{k}, \varepsilon^{-1} x_{2}\right) \in B_{a}\right\}, \mathbf{k} \in \mathbb{Z}$ for $d=2, B_{\varepsilon, a}^{\mathbf{k}}=\left\{x:\left(\varepsilon^{-1} x_{1}-k_{1}, \varepsilon^{-1} x_{2}-k_{2}, \varepsilon^{-1} x_{3}\right) \in B_{a}\right\}, \mathbf{k}=\left(k_{1}, k_{2}\right)$, $k_{j} \in \mathbb{Z}$ for $d=3, B_{\varepsilon, a}=\bigcup_{\mathbf{k}} B_{\varepsilon, a}^{\mathbf{k}}, \Gamma_{\varepsilon, a}=\partial B_{\varepsilon, a}$. Hereafter $\varepsilon$ is a small positive parameter, $\varepsilon=\varepsilon_{N}=\frac{1}{N}$, where $N \gg 1$ is a natural number. Define the domain $\Omega_{\varepsilon, a}$ as $\Omega \backslash \overline{B_{\varepsilon, a}}$ (see Fig. 1).

Consider the boundary value problems

$$
\begin{align*}
-\Delta U_{\varepsilon, a}=0 \quad \text { in } \Omega_{\varepsilon, a}, & U_{\varepsilon, a}
\end{align*}=0 \quad \text { on } \Gamma_{\varepsilon, a}, ~ \begin{array}{ll}
\frac{\partial U_{\varepsilon, a}}{\partial v}=\lambda U_{\varepsilon, a}+f & \text { on } \Gamma_{1},
\end{array} \frac{\partial U_{\varepsilon, a}}{\partial v}=0 \quad \text { on } \Gamma_{2} .
$$

where $\nu$ is an outer normal, $\lambda \in \mathbb{R}$, and the respective spectral problems

$$
\begin{gather*}
-\Delta u_{\varepsilon, a}=0 \quad \text { in } \Omega_{\varepsilon, a}, \quad u_{\varepsilon, a}=0 \quad \text { on } \Gamma_{\varepsilon, a} \\
\frac{\partial u_{\varepsilon, a}}{\partial v}=\lambda_{\varepsilon, a} u_{\varepsilon, a} \quad \text { on } \Gamma_{1}, \quad \frac{\partial u_{\varepsilon, a}}{\partial v}=0 \quad \text { on } \Gamma_{2}  \tag{3}\\
-\Delta u_{0}=0 \text { in } \Omega, \quad \frac{\partial u_{0}}{\partial v}+\sigma_{d} C_{d}(B) A u_{0}=\lambda_{0} u_{0} \text { on } \Gamma_{1}, \quad \frac{\partial u_{0}}{\partial v}=0 \text { on } \Gamma_{2} \tag{4}
\end{gather*}
$$

Hereafter $\sigma_{2}=2 \pi, \sigma_{3}=4 \pi, C_{2}(B)=1$ and $C_{3}(B)>0$ is the capacity of the domain $B$. In particular, $C_{3}(B)=1$, if $B$ is a unit ball.

We prove the following statements.

Theorem 2.1. Suppose that

$$
\begin{equation*}
-\frac{1}{\varepsilon \ln a} \longrightarrow A \neq \infty \quad \text { for } d=2, \quad \frac{a}{\varepsilon} \longrightarrow A \neq \infty \quad \text { for } d=3 \tag{5}
\end{equation*}
$$

$f \in L_{2}\left(\Gamma_{1}\right)$ and $\lambda$ is not an eigenvalue of the problem (4).
Then:

1) the boundary value problem (1) has a unique solution in $W_{2}^{1}\left(\Omega_{\varepsilon, a}\right)$ for any sufficiently small $\varepsilon$, and moreover the following uniform in $\varepsilon$ estimate:

$$
\begin{equation*}
\left\|U_{\varepsilon, a}\right\|_{W_{2}^{1}(\Omega)} \leqslant C\|f\|_{L_{2}\left(\Gamma_{1}\right)} \tag{6}
\end{equation*}
$$

holds true, where the function $U_{\varepsilon, a}$ is extended in $\overline{B_{\varepsilon, a}}$ by zero;
2) for the solution to problem (1) the following strong convergence

$$
\begin{equation*}
U_{\varepsilon, a} \rightarrow U_{\varepsilon \rightarrow 0} \quad \text { in } W_{2}^{1}(\Omega) \tag{7}
\end{equation*}
$$

takes place, if $A=0$, and the weak convergence

$$
\begin{equation*}
U_{\varepsilon, a} \underset{\varepsilon \rightarrow 0}{\rightharpoonup} U_{0} \quad \text { in } W_{2}^{1}(\Omega) \tag{8}
\end{equation*}
$$

holds true, if $A \neq 0$, where $U_{0}$ is a solution to the homogenized (limit) problem (2).
Theorem 2.2. Suppose that the condition (5) holds, and the multiplicity of the eigenvalue $\lambda_{0}$ to the problem (4) equals to $n$. Then there exist $n$ eigenfunctions $\lambda_{\varepsilon, a}^{(l)}$ of problem (3), $l=\overline{1, n}$ (with respect to their multiplicities) converging to $\lambda_{0}$ as $\varepsilon \rightarrow 0$.

Let $u_{\varepsilon, a}^{(l)}$ be orthonormalized in $L_{2}\left(\Gamma_{1}\right)$ eigenfunctions of problem (3), corresponding to $\lambda_{\varepsilon, a}^{(l)}$. Then from the sequence $\left\{\varepsilon_{k}=\frac{1}{k}\right\}_{k=1}^{\infty}$ and any sequence $\left\{a_{k}\right\}_{k=1}^{\infty}$ (for which the convergence (5) takes place) one can choose subsequences $\left\{\varepsilon_{k^{\prime}}\right\},\left\{a_{k^{\prime}}\right\}$ such that the convergence

$$
u_{\varepsilon, a}^{(l)} \underset{\varepsilon \rightarrow 0}{\longrightarrow} u_{*}^{(l)} \quad \text { in } W_{2}^{1}(\Omega)
$$

holds, if $A=0$, and weak convergence

$$
u_{\varepsilon, a}^{(l)} \underset{\varepsilon \rightarrow 0}{\rightharpoonup} u_{*}^{(l)} \quad \text { in } W_{2}^{1}(\Omega)
$$

holds, if $A \neq 0$, where the functions $u_{\varepsilon, a}^{(l)}$ are extended in $\overline{B_{\varepsilon, a}}$ by zero, and $u_{*}^{(l)}$ are orthonormalized in $L_{2}\left(\Gamma_{1}\right)$ eigenfunctions of problem (4), corresponding to $\lambda_{0}$ (they depend on the choice of the sequence $\left\{a_{k}\right\}_{k=1}^{\infty}$ and the subsequence).

The solution to the problems (1), (2), (3), and (4) is understood in a weak sense (in the sense of the integral identity, see, for instance, [16, Chapter IV]).

## 3. Construction of model functions in the semi-strip and in the semi-infinite parallelepiped

For $d=2$ we denote by $\mathbb{R}_{+}^{2}$ the semi-plane $x_{2} \geqslant 0$, and by $x_{0}^{(\mathbf{k})}$ the points with coordinates $x_{1}=b_{1}+\mathbf{k}, x_{2}=c$. Here $\mathbf{k} \in \mathbb{Z}$. For $d=3$ we denote by $\mathbb{R}_{+}^{3}$ the semi-space $x_{3} \geqslant 0$, and by $x_{0}^{(\mathbf{k})}$ the points with coordinates $x_{j}=b_{j}+k_{j}, j=1,2$, $x_{3}=c$. Here $\mathbf{k}=\left(k_{1}, k_{2}\right)$ and $k_{j} \in \mathbb{Z}$. Assume that $x_{0}:=x_{0}^{(\mathbf{0})}, y:=x-x_{0}$ and $\Sigma:=\left\{x: x_{1} \in(0,1), x_{2}=0\right\}$ for $d=2, \Sigma:=\{x$ : $\left.x_{1}, x_{2} \in(0,1), x_{3}=0\right\}$ for $d=3$.

Define $G_{2}(t):=\ln t, G_{3}(t):=-t^{-1}$.
It is easy to prove the following lemma.
Lemma 3.1. There exists a function $g_{d} \in C^{\infty}\left(\mathbb{R}_{+}^{d} \backslash \bigcup_{\mathbf{k}}\left\{x_{0}^{(\mathbf{k})}\right\}\right)$, 1-periodic in $x_{1}$ for $d=2$ and in $x_{1}, x_{2}$ for $d=3$, which satisfies the problem

$$
\begin{cases}\Delta g_{d}=0 & \text { if } x \in \Pi \backslash\left\{x_{0}\right\} \\ \frac{\partial g_{d}}{\partial v}=\sigma_{d} & \text { if } x \in \Sigma\end{cases}
$$

has the differentiable asymptotics

$$
g_{d}(x)=O\left(\mathrm{e}^{-2 \pi x_{d}}\right), \quad x_{d} \rightarrow+\infty
$$

and in a neighborhood of $x_{0}$ has the representation

$$
g_{d}(x)=G_{d}(|y|)+g_{d}^{(1)}(x)
$$

where $g_{d}^{(1)}(x)$ is an infinitely smooth function in the neighborhood of this point.
Corollary 3.2. The differentiable asymptotics

$$
g_{d}(x)=G_{d}(|y|)+c_{\Pi, d}+P_{1}^{\Pi, d}(y)+O\left(|y|^{2}\right), \quad y \rightarrow 0
$$

holds, where $c_{\Pi, d}$ is a constant and $P_{1}^{\Pi, d}(y)$ is a homogeneous polynomial of the first order.
Using [17, § 5.8] we prove the following lemma.
Lemma 3.3. There exist functions $V_{0}^{(d)}, V_{1}^{(d)} \in C^{\infty}\left(\mathbb{R}^{d} \backslash B\right)$, being solutions to the problems

$$
\Delta V_{i}^{(d)}=0, \quad x \in \mathbb{R}^{d} \backslash \bar{B}, \quad V_{i}^{(d)}=0, \quad x \in \partial B
$$

and having differentiable asymptotics

$$
\begin{array}{lc}
V_{0}^{(2)}(x)=\ln |x|+c_{B}+O\left(|x|^{-1}\right), & V_{0}^{(3)}(x)=1-C_{3}(B)|x|^{-1}+P_{1}^{B, 3}(x)|x|^{-3}+O\left(|x|^{-3}\right) \\
V_{1}^{(2)}(x)=P_{1}^{\Pi, 2}(x)+\tilde{c}+O\left(|x|^{-1}\right), & V_{1}^{(3)}(x)=P_{1}^{\Pi, 3}(x)-C_{3}(P, B)|x|^{-1}+O\left(|x|^{-2}\right)
\end{array}
$$

as $|x| \rightarrow \infty$, where $P_{1}^{B, 3}(y)$ is a homogeneous polynomial of the first order.
Lemma 3.4. Let $d=3$. Then there exists a 1-periodic in $x_{1}$, $x_{2}$ function $\widehat{g}_{3} \in C^{\infty}\left(\mathbb{R}_{+}^{3} \backslash \bigcup_{\mathbf{k}}\left\{x_{0}^{(\mathbf{k})}\right\}\right)$, which satisfies the problem

$$
\left\{\begin{array}{l}
\Delta \widehat{g}_{3}=0 \quad \text { if } x \in \Pi \backslash\left\{x_{0}\right\} \\
\frac{\partial \widehat{g}_{3}}{\partial v}=0 \quad \text { if } x \in \Sigma
\end{array}\right.
$$

has the differentiable asymptotics

$$
\widehat{g}_{3}(x)=O\left(\mathrm{e}^{-2 \pi x_{3}}\right), \quad x_{3} \rightarrow+\infty
$$

and in a neighborhood of $x_{0}$ has the representation

$$
\widehat{g}_{3}(x)=P_{1}^{B, 3}(y)|y|^{-3}+\widehat{g}_{3}^{(1)}(x)
$$

where $\widehat{g}_{3}^{(1)}(x)$ is an infinitely smooth function in the neighborhood of this point including this point.
Corollary 3.5. The differentiable asymptotics

$$
\begin{equation*}
\widehat{g}_{3}(x)=P_{1}^{B, 3}(y)|y|^{-3}+\widehat{c}_{B, \Pi}+O(|y|), \quad y \rightarrow 0 \tag{9}
\end{equation*}
$$

holds.

Denote $\Pi_{a}=\Pi \backslash \overline{B_{a}}$ (see Figure) and define in $\overline{\Pi_{a}}$ the following function:

$$
\begin{array}{rlr}
W_{a}(x):= & \left(1-\chi\left(\frac{|y|}{a^{\beta}}\right)\right)\left(1-\frac{1}{\ln a}\left(g_{2}(x)+c_{B}-c_{\Pi, 2}\right)\right) \\
& -\frac{1}{\ln a} \chi\left(\frac{|y|}{a^{\beta}}\right)\left(V_{0}^{(2)}\left(\frac{y}{a}\right)+a V_{1}^{(2)}\left(\frac{y}{a}\right)\right) & \text { for } d=2 \\
W_{a}(x):= & \left(1-\chi\left(\frac{|y|}{a^{\beta}}\right)\right)\left(1+a C_{3}(B)\left(g_{3}(x)-c_{\Pi, 3}\right)+a^{2}\left(C_{3}(P, B) g_{3}(x)+\widehat{g}_{3}(x)\right)\right)  \tag{10}\\
& +\chi\left(\frac{|y|}{a^{\beta}}\right)\left(V_{0}^{(3)}\left(\frac{y}{a}\right)+a V_{1}^{(3)}\left(\frac{y}{a}\right)\right) & \text { for } d=3
\end{array}
$$

where $\beta \in(0,1)$.
Now, denote $B_{a}^{\mathbf{k}}=\left\{x:\left(x_{1}-\mathbf{k}, x_{2}\right) \in B_{a}\right\}, \mathbf{k} \in \mathbb{Z}$ for $d=2, B_{a}^{\mathbf{k}}=\left\{x:\left(x_{1}-k_{1}, x_{2}-k_{2}, x_{3}\right) \in B_{a}\right\}, \mathbf{k}=\left(k_{1}, k_{2}\right), k_{j} \in \mathbb{Z}$ for $d=3$, $B^{a}=\bigcup_{\mathbf{k}} B_{a}^{\mathbf{k}}$ and extend the function $W_{a}(x)$ 1-periodically in $x_{1}$ for $d=2$ and in $x_{1}, x_{2}$ for $d=3$, keeping the same notation $W_{a}(x)$.

Using Lemmas 3.1 and 3.3, we prove the following theorem.

Theorem 3.6. The function $W_{a}(x) \in C^{\infty}\left(\mathbb{R}_{+}^{d} \backslash B^{a}\right)$ is 1-periodic in $x_{1}$ for $d=2$, and in $x_{1}, x_{2}$ for $d=3$, has the differentiable asymptotics

$$
\begin{array}{ll}
W_{a}(x)=1-\frac{1}{\ln a}\left(c_{B}-c_{\Pi, 2}+O\left(\mathrm{e}^{-2 \pi x_{2}}\right)\right) & \text { as } x_{2} \rightarrow \infty \text { for } d=2 \\
W_{a}(x)=1-a\left(C_{3}(B) c_{\Pi, 3}+O\left(\mathrm{e}^{-2 \pi x_{3}}\right)\right) & \text { as } x_{3} \rightarrow \infty \text { for } d=3
\end{array}
$$

uniform in $a$, and satisfies the problem

$$
\left\{\begin{array}{l}
\Delta W_{a}=F_{a} \quad \text { if } x \in \Pi_{a} \\
\frac{\partial W_{a}}{\partial v}=-\frac{2 \pi}{\ln a} \quad \text { for } d=2, \quad \frac{\partial W_{a}}{\partial v}=a 4 \pi C_{3}(B) \quad \text { for } d=3, \quad \text { if } x \in \Sigma \\
W_{a}=0 \quad \text { if } x \in \partial B_{a}
\end{array}\right.
$$

where $F_{a} \in C_{0}^{\infty}\left(\Pi_{a}\right)$.
Moreover,

$$
\begin{gathered}
\left\|1-W_{a}\right\|_{L_{2}(\Sigma)}=O\left(\frac{1}{|\ln a|}\right) \quad \text { for } d=2, \quad\left\|1-W_{a}\right\|_{L_{2}(\Sigma)}=0(a) \quad \text { for } d=3 \\
\left\|1-\left(W_{a}+\frac{1}{\ln a}\left(c_{B}-c_{\Pi, 2}\right)\right)\right\|_{L_{2}\left(\Pi_{a}\right)}=O\left(\frac{1}{|\ln a|}\right) \quad \text { for } d=2 \\
\left\|1-\left(W_{a}+a C_{3}(B) c_{\Pi, 3}\right)\right\|_{L_{2}\left(\Pi_{a}\right)}=0\left(a+a^{\frac{3 \beta}{2}}\right) \quad \text { for } d=3 \\
\left\|F_{a}\right\|_{L_{2}\left(\Pi_{a}\right)}=0\left(\frac{1}{|\ln a|}\left(a^{\beta}+a^{1-2 \beta}\right)\right) \quad \text { for } d=2 \\
\left\|F_{a}\right\|_{L_{2}\left(\Pi_{a}\right)}=0\left(a^{1+\frac{3}{2} \beta}+a^{2-\frac{1}{2} \beta}+a^{3-\frac{7}{2} \beta}\right) \quad \text { for } d=3
\end{gathered}
$$

## 4. Proof of Theorem 2.1

Before proving Theorem 2.1, we prove some auxiliary propositions. Using Theorem 3.6, we prove the following lemma.
Lemma 4.1. There exists a function $W_{\varepsilon, a}(x)$ from $W_{2}^{1}\left(\Omega_{\varepsilon, a} ; \Gamma_{\varepsilon, a}\right) \cap W_{2}^{2}\left(\Omega_{\varepsilon, a}\right)$, such that the relations

$$
\begin{aligned}
&\left.\frac{\partial W_{\varepsilon, a}}{\partial v}\right|_{\Gamma_{1}}=-\frac{2 \pi}{\varepsilon \ln a} \quad \text { ford } d=2,\left.\quad \frac{\partial W_{\varepsilon, a}}{\partial v}\right|_{\Gamma_{1}}=\frac{4 \pi C_{3}(B) a}{\varepsilon} \text { for } d=3 \\
&\left\|1-W_{\varepsilon, a}\right\|_{L_{2}(\Omega)} \underset{\varepsilon \rightarrow 0}{\longrightarrow} 0 \\
&\left\|\Delta W_{\varepsilon, a}\right\|_{L_{2}\left(\Omega_{\varepsilon, a}\right)}^{\longrightarrow} 0 \\
&\left\|1-W_{\varepsilon, a}\right\|_{L_{2}(\Gamma)}^{\underset{\varepsilon \rightarrow 0}{\longrightarrow} 0} 0 \\
&\left\|\frac{\partial W_{\varepsilon, a}}{\partial v}\right\|_{L_{2}\left(\Gamma_{2}\right)} \xrightarrow[\varepsilon \rightarrow 0]{\longrightarrow} 0
\end{aligned}
$$

hold if (5) is true.
Next lemma follows from Lemma 4.1.
Lemma 4.2. Let condition (5) hold. Assume also that the function $U_{\varepsilon, a} \in W_{2}^{1}\left(\Omega_{\varepsilon, a} ; \Gamma_{\varepsilon, a}\right)$ converges weakly

$$
U_{\varepsilon, a} \underset{\varepsilon \rightarrow 0}{\rightharpoonup} U^{*} \quad \text { in } W_{2}^{1}(\Omega)
$$

Then for any functions $v \in C^{\infty}(\bar{\Omega})$ the convergences

$$
\begin{aligned}
\int_{\Omega_{\varepsilon, a}} U_{\varepsilon, a}\left(v W_{\varepsilon, a}\right) \mathrm{d} x & \rightarrow \int_{\Omega} U^{*} v \mathrm{~d} x \\
\int_{\Omega_{\varepsilon, a}} \nabla U_{\varepsilon, a} \nabla\left(v W_{\varepsilon, a}\right), \mathrm{d} x & \rightarrow \int_{\Omega} \nabla U^{*} \nabla v \mathrm{~d} x+\sigma_{d} C_{d}(B) A \int_{\Gamma_{1}} U^{*} v \mathrm{~d} x_{1}
\end{aligned}
$$

take place.

Proof of Theorem 2.1. Since the Fredholm alternative for problem (1) holds (see, for instance, [16, Chapter II, § 3]), then it is sufficient to show the estimate (6) to prove 1).

The standard norm in $W_{2}^{1}(\Omega)$ is equivalent to the norm in $\|u\|_{H^{1}(\Omega)}$, generated by the following scalar product:

$$
(u, v)_{H^{1}(\Omega)}=\int_{\Omega} \nabla u \nabla v \mathrm{~d} x+\int_{\Gamma_{1}} u v \mathrm{~d} s
$$

(see, for instance, [16, Ch. III, §5.6]).
Using the identity of the problems (1), (2), and applying the Cauchy inequality, we get, for any fixed $\lambda \in \mathbb{R}$, the a priori estimate, uniform in $\varepsilon$ and $a$, for the solution to problem (1), of the form

$$
\begin{equation*}
\left\|U_{\varepsilon, a}\right\|_{H^{1}(\Omega)} \leq C\left(\left\|U_{\varepsilon, a}\right\|_{L_{2}\left(\Gamma_{1}\right)}+\|f\|_{L_{2}\left(\Gamma_{1}\right)}\right) \tag{11}
\end{equation*}
$$

If $U_{\varepsilon_{m(k)}, a_{k}}=0$ on $\Gamma_{1}$, then estimate (6) follows from (11).
Then, applying Lemma 4.2, we prove the statement 2) and complete the proof.

In an analogous way, using Lemma 4.2, one can derive the following assertion.
Lemma 4.3. Let the condition (5) hold, assume also that $\lambda$ is not an eigenvalue of problem (4), $U_{\varepsilon, a}$ is the solution to problem (1) for $f=f_{\varepsilon, a}, U_{0}$ is the solution to problem (2) for $f=f_{0}$ and the weak convergence:

$$
\begin{equation*}
f_{\varepsilon, a} \underset{\varepsilon \rightarrow 0}{\longrightarrow} f_{0} \quad \text { in } L_{2}\left(\Gamma_{1}\right) \tag{12}
\end{equation*}
$$

holds.
Then the convergence (7) and (8) take place.
Obviously the following proposition holds true.

Lemma 4.4. Suppose that $\lambda$ is not an eigenvalue of the problem (4), $U^{\varepsilon, a}$ be is the solution to problem (2) for $f=f_{\varepsilon, a}, U_{0}$ is the solution to problem (2) for $f=f_{0}$ and the weak convergence (12) holds.

Then the weak convergence

$$
U^{\varepsilon, a} \underset{\varepsilon \rightarrow 0}{\longrightarrow} U_{0} \quad \text { in } W_{2}^{1}(\Omega)
$$

takes place.

## 5. Proof of Theorem 2.2

Denote by $\mathcal{P}_{\varepsilon, a}$ and $\mathcal{P}_{0}$ operators $\mathcal{P}_{\varepsilon, a}, \mathcal{P}_{0}: L_{2}\left(\Gamma_{1}\right) \rightarrow L_{2}\left(\Gamma_{1}\right)$, mapping $f$ to the traces on $\Gamma_{1}$ of solutions to boundary value problems (1) and (2), respectively, for $\lambda=-1$. It follows from the definition that these operators are compact, selfadjoint and positive.

Lemmas 4.3 and 4.4 lead to the following statements.
Lemma 5.1. If condition (5) and weak convergence (12) hold, then

$$
\mathcal{P}_{\varepsilon, a} f_{\varepsilon, a} \underset{\varepsilon \rightarrow 0}{\longrightarrow} \mathcal{P}_{0} f_{0}, \quad \mathcal{P}_{0} f_{\varepsilon, a} \underset{\varepsilon \rightarrow 0}{\longrightarrow} \mathcal{P}_{0} f_{0} \quad \text { in } \quad L_{2}\left(\Gamma_{1}\right)
$$

strongly.
Lemma 5.2. If condition (5) holds, then

$$
\mathcal{P}_{\varepsilon, a} \underset{\varepsilon \rightarrow 0}{\rightrightarrows} \mathcal{P}_{0}
$$

uniformly.
Denote by $\mathcal{L}_{\varepsilon, a}, \mathcal{L}_{0}: L_{2}\left(\Gamma_{1}\right) \rightarrow L_{2}\left(\Gamma_{1}\right)$ the operators inverse to $\mathcal{P}_{\varepsilon, a}, \mathcal{P}_{0}: L_{2}\left(\Gamma_{1}\right) \rightarrow L_{2}\left(\Gamma_{1}\right)$. From Lemma 5.2 and [18, Ch.IV, §3.4] (see also [19, Ch. 9, §4]), we easily derive the following proposition.

Lemma 5.3. Suppose that the condition (5) holds, and the multiplicity of the eigenvalue $\Lambda_{0}$ to the operator $\mathcal{L}_{0}$ equals to $n$. Then there exist $n$ eigenvalues $\Lambda_{\varepsilon, a}^{(l)}$ of the operator $\mathcal{L}_{\varepsilon, a}, l=\overline{1, n}$ (with respect to their multiplicities) converging to $\Lambda_{0}$ as $\varepsilon \rightarrow 0$.

Since obviously $\Lambda_{\varepsilon, a}^{(l)}=\lambda_{\varepsilon, a}^{(l)}+1, \Lambda_{0}=\lambda_{0}+1$, then the next assertion follows.
Lemma 5.4. Suppose that the condition (5) holds, and the multiplicity of the eigenvalue $\lambda_{0}$ to the problem (4) equals to $n$. Then there exist $n$ eigenvalues $\lambda_{\varepsilon, a}^{(l)}$ of problem (3), $l=\overline{1, n}$ (with respect to their multiplicities) converging to $\lambda_{0}$ as $\varepsilon \rightarrow 0$.

Using the trick from the proof of Theorem 2.1, we prove the next lemma.
Lemma 5.5. Assume that condition (5) holds, $\lambda_{\varepsilon, a}^{(l)}, l=\overline{1, n}$ are eigenvalues of problem (3), converging to $n$-multiple eigenvalue $\lambda_{0}$ of the limit problem (4) and $u_{\varepsilon, a}^{(l)}$ are the respective normalized in $L_{2}\left(\Gamma_{1}\right)$ eigenfunctions. Then, from a sequence $\varepsilon_{k}=\frac{1}{k}$ and any sequence $a_{k} \longrightarrow 0$ as $k \rightarrow \infty$, one can chose subsequences $\left\{\varepsilon_{k^{\prime}}\right\},\left\{a_{k^{\prime}}\right\}$, such that the strong convergence

$$
u_{\varepsilon, a}^{(l)} \underset{\varepsilon \rightarrow 0}{\longrightarrow} u_{*}^{(l)} \quad \text { in } W_{2}^{1}(\Omega)
$$

holds, if $A=0$ and a weak convergence

$$
u_{\varepsilon, a}^{(l)} \underset{\varepsilon \rightarrow 0}{\rightharpoonup} u_{*}^{(l)} \quad \text { in } W_{2}^{1}(\Omega)
$$

holds if $A \neq 0$, where $u_{*}^{(l)}$ are orthonormalized in $L_{2}(\Omega)$ eigenfunctions of the limit problem (4), corresponding to $\lambda_{0}$ (which in general depend on the choice of the sequence $a_{k} \longrightarrow 0$ as $k \rightarrow \infty$, and the subsequence).

Proof of Theorem 2.2. The proof lies on Lemmas 5.4 and 5.5.

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## References

[1] W. Stekloff, Sur les problèmes fondamentaux de la physique mathématique, Ann. Sci. Éc. Norm. Super. 19 (3) (1902) 191-259, 455-490 (in French).
[2] M.E. Pérez, On periodic Steklov type eigenvalue problems on half-bands and the spectral homogenization problem, Discrete Contin. Dyn. Syst., Ser. B 7 (4) (2007) 859-883.
[3] S.A. Nazarov, J. Taskinen, On the spectrum of the Steklov problem in a domain with a peak, Vestn. St. Petersbg. Univ., Ser. 11 (2008) 56-65 (English translation: Vestn. St. Petersbg. Univ., Math. 41 (1) (2008) 45-52).
[4] A.G. Chechkina, Convergence of solutions and eigenelements of Steklov type boundary value problems with boundary conditions of rapidly varying type, Probl. Mat. Anal. 42 (2009) 129-143 (English translation: J. Math. Sci. (N.Y.) 162 (3) (2009) 443-458).
[5] V.A. Sadovnichii, A.G. Chechkina, On estimate of eigenfunctions to the Steklov-type problem with small parameter in case of limit spectrum degeneration, Ufim. Mat. Zh. 3 (3) (2011) 127-139 (English translation: Ufa Math. J. 3 (3) (2011) 122-134).
[6] E. Sánchez-Palencia, Boundary value problems in domains containing perforated walls, in: Nonlinear Partial Differential Equations and Their Applications, vol. 3, Paris, 1980-1981, pp. 309-325.
[7] A.G. Belyaev, Homogenization of mixed boundary value problem for the Poisson equation in a domain perforated along the boundary, Usp. Mat. Nauk 45 (4) (1990) 123 (in Russian).
[8] M. Lobo, O.A. Oleinik, M.E. Pérez, T.A. Shaposhnikova, On homogenization of solutions of boundary value problems in domains perforated along manifolds, Ann. Sc. Norm. Super. Pisa, Cl. Sci. (4) 25 (3-4) (1997) 611-629.
[9] G.A. Chechkin, T.P. Chechkina, C. D' Apice, U. De Maio, Homogenization in domains randomly perforated along the boundary, Discrete Contin. Dyn. Syst., Ser. B 12 (4) (2009) 713-730.
[10] G.A. Chechkin, Yu.O. Koroleva, A. Meidell, L.-E. Persson, On the Friedrichs inequality in a domain perforated along the boundary. Homogenization procedure. Asymptotics in parabolic problems, Russ. J. Math. Phys. 16 (1) (2009) 1-16.
[11] R.R. Gadyl'shin, Yu.O. Koroleva, G.A. Chechkin, On the convergence of solutions and eigenelements of a boundary value problem in a domain perforated along the boundary, Differ. Uravn. 46 (5) (2010) 665-677 (English translation: Differ. Equ. 46 (5) (2010) 667-680).
[12] G.A. Chechkin, Yu.O. Koroleva, L.-E. Persson, P. Wall, A new weighted Friedrichs-type inequality for a perforated domain with a sharp constant, Eurasian Math. J. 2 (1) (2011) 81-103.
[13] R.R. Gadyl'shin, Yu.O. Koroleva, G.A. Chechkin, On the asymptotic behavior of a simple eigenvalue of a boundary value problem in a domain perforated along the boundary, Differ. Uravn. 47 (6) (2011) 819-828 (English translation: Differ. Equ. 47 (6) (2011) 822-831).
[14] R.R. Gadyl'shin, D.V. Kozhevnikov, G.A. Chechkin, On the spectral problem in a domain perforated along the boundary. Perturbation of multiple eigenvalue, Probl. Mat. Anal. 73 (2013) 31-45 (English translation: J. Math. Sci. (N.Y.) 196 (3) (2014) 276-292).
[15] R.R. Gadyl'shin, D.V. Kozhevnikov, Homogenization of the boundary value problem in a domain perforated along a part of the boundary, Probl. Mat. Anal. 75 (2014) 41-60 (English translation: J. Math. Sci. (N.Y.) 198 (6) (2014) 701-723).
[16] V.P. Mikhailov, Partial Differential Equations, Mir, Moscow, 1978.
[17] V.S. Vladimirov, Equazioni della Fisica Matematica (Equations of Mathematical Physics) (in Italian), translated from the fourth Russian edition by Ernest Kozlov Mir, Moscow, 1987.
[18] T. Kato, Perturbation Theory for Linear Operators, Springer, Heidelberg, Berlin, 1966.
[19] M.Sh.Birman, M.Z. Solomyak, Spectral Theory of Selfadjoint Operators in Hilbert Space, translated from Russian Mathematics and Its Applications (Soviet Series), D. Reidel Publishing Co., Dordrecht, The Netherlands, 1987.


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