



A new analytical formula to compute the step length of Padé approximants in the ANM: Application to buckling structures

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ABSTRACT

In this paper, we propose a new analytical formula to define the next branch in the Asymptotic Numerical Method (ANM) using the Padé approximants. The proposed formula is based on the computation of the relative error of two consecutive Padé approximants. This formula is obtained by developing the relative error with respect to the path parameter. An appropriate matrix formulation is adopted for the computation of this relative error. A comparison between the analytical formula proposed in this paper and the classical continuation Padé approximants using the step length computed numerically using dichotomy method is presented for examples of buckling structures.

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1. Introduction

The analysis of physical phenomena has always required the resolution of systems of nonlinear equations. The most powerful numerical algorithm enabling us to solve these systems of equations is the Newton–Raphson one [1–3]. This method is based on two stages; a prediction and a correction. The principle of these strategies is to follow the branch in a stepwise manner via a succession of linearizations and some iterations to achieve the equilibrium. The complete solution is sought by a continuation type point by point. This method has been applied in many nonlinear problems of structures [2,4–8].

An alternative to this method is the Asymptotic Numeric Method (ANM), which is a continuation procedure that allows one to obtain the solution to nonlinear problems in a succession of analytical branches [9,10]. This method has been successfully applied in several works [11–18]. Each branch is determined by a representation in vectorial series over a validity range $[0, a_{\max}^s]$ [10–12]. In these works, the step length a_{\max}^s is computed by an explicit formula [10]. In order to increase this validity range, these vectorial series are often replaced by an appropriate vectorial Padé approximant [19–21].

The continuation Padé approximants is used for the first time by Elhage-Hussein et al. in [13]. In this work, the step length a_{\max}^p is computed numerically by the dichotomy method. This computation of the step length a_{\max}^p is based on the relative error $E_M(a)$ of two consecutive Padé approximants of order $M - 1$, and M , respectively.

In this work, we propose for the first time to compute analytically the step length a_{\max}^p of these Padé approximants by using a development with respect to the path parameter a of the relative error $E_M(a)$. This development is done by using

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an appropriate matrix formulation of the relative error. The smallest positive root of this polynomial corresponds to a good approximation of the step length a_{\max}^p . In addition, we have made a comparison between the continuation Padé approximants used in the article [13] and the continuation Padé approximants proposed in this paper. The difference between these two continuations lies in the computation of the validity range a_{\max}^p . In the continuation Padé approximants of paper [13], the validity range a_{\max}^p is numerically computed by the dichotomy method and, in the proposed continuation Padé approximants, it is given directly by the new explicit analytical formula. A comparison between this analytical formula and that used in [13] is made on examples of buckling structures.

This paper is organized as follows. In Section 2, the Asymptotic Numerical Method (ANM) and vectorial Padé approximants are presented. The proposition of an analytical formula of the step length a_{\max}^p is detailed in Section 3. In Section 4, two benchmarks are discussed in order to compare the continuation Padé approximants using the new analytical formula of the step length a_{\max}^p with the classical Padé approximant and series continuations. The first concerns the nonlinear bending of a plate and the second concerns the buckling of a cylindrical shell loaded in the center.

2. Asymptotic Numerical Method (ANM) and vectorial Padé approximant

Many engineering problems can be reduced to solve nonlinear problems depending on a control parameter λ . These problems are written in the following general form:

$$\{R(\{U\}, \lambda)\} = \{0\} \tag{1}$$

where $\{U\}$ is the unknown vector of \mathbb{R}^n , λ is a parameter and $\{R\}$ is a vector function with values in \mathbb{R}^n assumed to be sufficiently regular with respect to its arguments $\{U\}$, and λ . The Asymptotic Numerical Method (ANM) [16,22] is a family of algorithms for path-following problems. This principle is simply to expand the unknown $(\{U\}, \lambda)$ of the nonlinear problem (1) in power series with respect to a path parameter a :

$$\begin{cases} \{U^S\}(a) - \{U^J\} &= \sum_{k=1}^N a^k \{U_k\} \\ \lambda^S(a) - \lambda^J &= \sum_{k=1}^N a^k \lambda_k \end{cases}, \quad a \in [0, a_{\max}^S] \tag{2}$$

where $(\{U^J\}, \lambda^J)$ is a started vectorial solution, $(\{U_k\}, \lambda_k)$ is a vectorial unknown and N is the truncation order of the series. The validity range $[0, a_{\max}^S]$ is deduced from the computation of the truncated vector series (2). So, the step length a_{\max}^S is computed a posteriori by the following estimation, which has been proposed in [16]:

$$a_{\max}^S = \left(\varepsilon_s \frac{\|\{U_1\}\|}{\|\{U_N\}\|} \right)^{\frac{1}{N-1}} \tag{3}$$

where ε_s is a given tolerance parameter and $\|\cdot\|$ indicates a standard norm. The step lengths depend on the definition of the path parameter a and we must add an auxiliary equation to define this parameter [16–18]. By using the evaluation of the series at $a = a_{\max}^S$, we obtain a new starting point and define, in this way, the ANM continuation procedure. This continuation method has been proved to be an efficient method to compute the solution to nonlinear partial differential equations [9,11,12,16–18,22].

In order to increase the validity range of each solution branch, we replace the vectorial representation (2) by the vectorial Padé approximant of order $M = N - 1$:

$$\begin{cases} \{U^P\}(a) - \{U^J\} &= \sum_{k=1}^M \frac{\Delta_{M-k}^M(a)}{\Delta_M^M(a)} a^k \{U_k\} \\ \lambda^P(a) - \lambda^J &= \sum_{k=1}^M \frac{\Delta_{M-k}^M(a)}{\Delta_M^M(a)} a^k \lambda_k \end{cases}, \quad a \in [0, a_{\max}^P] \tag{4}$$

where $\Delta_k^M(a) = \sum_{i=0}^k b_i^M a^i$, $0 \leq k \leq M$, with b_0^M is equal to 1, b_k^M , $1 \leq k \leq M$, are computed by a Gram–Schmidt orthonormalization technique of the vectors $\{U_k\}$, $1 \leq k \leq M$, and a_{\max}^P is a step length that determines the validity range of this approximant. Recall, for a given tolerance parameter ε_p , that the computation of the step length a_{\max}^P is done numerically by solving the following equation [13]:

$$\frac{\|\{U_M^P\}(a_{\max}^P) - \{U_{M-1}^P\}(a_{\max}^P)\|_2}{\|\{U_M^P\}(a_{\max}^P) - \{U^J\}\|_2} - \varepsilon_p = 0 \tag{5}$$

This Eq. (5) has been solved by the dichotomy method [16]. In the same way as in the case of the series continuation, we have proceed to a continuation of each branch by replacing $(\{U^J\}, \lambda^J)$ by $(\{U\}(a_{\max}^P), \lambda(a_{\max}^P))$.

In order to facilitate the computation of the relative error $E_M(a)$, we propose, in this paper, to introduce a matrix formulation used recently in [23] of the rational representation (4) in the following form:

$$\{U^P\}(a) - \{U^J\} = \frac{1}{\Delta_M^M(a)} [U_M][B_M]\{a_M\} \quad ; \quad a \in [0, a_{\max}^P] \tag{6}$$

where $[U_M] = [\{U_1\} \cdots \{U_M\}]$ is the matrix of order $NDL \times M$ (NDL is the number of degrees of freedom) whose columns are the vectors $\{U_1\} \cdots \{U_M\}$, $[B_M]$ is the upper triangular matrix of order M defined by:

$$[B_M] = \begin{bmatrix} 1 & b_1^M & \cdots & b_{M-1}^M \\ \vdots & 1 & \cdots & \vdots \\ \vdots & & \ddots & b_1^M \\ 0 & \cdots & 0 & 1 \end{bmatrix} \tag{7}$$

and ${}^t\{a_M\} = \langle a, a^2, \dots, a^M \rangle$.

The objective of the following section is to search analytically the step length a_{\max}^P by defining a relative error between two Padé approximants truncated at the consecutive orders $M - 1$, and M .

3. Proposition of an analytical formula of the step length a_{\max}^P

To establish an analytical formula of the step length a_{\max}^P , let us recall that the relative error is given by:

$$E_M(a) = \frac{\|\{U_M^P\}(a) - \{U_{M-1}^P\}(a)\|_2}{\|\{U_M^P\}(a) - \{U^J\}\|_2} \tag{8}$$

With the help of the matrix formulation (6), it is shown that the formula (8) can be rewritten in the following form:

$$E_M(a) = a^{M-1} \frac{\|\sum_{k=0}^{M-1} a^k [R_M]\{\varphi^k\}\|_2}{\|\sum_{k=0}^{2M-2} a^k [R_M]\{\eta^k\}\|_2} \tag{9}$$

where $[R_M] = (\alpha_{ij})_{1 \leq i \leq j \leq M}$ is an upper triangular square matrix of order M obtained by a factorization method QR from the matrix $[U_M]$; with α_{ij} are the orthogonalization coefficients of Gram–Schmidt, the vectors $\{\varphi^k\}$, $0 \leq k \leq M - 1$, and $\{\eta^k\}$, $0 \leq k \leq 2M - 2$ are functions of the coefficients b_k^M ($1 \leq k \leq M$) of the polynomial Δ_k^M of the Padé approximants truncated at orders M , and $M - 1$, whose expressions are detailed in Appendix A.

By a simple computation (see Appendix B), we show that the development of the relative error $E_M(a)$ at order 3 with respect to the path parameter a in the neighborhood of 0 is written as:

$$E_M(a) = a^{M-1}(e_0 + e_1 a + e_2 a^2 + e_3 a^3 + O(a^4)) \tag{10}$$

where the computations of e_0 , e_1 , e_2 , and e_3 are detailed in Appendix B. Taking into account Eqs. (5) and (10), the step length a_{\max}^P is solution to the following equation:

$$E_M(a) - \varepsilon_p = 0 \tag{11}$$

The resolution of Eq. (11) consists in computing the roots of the polynomial $P(a)$ given by (see Appendix B):

$$P(a) = a^3 + \beta_2 a^2 + \beta_1 a + \beta_0 \tag{12}$$

where the coefficients β_0 , β_1 , and β_2 are given in Appendix B. In order to determine the roots of the polynomial $P(a)$, we replace the parameter a by $z - \frac{\beta_2}{3}$, which allows us to obtain the following equation verified by z :

$$z^3 + pz + q = 0 \tag{13}$$

where $p = \beta_1 - \frac{\beta_2^2}{3}$, and $q = \frac{2\beta_2^3}{27} - \frac{\beta_1\beta_2}{3} + \beta_0$. By introducing u , and v such that:

$$z = u + v, \tag{14}$$

Eq. (13) becomes:

$$u^3 + v^3 + (u + v)(3uv + p) + q = 0 \tag{15}$$

Therefore, if u , and v verify the following system

$$\begin{cases} u^3 + v^3 = -q \\ uv = -\frac{p}{3} \end{cases} \tag{16}$$

then u^3 , and v^3 are the roots of the trinomial:

$$X^2 + qX - \frac{p^3}{27} = 0 \tag{17}$$

Then the roots of (17) for a positive discriminant $\Delta_0 = q^2 + 4\frac{p^3}{27} > 0$ are given by:

$$\begin{cases} u^3 = \frac{-q + \sqrt{\Delta_0}}{2} \\ v^3 = \frac{-q - \sqrt{\Delta_0}}{2} \end{cases} \tag{18}$$

Knowing the roots u , and v , then $z_1 = u + v$ is a root of (13). In the case where the discriminant is negative, there exists $\theta \in \mathbb{R}$ such that u^3 , and v^3 verify

$$\begin{cases} u^3 = \frac{-q + i\sqrt{-\Delta_0}}{2} = \sqrt{\frac{q^2 - \Delta_0}{4}} \exp^{i\theta} \\ v^3 = \frac{-q - i\sqrt{-\Delta_0}}{2} = \sqrt{\frac{q^2 - \Delta_0}{4}} \exp^{-i\theta} \end{cases} \tag{19}$$

Therefore, $z_1 = \left(\frac{q^2 - \Delta_0}{4}\right)^{\frac{1}{6}} \left[\exp\left(\frac{i\theta}{3}\right) + \exp\left(\frac{-i\theta}{3}\right)\right] = 2\left(\frac{q^2 - \Delta_0}{4}\right)^{\frac{1}{6}} \cos\left(\frac{\theta}{3}\right)$ is a root of (13). Note that the equalities in (19) show that:

$$\theta = \arccos\left(\frac{-q}{\sqrt{q^2 - \Delta_0}}\right) \tag{20}$$

Knowing a root z_1 of (13), the other two roots z_2 , and z_3 of (13) are those of the polynomial of degree 2 obtained by the Euclidean division of (13) by $z - z_1$. So, the roots of the polynomial $P(a)$ are given by:

$$\begin{cases} a_1 = z_1 - \frac{\beta_2}{3} \\ a_2 = z_2 - \frac{\beta_2}{3} \\ a_3 = z_3 - \frac{\beta_2}{3} \end{cases} \tag{21}$$

Finally, the step length a_{\max}^p is given by the smallest positive root:

$$a_{\max}^p = \min(a_1, a_2, a_3) \tag{22}$$

4. Numerical applications

In this section, we have computed analytically the smallest positive real root of (19) that we have compared with the value of a_{\max}^p obtained by the dichotomy method applied to Eq. (7) for the same residual. The used coefficients b_i^M of the Padé approximant are those computed by the Gram–Schmidt orthonormalization technique of the vectors $\{U_i\}$ ($1 \leq i \leq M$). A comparison between this new formula (19) and the one determined numerically by the dichotomy method is illustrated on examples of bending and buckling of structures. The DKT18 element is used for finite element analysis [24,25].

4.1. Nonlinear bending of a square plate

We consider the nonlinear bending of an elastic, homogeneous and isotropic square plate of length $L = 200$ mm and thickness $h = 1$ mm, embedded at the four edges and subjected to vertical loading λF in its center, with $F = 4$ N. The

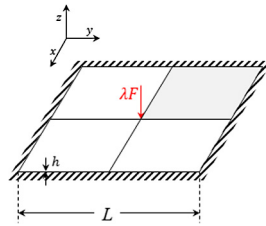
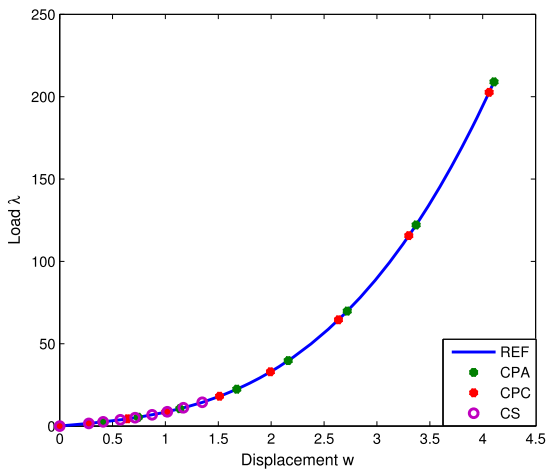


Fig. 1. Geometrical characteristics of the plate.

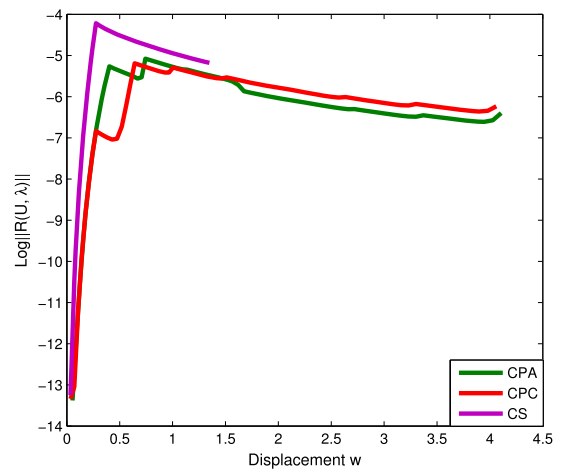
Table 1

Comparison between the results obtained by 8 ANM steps of the three types of continuation CS, CPC, and CPA for the order $N = 10$ and for the tolerance parameters $\epsilon_s = 1 \cdot 10^{-7}$, $\epsilon_p = 7.98 \cdot 10^{-5}$, and $\epsilon_p = 1.5 \cdot 10^{-3}$, respectively.

Order	CS, $\epsilon_s = 1.0 \cdot 10^{-7}$			CPC, $\epsilon_p = 7.89 \cdot 10^{-5}$			CPA, $\epsilon_p = 1.5 \cdot 10^{-3}$		
	N (steps)	λ	w	N (steps)	λ	w	N (steps)	λ	w
10	8	14.41	1.34	8	202.6	4.06	8	208.9	4.10



(a) Load-displacement curve



(b) Residual-displacement curve

Fig. 2. Load-displacement and residual-displacement curves obtained by 8 ANM steps of the three algorithms CS, CPC, and CPA for the truncation order $N = 10$.

mechanical characteristics of this plate are: Young’s modulus $E = 3 \cdot 10^4$ MPa and Poisson’s ratio $\nu = 0.3$ (see Fig. 1). The plate is discretized in 32 DKT18 finite elements, i.e. 150 degrees of freedom.

In this numerical experiment, the obtained results are compared between three types of continuations. The first is a continuation series CS using the representation (2), the second is a continuation Padé approximant using the representation (4), where the step length a_{max}^p is computed by the dichotomy method CPC and the third is a continuation Padé approximant using the representation (4) where the step length a_{max}^p is defined by the analytical formula (19) CPA.

In Table 1, we report the results obtained by 8 ANM steps with the truncation orders $N = 10$ for the same residual. Note that with the continuation series CS, we have reached a deflection $w = 1.34$ mm, however with the continuation Padé of type CPC, we have reached a deflection $w = 4.06$ mm. If we use the continuation Padé of type CPA with the new formula of step length a_{max}^p , we have reached a deflection $w = 4.10$ mm. From the obtained results, we note that for $N = 10$, the proposed analytical formula works well and gives same results as the dichotomy method.

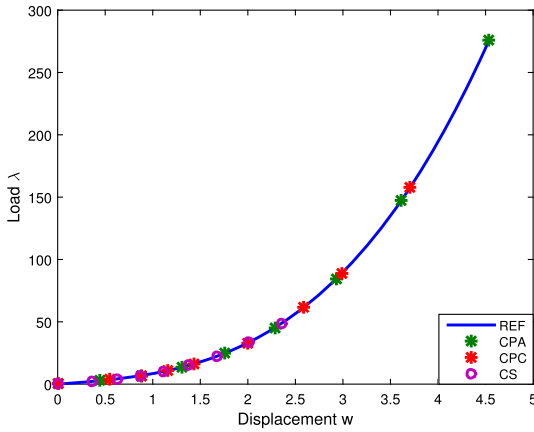
In Fig. 2, we present a comparison between the results obtained at the forced node by the continuation series CS and the continuation Padé approximant using the two definitions of the step length a_{max}^p CPC, and CPA. These results are obtained with the truncation order $N = 10$ in 8 ANM steps and with a tolerance parameter that gives the same quality of the solution (see Fig. 2b). From these results, we note that the two continuations of Padé approximants are comparable.

In Table 2, we show a comparison between the results obtained by the continuation series CS and the continuation Padé where the step length a_{max}^p is computed by the dichotomy method and by the proposed analytical formula (19). These results are obtained with a truncation order $N = 16$ in 8 ANM steps. According to this table, we note that, also for this

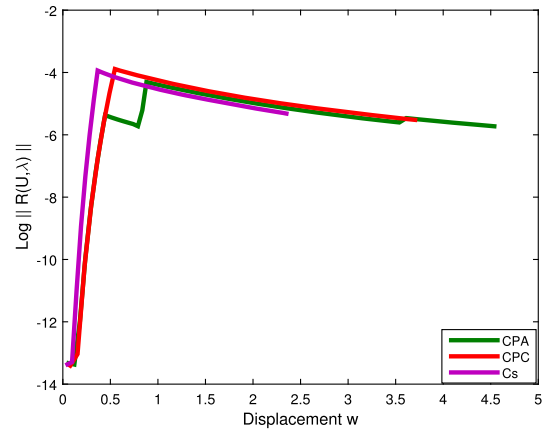
Table 2

Comparison between the results obtained by 8 ANM steps of the three types of continuation *CS*, *CPC*, and *CPA* for the order $N = 16$ and for the tolerance parameters $\varepsilon_s = 1.0 \cdot 10^{-7}$, $\varepsilon_p = 1.0 \cdot 10^{-4.2}$, and $\varepsilon_p = 1.1 \cdot 10^{-3}$, respectively.

Order	<i>CS</i> , $\varepsilon_s = 1.0 \cdot 10^{-7}$			<i>CPC</i> , $\varepsilon_p = 1.0 \cdot 10^{-4.2}$			<i>CPA</i> , $\varepsilon_p = 1.1 \cdot 10^{-3}$		
	N (steps)	λ	w	N (steps)	λ	w	N (steps)	λ	w
16	8	48.59	2.35	8	157.5	3.70	8	275.8	4.53



(a) Load-displacement curve



(b) Residual-displacement curve

Fig. 3. Load-displacement and Residual-displacement curves obtained by 8 ANM steps of the three algorithms *CS*, *CPC*, and *CPA* for the truncation order $N = 16$.

Table 3

Comparison between the results obtained by 8 ANM steps of the three types of continuation *CS*, *CPC*, and *CPA* for the order $N = 20$ and for the tolerance parameters $\varepsilon_s = 10^{-7}$, $\varepsilon_p = 5 \cdot 10^{-5}$, and $\varepsilon_p = 3.211 \cdot 10^{-4}$, respectively.

Order	<i>CS</i> , $\varepsilon_s = 10^{-7}$			<i>CPC</i> , $\varepsilon_p = 5 \cdot 10^{-5}$			<i>CPA</i> , $\varepsilon_p = 3.211 \cdot 10^{-4}$		
	N (steps)	λ	w	N (steps)	λ	w	N (steps)	λ	w
20	8	76.08	2.81	8	178.2	3.87	8	272.9	4.52

choice of order $N = 16$, the continuation Padé approximant using the analytical formula (19) to compute the step length a_{\max}^p gives good results. With 8 ANM steps, we reach a deflection of $w = 4.53$.

Fig. 3 illustrates the load-displacement and residual-displacement curves at the forced node obtained by 8 ANM steps of the three continuations *CS*, *CPC*, and *CPA* for the truncation order $N = 16$ and for the tolerance parameter that gives the same quality of solution.

Table 3 represents a comparison between the results obtained by the three continuations *CS*, *CPC*, and *CPA* for the truncation order $N = 20$ in 8 ANM steps and for tolerance parameters that give the same quality of solution. From the obtained results, it is clear that the step length of the continuation *CPA* are greater here than those of the two other continuations.

Fig. 4 illustrates the load-displacement and residual-displacement curves at the forced node obtained by 8 ANM steps of the three continuations *CS*, *CPC*, and *CPA* for the truncation order $N = 20$ and for the tolerance parameter that gives the same quality of solution.

4.2. Buckling of a cylindrical shell loaded in the center

This second application concerns the buckling of a cylindrical shell articulated along the two opposite edges and free on the two other ones. This structure is of length $2L = 504$ mm, of radius $R = 2540$ mm, and has an angle of half opening $\theta = 0.1$ rad, made of a homogeneous, elastic, and isotropic material of Young’s modulus $E = 3102.75$ MPa and Poisson’s coefficient $\nu = 0.3$, subjected to a vertical loading λF applied at its center, with $F = 1000$ N. Due to the symmetry, only a quarter of the structure is discretized in 200 triangular elements of type *DKT18*, i.e. 726 degrees of freedom. Two types of thickness are studied in this application: $h = 12.7$ mm and $h = 6.35$ mm (see Fig. 5). For this benchmark example, the response curve based on the value of the thickness h has limit points in λ , and w . This is a good test for the continuation methods [12,22].

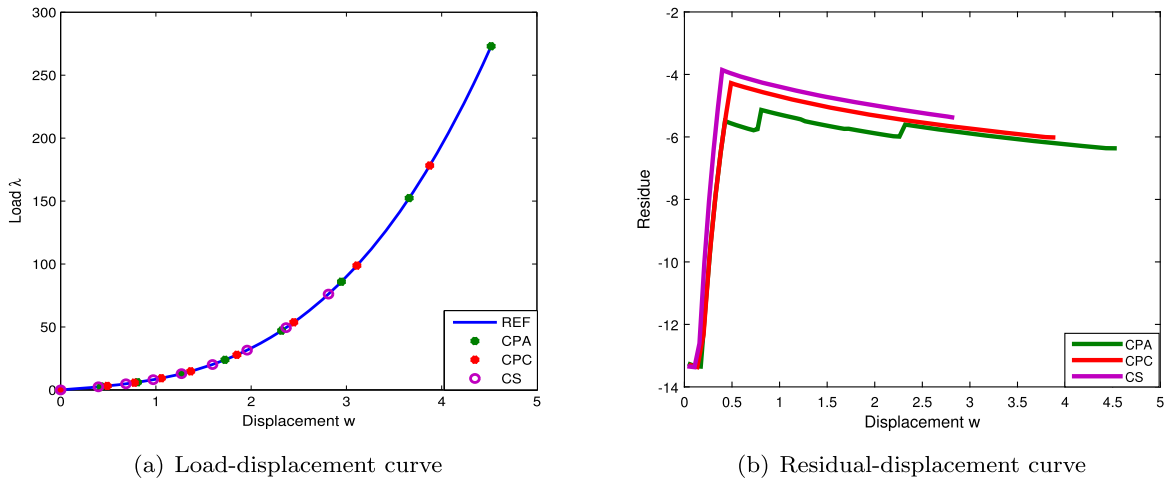
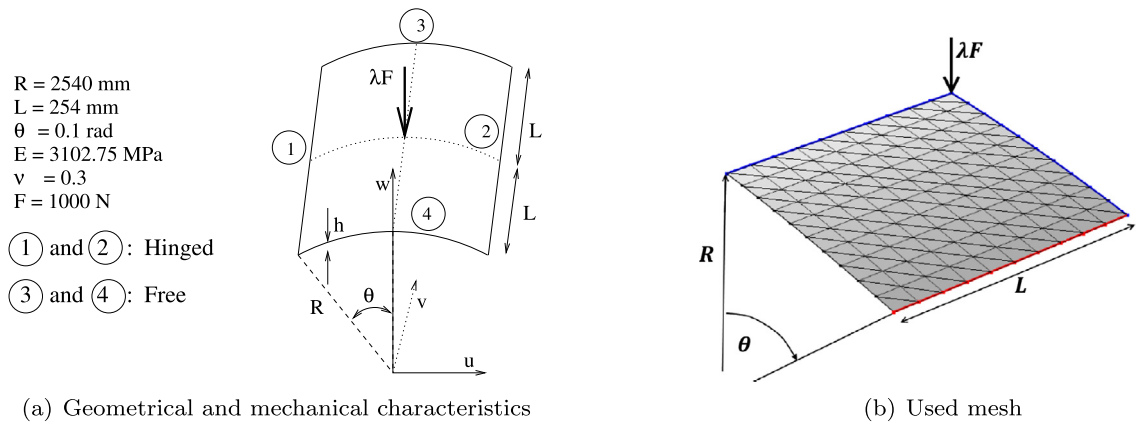


Fig. 4. Load-displacement and residual-displacement curves obtained by 8 ANM steps of the three algorithms CS, CPC, and CPA for the truncation order $N = 20$.



$R = 2540$ mm
 $L = 254$ mm
 $\theta = 0.1$ rad
 $E = 3102.75$ MPa
 $\nu = 0.3$
 $F = 1000$ N

① and ②: Hinged
 ③ and ④: Free

Fig. 5. Geometrical and mechanical characteristics of the studied structure and the used mesh.

Table 4

Comparison between the results obtained by 6 ANM steps of the three types of continuation CS, CPC, and CPA for the order $N = 10$ and for the tolerance parameters $\epsilon_s = 3 \cdot 10^{-6}$, $\epsilon_p = 10^{-6}$, and $\epsilon_p = 1.7 \cdot 10^{-5.1}$, respectively.

Order	CS, $\epsilon_s = 3 \cdot 10^{-6}$			CPC, $\epsilon_p = 10^{-6}$			CPA, $\epsilon_p = 1.7 \cdot 10^{-5.1}$		
	N (steps)	λ	w	N (steps)	λ	w	N (steps)	λ	w
10	6	0.75	22.01	6	0.78	22.2	6	1.27	24.4

In this second numerical experiment, we treat first the buckling of a thick cylindrical shell ($h = 12.7$ mm) while comparing the results obtained by the three types of continuations CS, CPC, and CPA. We made the same numerical study as in the first example.

In Table 4, we report the results obtained by 6 ANM steps with the truncation order $N = 10$ for the same residual. The parameters ϵ_s , and ϵ_p have been chosen in order to get the same residual. Note that with the continuation series CS, we have reached a deflection $w = 22.01$ mm, however with the continuation Padé approximant CPC, we have reached a deflection $w = 22.2$ mm. If we use the continuation Padé approximant CPA with the new proposed formula (19) of a_{max}^p , we have reached a deflection $w = 24.4$ mm. From the obtained results, we note that for $N = 10$, the proposed analytical formula is less better than the two other in this example.

In Fig. 6, we present a comparison between the results obtained at the forced node by the continuation series and the continuation Padé using the two definitions of the step length a_{max}^p . These results are obtained with the truncation order $N = 10$ in 6 ANM steps and with a tolerance parameter that gives the same quality of solution (see Fig. 6b).

Table 5 represents a comparison between the results obtained by the three continuations CS, CPC, and CPA for the truncation order $N = 15$ in 5 ANM steps and for the tolerance parameters ϵ_s , and ϵ_p that give the same quality of solution.

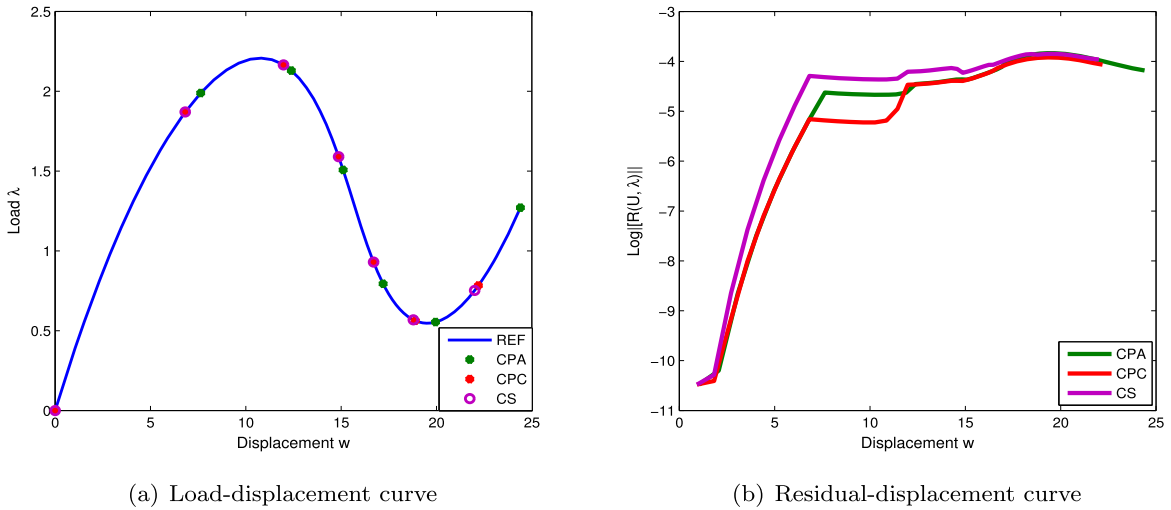


Fig. 6. Load-displacement and residual curves for $h = 12.7$ mm and $N = 10$ obtained by 6 ANM steps of the three continuations CS, CPC, and CPA.

Table 5

Comparison between the results obtained by 5 ANM steps of the three types of continuation CS, CPC, and CPA for the order $N = 15$ and for the tolerance parameters $\epsilon_s = 2.0 \cdot 10^{-6}$, $\epsilon_p = 10^{-6}$, and $\epsilon_p = 8.0 \cdot 10^{-5}$, respectively.

Order	CS, $\epsilon_s = 2.0 \cdot 10^{-6}$			CPC, $\epsilon_p = 10^{-6}$			CPA, $\epsilon_p = 8.0 \cdot 10^{-5}$		
	N (steps)	λ	w	N (steps)	λ	w	N (steps)	λ	w
15	5	2.47	27.8	5	17.38	44.44	5	28.53	51.03

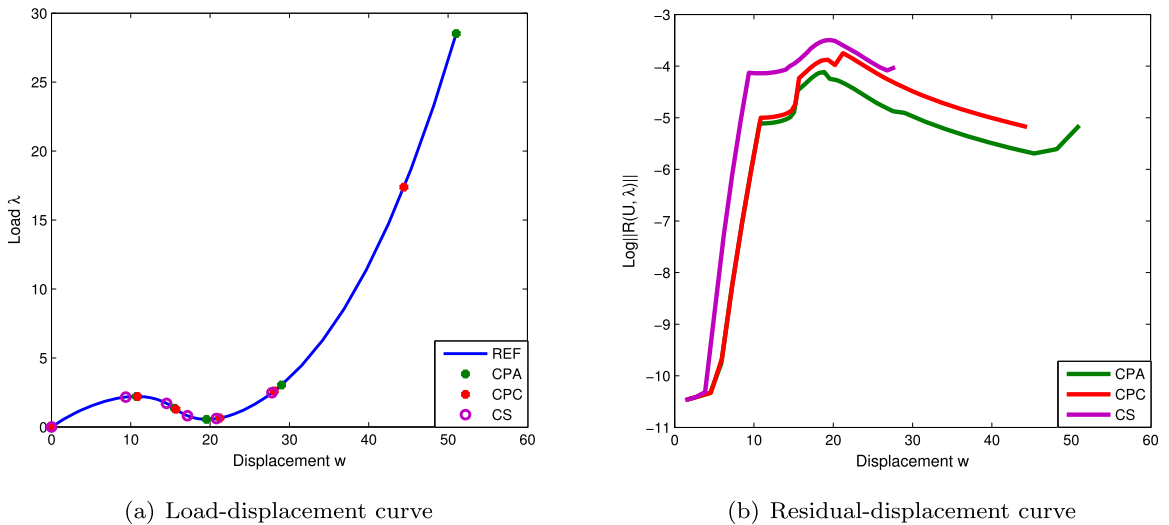


Fig. 7. Load-displacement and residual curves for $h = 12.7$ mm and $N = 15$ obtained by 5 ANM steps of the three continuations CS, CPC, and CPA.

From the obtained results, it is clear that the step length of the continuation Padé approximant CPA is greater here than those of the two other continuations.

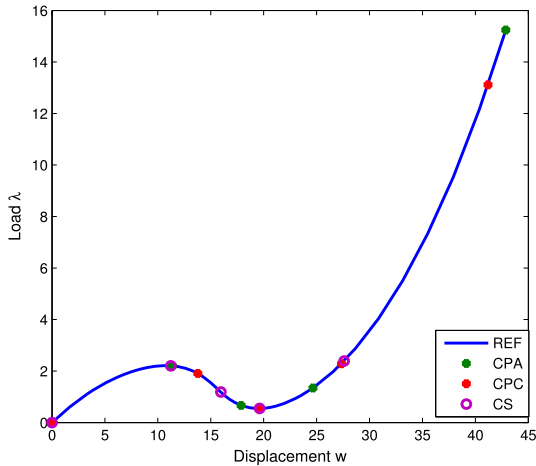
In Fig. 7, we present a comparison between the results obtained at the forced node by the continuation series and the continuation Padé approximant using the two definitions of the step length a_{max}^p . These results are obtained with the truncation order $N = 15$ in 5 ANM steps and with a tolerance parameter that gives the same quality of solution (see Fig. 7b).

Table 6 provides a comparison between the results obtained by the three continuations CS, CPC, and CPA for the truncation order $N = 20$ in 4 ANM steps and for the tolerance parameters ϵ_s , and ϵ_p that give the same quality of solution. From the obtained results, it is clear that the step length of the continuation CPA is greater here than those of the two other continuations.

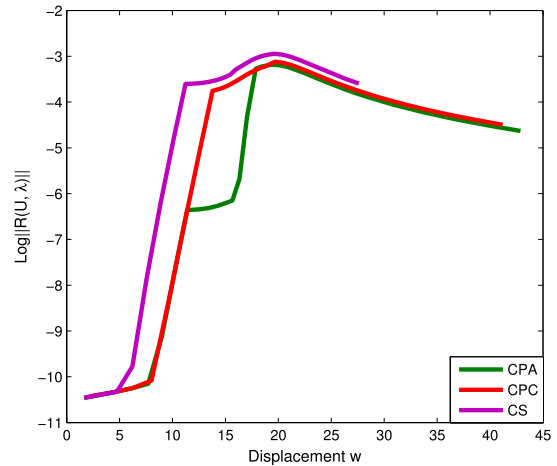
Table 6

Comparison between the results obtained by 4 ANM steps of the three types of continuation *CS*, *CPC*, and *CPA* for the order $N = 20$ and for the tolerance parameters $\varepsilon_s = 2.0 \cdot 10^{-6}$, $\varepsilon_p = 5.0 \cdot 10^{-6}$, and $\varepsilon_p = 10^{-4.1}$, respectively.

Order	<i>CS</i> , $\varepsilon_s = 2.0 \cdot 10^{-6}$			<i>CPC</i> , $\varepsilon_p = 5.0 \cdot 10^{-6}$			<i>CPA</i> , $\varepsilon_p = 10^{-4.1}$		
	N (steps)	λ	w	N (steps)	λ	w	N (steps)	λ	w
20	4	2.39	27.6	4	13.11	41.19	4	15.24	42.88



(a) Load-displacement curve



(b) Residual-displacement curve

Fig. 8. Load-displacement and residual-displacement curves for $h = 12.7$ mm and $N = 20$ obtained by 4 ANM steps of the three continuations *CS*, *CPC*, and *CPA*.

Table 7

Comparison between the results obtained by 22 ANM steps of the three types of continuation *CS*, *CPC*, and *CPA* for the order $N = 10$ and for the tolerance parameters $\varepsilon_s = 6.0 \cdot 10^{-7}$, $\varepsilon_p = 12.0 \cdot 10^{-7}$, and $\varepsilon_p = 9.3 \cdot 10^{-6}$, respectively.

Order	<i>CS</i> , $\varepsilon_s = 6.0 \cdot 10^{-7}$			<i>CPC</i> , $\varepsilon_p = 12.0 \cdot 10^{-7}$			<i>CPA</i> , $\varepsilon_p = 9.3 \cdot 10^{-6}$		
	N (steps)	λ	w	N (steps)	λ	w	N (steps)	λ	w
10	22	-0.36	17.38	22	0.67	29.97	22	0.64	29.80

In Fig. 8, we present a comparison between the results obtained at the forced node by the continuation series and the continuation Padé approximant using the two definitions of the step length a_{max}^p . These results are obtained with the truncation order $N = 20$ in 4 ANM steps and with tolerance parameters ε_s , and ε_p that give the same quality of solution (see Fig. 8b).

In this second numerical experiment, we treat also the buckling of a thin cylindrical shell ($h = 6.35$ mm) while comparing the results obtained by the three continuations *CS*, *CPC*, and *CPA*.

In Table 7, we report the results obtained by 22 ANM steps with the truncation order $N = 10$ for the same residual. Note that with the continuation series *CS*, we have reached a deflection $w = 17.38$ mm, however with the continuation Padé approximant *CPC*, we have reached a deflection $w = 29.97$ mm. If we use the continuation Padé *CPA* with the new formula (19) of the step length a_{max}^p , we have reached a deflection $w = 29.80$ mm.

In Fig. 9, we present a comparison between the results obtained at the forced node by 22 ANM steps of the three continuations *CS*, *CPC*, and *CPA* with the truncation order $N = 10$ and with a tolerance parameter that gives the same quality of solution (see Fig. 9b).

Table 8 presents a comparison between the results obtained by 16 ANM steps of the three continuations *CS*, *CPC*, and *CPA* for the truncation order $N = 15$ and for the tolerance parameter that gives the same quality of solution. From the obtained results, it is clear that the step length of the continuation *CPA* is less here than those of the two other continuations.

In Fig. 10, we present a comparison between the results obtained at the forced node by 16 ANM steps of the three continuations *CS*, *CPC*, and *CPA*. These results are obtained with the truncation order $N = 15$ and with a tolerance parameter that gives the same quality of solution (see Fig. 10b).

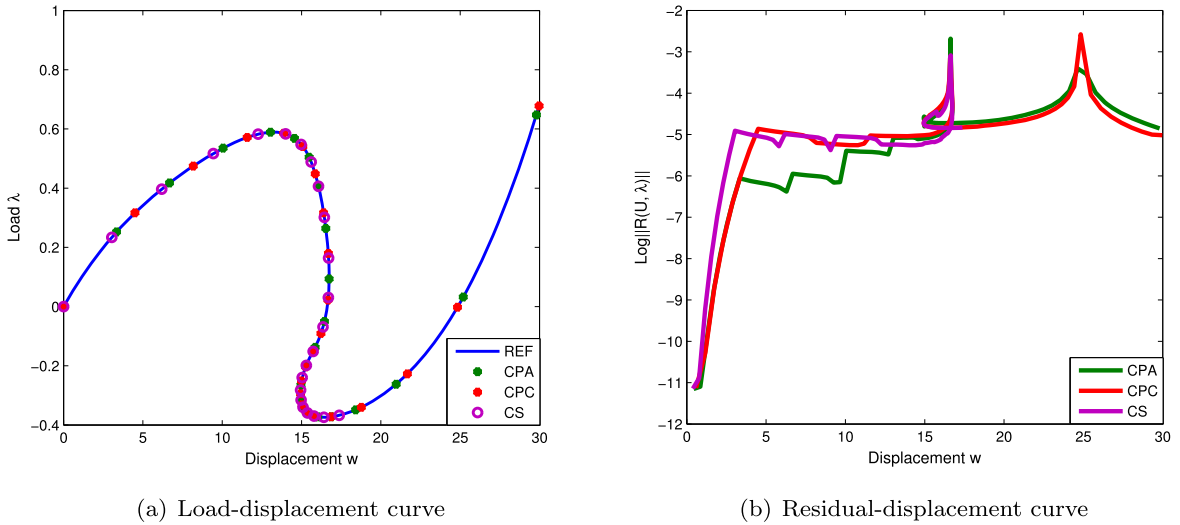


Fig. 9. Load-displacement and residual-displacement curves for $h = 6.35$ mm and $N = 10$ obtained by 22 ANM steps of the three continuations CS, CPC, and CPA.

Table 8

Comparison between the results obtained by 16 ANM steps of the three types of continuation CS, CPC, and CPA for the order $N = 15$ and for the tolerance parameters $\epsilon_s = 6.0 \cdot 10^{-8}$, $\epsilon_p = 6.0 \cdot 10^{-8}$, and $\epsilon_p = 12.0 \cdot 10^{-7}$, respectively.

Order	CS, $\epsilon_s = 6.0 \cdot 10^{-8}$			CPC, $\epsilon_p = 6.0 \cdot 10^{-8}$			CPA, $\epsilon_p = 12.0 \cdot 10^{-7}$		
	N (steps)	λ	w	N (steps)	λ	w	N (steps)	λ	w
15	16	-0.37	16.34	16	0.31	27.62	16	0.18	26.58

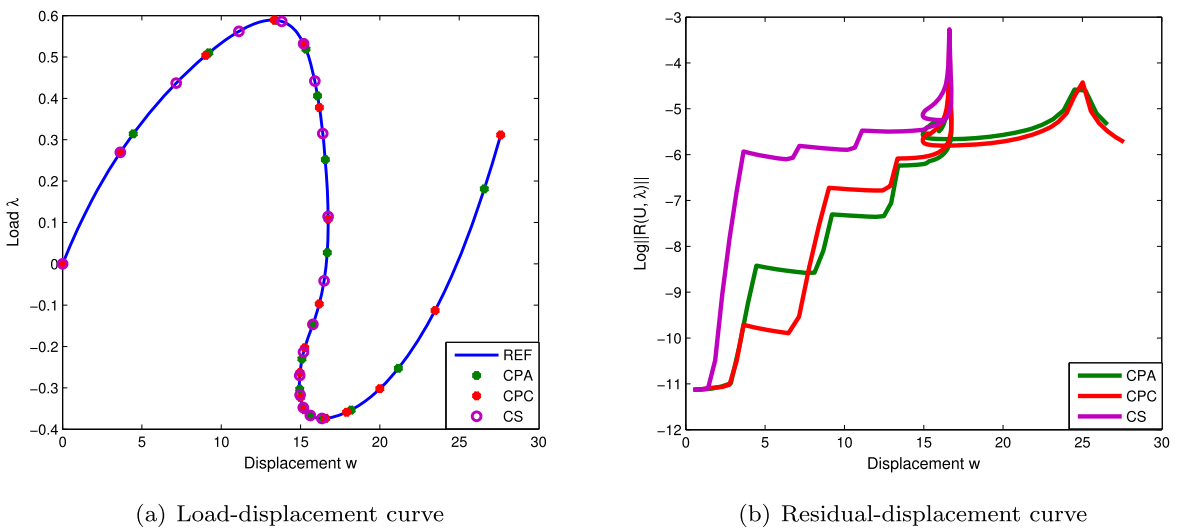


Fig. 10. Load-displacement and residual curves for $h = 6.35$ mm and $N = 15$ obtained by 16 ANM steps of the three continuations CS, CPC, and CPA.

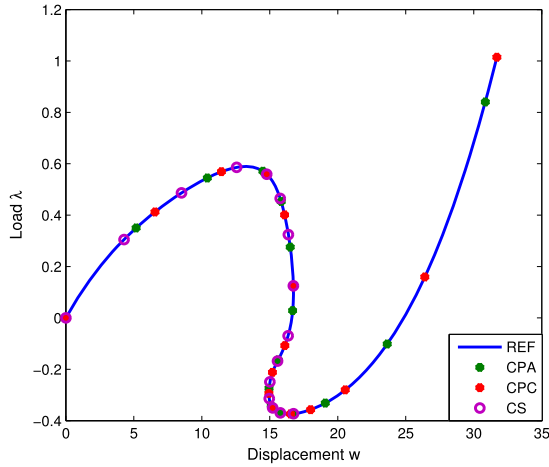
Table 9 presents a comparison between the results obtained by 14 ANM steps of the three continuations CS, CPC, and CPA for the truncation order $N = 20$ and for the tolerance parameter that gives the same quality of solution. From the obtained results, it is clear that the step length of the continuation CPA is less here than those of the two other continuations.

In Fig. 11, we present a comparison between the results obtained at the forced node by 14 ANM steps of the three continuations CS, CPC, and CPA. These results are obtained with the truncation order $N = 20$ and with a tolerance parameter that gives the same quality of solution (see Fig. 11b).

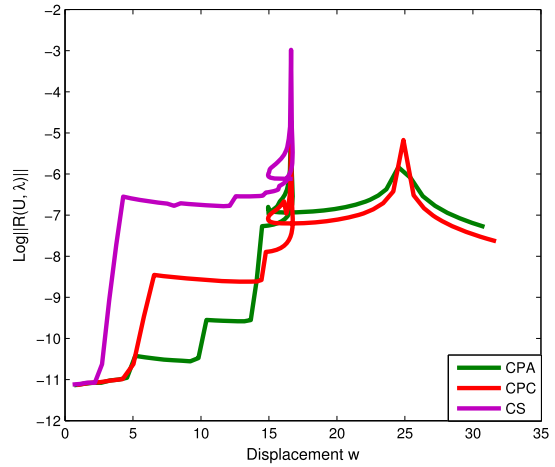
Table 9

Comparison between the results obtained by 14 ANM steps of the three continuations *CS*, *CPC*, and *CPA* for the order $N = 20$ and for the tolerance parameters $\varepsilon_s = 10^{-8}$, $\varepsilon_p = 10^{-9.5}$, and $\varepsilon_p = 10^{-6.14}$, respectively.

Order	<i>CS</i> , $\varepsilon_s = 10^{-8}$			<i>CPC</i> , $\varepsilon_p = 10^{-9.5}$			<i>CPA</i> , $\varepsilon_p = 10^{-6.14}$		
	N (steps)	λ	w	N (steps)	λ	w	N (steps)	λ	w
20	14	-0.37	16.74	14	1.01	31.7	14	0.84	30.85



(a) Load-displacement curve



(b) Residual-displacement curve

Fig. 11. Load-displacement and residual curves for $h = 6.35$ mm and $N = 20$ obtained by 14 ANM steps of the three continuations *CS*, *CPC*, and *CPA*.

It should be noted that the continuation *CPC* uses the numerical solution a_{\max}^p to Eq. (5) when the pole of the denominator of the Padé approximants is greater than a_{\max}^s of Eq. (3) and otherwise it uses a_{\max}^s , whereas the continuation *CPA* always uses a_{\max}^p of Eq. (22).

5. Conclusion

In this work, we have proposed an analytical formula of the step length a_{\max}^p that defines the validity range of Padé approximants in the ANM algorithm. In order to determine this formula, we have defined the relative error between two consecutive vectorial Padé approximants. Then, with the help of a matrix formulation, we have developed this relative error at the order 3. From the obtained results, it can be seen that, when this new analytical formula, is used it is possible to ensure a purely Padé continuation. Moreover, it is clear that the proposed formula makes it possible to have comparable step lengths with the same residual compared to continuations of series *CPS* and Padé using the dichotomy method *CPC*. It should be noted also that the continuation *CPC* requires, in addition, the computation time of a_{\max}^p from Eq. (5).

Appendix A. Computation of the vectors $\{\varphi^k\}$, and $\{\eta^k\}$

The components φ_i^k , $1 \leq i \leq M$ of vectors $\{\varphi^k\}$, $0 \leq k \leq M - 1$ that appear in Eq. (9) are such that:

$$\varphi_i^0 = \begin{cases} b_{M-i}^{M-1} & \text{for } 1 \leq i \leq M - 1 \\ 1 & \text{for } i = M \end{cases}, \varphi_i^{M-1} = \begin{cases} b_{M-1}^{M-1} b_{M-i}^M - b_M^M b_{M-i-1}^{M-1} & \text{for } 1 \leq i \leq M - 1 \\ b_{M-1}^{M-1} & \text{for } i = M \end{cases} \quad (23)$$

and, for $1 \leq k \leq M - 2$,

$$\varphi_i^k = \begin{cases} b_k^{M-1} b_{M-i}^M + b_{M-1}^{M-1} b_{k-i+1}^M + \sum_{n=k+1}^{M-2} (b_n^{M-1} b_{M+k-i-n}^M - b_n^M b_{M+k-i-n}^{M-1}) \\ \quad - b_{M-1}^M b_{k-i+1}^{M-1} - b_M^M b_{k-i}^{M-1} & \text{for } 1 \leq i \leq k, \\ b_k^{M-1} b_{M-i}^M + \sum_{n=k+1}^{M+k-i} (b_n^{M-1} b_{M+k-i-n}^M - b_n^M b_{M+k-i-n}^{M-1}) & \text{for } k + 1 \leq i \leq M - 1 \\ b_k^{M-1} & \text{for } i = M \end{cases} \quad (24)$$

The other vectors $\{\eta^k\}$ for $0 \leq k \leq 2M - 2$ are given by:

$$\{\eta^k\} = \begin{cases} \langle 1, 0, \dots, 0 \rangle & \text{for } k = 0 \\ \langle f_k, f_{k-1}, \dots, f_1, 1, 0, \dots, 0 \rangle & \text{for } 1 \leq k \leq M - 1 \\ \langle g_{1,k-M}^M, g_{2,k-M}^M, \dots, g_{M,k-M}^M \rangle & \text{for } M \leq k \leq 2M - 1 \end{cases} \tag{25}$$

where the components f_k , and $g_{i,k-M}^M$ are given by:

$$\begin{cases} f_k = \sum_{n=0}^k b_n^M b_{k-n}^{M-1} & \text{for } 1 \leq k \leq M - 1 \\ g_{i,k}^M = \sum_{n=k}^{\min(M-1; M+k-i)} b_n^{M-1} b_{M+k-i-n}^M & \text{for } 1 \leq i \leq M, \quad 1 \leq k \leq M - 1 \end{cases} \tag{26}$$

as explained in the paper [23].

Appendix B. Computation of the polynomial $P(a)$

In order to approach $E_M(a) - \varepsilon_p$ of Eq. (11) by its truncated development at the order 3 with respect to the parameter a in the neighborhood of 0, let

$$\begin{aligned} \{\gamma^k\} &= [R_M]\{\varphi^k\} \text{ for } 0 \leq k \leq M - 1 \\ \{\xi^k\} &= [R_M]\{\eta^k\} \text{ for } 0 \leq k \leq 2M - 2 \end{aligned} \tag{27}$$

We have

$$E_M^2(a) = a^{2M-2} \frac{H(a)}{Q(a)} \tag{28}$$

where $H(a)$, and $Q(a)$ are given by:

$$\begin{cases} H(a) = \|\{\gamma^0\}\|_2^2 + (2 \langle \gamma^0 \rangle \{\gamma^1\})a + (\langle \gamma^0 \rangle \{\gamma^2\} + \|\{\gamma^1\}\|_2^2)a^2 + 2(\langle \gamma^1 \rangle \{\gamma^2\} + \langle \gamma^0 \rangle \{\gamma^3\})a^3 + O(a^4) \\ Q(a) = \|\{\xi^0\}\|_2^2 + (2 \langle \xi^0 \rangle \{\xi^1\})a + (\langle \xi^0 \rangle \{\xi^2\} + \|\{\xi^1\}\|_2^2)a^2 + 2(\langle \xi^1 \rangle \{\xi^2\} + \langle \xi^0 \rangle \{\xi^3\})a^3 + O(a^4) \end{cases} \tag{29}$$

Subsequently, for a small enough, we have:

$$E_M(a) = a^{M-1} (e_0 + e_1 a + e_2 a^2 + e_3 a^3 + O(a^4)) \tag{30}$$

where e_0, e_1, e_2 , and e_3 are given by:

$$\begin{cases} e_0 = \left(\frac{\|\{\gamma^0\}\|_2^2}{\|\{\xi^0\}\|_2^2} \right)^{\frac{1}{2}} \\ e_1 = \frac{e_0 B}{2} \\ e_2 = e_0 \left(\frac{G}{2} - \frac{B^2}{8} \right) \\ e_3 = \frac{e_0 H}{2} - \frac{e_0 B G}{4} + \frac{e_0 H^3}{16} \end{cases} \tag{31}$$

with

$$\left\{ \begin{aligned}
 B &= \frac{2 \langle \gamma^0 \rangle \{\gamma^1\}}{\|\{\gamma^0\}\|_2^2} - \frac{2 \langle \xi^0 \rangle \{\xi^1\}}{\|\{\xi^0\}\|_2^2} \\
 C &= \left(\frac{2 \langle \xi^0 \rangle \{\xi^1\}}{\|\{\xi^0\}\|_2^2} \right)^2 - \frac{2 \langle \xi^0 \rangle \{\xi^2\} + \|\{\xi^1\}\|_2^2}{\|\{\xi^0\}\|_2^2} \\
 D &= -4 \frac{\langle \xi^0 \rangle \{\xi^1\} \langle \gamma^0 \rangle \{\gamma^1\}}{\|\{\xi^0\}\|_2^2} \\
 E &= 2 \langle \gamma^0 \rangle \{\gamma^2\} + \|\{\gamma^1\}\|_2^2 \\
 F &= \|\{\gamma^0\}\|_2^2 \left(4 \frac{\langle \xi^0 \rangle \{\xi^1\}}{\|\{\xi^0\}\|_2^2} \frac{(2 \langle \xi^0 \rangle \{\xi^2\} + \|\{\xi^1\}\|_2^2)}{\|\{\xi^0\}\|_2^2} - 8 \left(\frac{\langle \xi^0 \rangle \{\xi^1\}}{\|\{\xi^0\}\|_2^2} \right)^3 - 2 \frac{\langle \xi^1 \rangle \{\xi^2\} + \langle \xi^0 \rangle \{\xi^3\}}{\|\{\xi^0\}\|_2^2} \right) \\
 &\quad + 2 (\langle \gamma^1 \rangle \{\gamma^2\} + \langle \gamma^0 \rangle \{\gamma^3\}) \\
 G &= \frac{\|\{\gamma^0\}\|_2^2 C + E + D}{\|\{\gamma^0\}\|_2^2} \\
 H &= \frac{2 \langle \gamma^0 \rangle \{\gamma^1\} C - 2 \frac{\langle \xi^0 \rangle \{\xi^1\}}{\|\{\xi^0\}\|_2^2} E + F}{\|\{\gamma^0\}\|_2^2}
 \end{aligned} \right. \tag{32}$$

As $E_M(a) = \varepsilon_p$, then

$$a^{M-1} (e_0 + e_1 a + e_2 a^2 + e_3 a^3 + O(a^4)) = \varepsilon_p \tag{33}$$

we deduce

$$a (e_0 + e_1 a + e_2 a^2 + e_3 a^3 + O(a^4))^{\frac{1}{M-1}} - \varepsilon_p^{\frac{1}{M-1}} = 0 \tag{34}$$

which gives, after development truncated at the order 3 with respect to the parameter a , the following equality

$$A_3 a^3 + A_2 a^2 + a - A_1 + O(a^4) = 0 \tag{35}$$

where A_1 , A_2 , and A_3 are given by:

$$\left\{ \begin{aligned}
 A_1 &= \left(\varepsilon_p^2 \frac{\|\{\xi^0\}\|_2^2}{\|\{\gamma^0\}\|_2^2} \right)^{\frac{1}{2M-2}} \\
 A_2 &= \frac{B}{2M-2} \\
 A_3 &= \frac{1}{2M-2} \left(\frac{\|\{\gamma^0\}\|_2^2 C + E + D}{\|\{\gamma^0\}\|_2^2} \right) + \frac{3-2M}{8(M-1)^2} B
 \end{aligned} \right. \tag{36}$$

By neglecting in (35) the term $O(a^4)$, we get the following equation:

$$A_3 a^3 + A_2 a^2 + a - A_1 = 0 \tag{37}$$

which is written as:

$$P(a) = a^3 + \beta_2 a^2 + \beta_1 a + \beta_0 = 0 \tag{38}$$

where β_0 , β_1 , and β_2 are given by:

$$\begin{cases} \beta_0 = -\frac{A_1}{A_3} \\ \beta_1 = \frac{1}{A_3} \\ \beta_2 = \frac{A_2}{A_3} \end{cases} \quad (39)$$

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